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LOCALISATION IN THE BOUCHAUD-ANDERSON MODEL

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ABSTRACT. It is well-known that both random branching and trapping mechanisms can induce localisation phenomena in random walks; the prototypical examples being the parabolic Anderson and Bouchaud trap models respectively. Our aim is to investigate how these localisation phenomena interact in a hybrid model combining the dynamics of the parabolic Anderson and Bouchaud trap models. Under certain natural assumptions, we show that the localisation effects due to random branching and trapping mechanisms tend to (i) mutually reinforce, and (ii) induce a local correlation in the random fields (the 'fit and stable' hypothesis of population dynamics).

Minor revision to published version. This is an updated version of [23] containing the following minor revisions:

- A typo in the statement of Proposition 3.14 has been corrected;
- The coupling used in Section 5 has been slightly modified to fix a gap in its original statement; we thank Renato Soares dos Santos for pointing this out to us;
- A slight correction has been made to the proof of Proposition 4.11; and
- The bibliography has been updated.

1. INTRODUCTION

1.1. The Bouchaud–Anderson model. This paper studies a certain random walk model on \mathbb{Z}^d that is a hybrid of the well-known *parabolic Anderson* (PAM) and *Bouchaud trap* (BTM) models. To introduce this model, first recall the PAM, which describes the evolution of a diffusive particle in a random potential field (or, equivalently, a random branching environment; see below). Precisely, the PAM is the Cauchy problem on the lattice \mathbb{Z}^d

$$\frac{\partial u(t,z)}{\partial t} = (\Delta + \xi) u(t,z), \qquad (t,z) \in [0,\infty) \times \mathbb{Z}^d; \qquad (1)$$

$$u(0,z) = \mathbb{1}_{\{0\}}(z), \qquad z \in \mathbb{Z}^d;$$

where $\xi = \{\xi(z)\}_{z \in \mathbb{Z}^d}$ is a collection of independent identically distributed (i.i.d.) random variables known as the (random) *potential field* and Δ is the *discrete Laplacian* defined by $(\Delta f)(z) = \sum_{|y-z|=1} (2d)^{-1} (f(y) - f(z))$, where $|\cdot|$ denotes the ℓ_1 -norm. For a large class of potential field distributions,¹ equation (1) has a unique non-negative solution defined for all time t. For general background information on the PAM, including its origins in the statistical physics literature and its interpretation in terms of a system of branching diffusive particles, see [13].

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¹More specifically, those satisfying a certain integrability condition on the upper-tail; see [13].

Recall also the BTM, which describes the evolution of a diffusive particle in a random trapping landscape. Precisely, the BTM is the continuous-time Markov chain on \mathbb{Z}^d defined by the jump rates

$$w_{z \to y} := \begin{cases} (2d\sigma(z))^{-1}, & \text{if } |y - z| = 1, \\ 0, & \text{otherwise}, \end{cases}$$
(2)

where $\sigma = {\sigma(z)}_{z \in \mathbb{Z}^d}$ is a collection of strictly-positive i.i.d. random variables known as the (random) *trapping landscape*. Remark that the density of the BTM satisfies the equation

$$\frac{\partial u(t,z)}{\partial t} = \Delta \sigma^{-1} u(t,z), \qquad (t,z) \in [0,\infty) \times \mathbb{Z}^d, \qquad (3)$$

where, for clarity, we stress that the operator $\Delta \sigma^{-1}$ acts as

$$(\Delta \sigma^{-1} f)(z) = \sum_{|y-z|=1} (2d\sigma(y))^{-1} f(y) - \sigma^{-1}(z) f(z) \,.$$

For general background information on the BTM, including its origins in the study of spin-glasses dynamics and its broad utility as a simple model for a variety of trapping behaviour, see [4].

The PAM and BTM are of great interest in the theory of random processes because they exhibit *intermittency*, that is, unlike other commonly studied models of diffusion, their long-term behaviour cannot, in general, be described with a simple averaging principle (see [13] and [4] for a general overview of the PAM and BTM respectively.) Instead, extremes in the respective random environments may create concentration effects, which can result in the eventual *localisation* of the solution to equations (1) and (3) respectively over long periods of time. In the most extreme cases, the solution localises on just a few sites.

Our aim is to study how the localisation phenomena in the PAM and the BTM interact. To do this, we consider the Cauchy problem on the lattice \mathbb{Z}^d

$$\frac{\partial u(t,z)}{\partial t} = (\Delta \sigma^{-1} + \xi) u(t,z), \qquad (t,z) \in [0,\infty) \times \mathbb{Z}^d; \qquad (4)$$
$$u(0,z) = \mathbb{1}_{\{0\}}(z), \qquad z \in \mathbb{Z}^d;$$

derived by replacing the discrete Laplacian in equation (1) with the generator of the BTM in equation (3). We refer to equation (4) as the *Bouchaud–Anderson model* (BAM).

By analogy with the PAM (see [13], Section 1.2), the solution to equation (4) has a natural interpretation as the expected number of particles in a system of continuously-branching diffusive particles on the lattice \mathbb{Z}^d specified by:

- *Initialisation*: A single particle at the origin;
- Branching: The local branching rate for a particle at a site z is given by $\xi(z)$;
- Trapping: Each particle evolves as an independent BTM, that is, the waiting time at each visit to a site z is independent and distributed exponentially with mean $\sigma(z)$, with the subsequent site chosen uniformly from among the nearest neighbours.

This interpretation can be formalised in the Feynman-Kac representation of the solution to (4):

$$u(t,z) := \mathbb{E}_0 \left[\exp\left\{ \int_0^t \xi(X_s) ds \right\} \mathbb{1}_{\{X_t=z\}} \right],$$
(5)

where X is the BTM and, for $z \in \mathbb{Z}^d$, \mathbb{E}_z denotes the expectation over X given that $X_0 = z$. As we shall see, the interaction between the random branching and trapping mechanisms makes the localisation behaviour of the BAM highly non-trivial.

1.2. Localisation in the PAM and BTM. The PAM and BTM are said to *localise* if, as $t \to \infty$, the solution of equations (1) and (3) respectively are eventually concentrated on a small number of sites with overwhelming probability, i.e. if there exists a (random) *localisation set* Γ_t such that, as $t \to \infty$, $|\Gamma_t| = t^{o(1)}$ and

$$\frac{\sum_{z \in \Gamma_t} u(t, z)}{U(t)} \to 1 \qquad \text{in probability}, \tag{6}$$

where $U(t) := \sum_{z \in \mathbb{Z}^d} u(t, z)$ is the total mass of the solution (in the BTM, this is identically one); see Section 1.8 for the definition of the asymptotic notation used here and throughout the paper.

Naturally, the primary measure of the strength of localisation in the PAM and BTM is the cardinality of the localisation set Γ_t . As such, the most extreme form of localisation is *complete localisation*, which occurs if the total mass is eventually concentrated at just one site, i.e. if Γ_t can be chosen in equation (6) such that $|\Gamma_t| = 1$. A finer measure of the strength of localisation is the *radius of influence*, which measures the extent to which localisation sites themselves are determined by purely local features of the random environment. More precisely, the radius of influence ρ is the smallest integer for which the localisation sites can be determined by maximising a functional on \mathbb{Z}^d that depends on the random environments only through their values in balls of radius ρ around each site.

Broadly speaking, localisation in the PAM and BTM is generated by the structure-forming effects of extremes in the respective random environment. If these extremes are both sufficiently pronounced and sufficiently regular, over long periods of time the model will come to adopt the structure present in the environment, with localisation the most extreme manifestation of this. Naturally then, the strength of localisation in the PAM and BTM should depend on (i) the asymptotic rate of decay, and (ii) the regularity of the upper-tail of the random variables $\xi(0)$ and $\sigma(0)$. In this context, it is convenient to restrict $\xi(0)$ and $\sigma(0)$ to be strictly-positive and to characterise these random variables by their *exponential tail decay rate* function

$$g_{\xi}(x) := -\log(\mathbb{P}(\xi(0) > x)) \quad \text{and} \quad g_{\sigma}(x) := -\log(\mathbb{P}(\sigma(0) > x))$$

for then (i) and (ii) translate to the asymptotic growth and regularity of the non-decreasing functions g_{ξ} and g_{σ} .

We briefly outline some known results on localisation in the PAM and BTM. For simplicity, we shall assume all necessary regularity conditions without further specification.

1.2.1. Localisation in the parabolic Anderson model. The conditions under which the PAM completely localises in the sense of equation (6) has been the subject of intense and ongoing research over the last 25 years. The current understanding is that double-exponential tail decay $(g_{\xi}(x) \approx e^x)$ forms the boundary of the complete localisation universality class. More precisely, it is conjectured that the PAM exhibits complete localisation as long as $\log g_{\xi}(x) \ll x$. This has been proven (in [20]) in the extremal² case of Pareto-like tail decay $(g_{\xi}(x) \sim \gamma \log x, \text{ for } \gamma > d)$, and more recently (in [26] and [10]) in the case of Weibull-like tail decay $(g_{\xi}(x) \sim x^{\gamma})$. On the other hand, if $\log g_{\xi}(x) \gg x$, then complete localisation is known not to hold (see [12]). What occurs in the interface regime of double-exponential tail decay $(\log g_{\xi}(x) \sim cx, \text{ for } c > 0)$ is not currently well-understood.

As for the radius of influence of the potential field, ρ_{PAM} , in the case of Pareto-like tail decay it has been shown (see [20]) that $\rho_{\text{PAM}} = 0$, in other words, the localisation site can be determined by maximising a functional that depends on the potential field ξ only through its value at individual lattice sites, with interactions between neighbouring lattice sites having no influence on localisation. On the other hand, in the case of Weibull-like tail decay $(g_{\xi}(x) \sim x^{\gamma})$, the radius of influence has been shown (see [10]) to be $\rho_{\text{PAM}} = [(\gamma - 1)/2]^+$, where [x] and x^+ denote the integer and positive parts of x respectively. Clearly this implies that $\rho_{\text{PAM}} = 0$ if and only if $\gamma < 3$, and also that $\rho_{\text{PAM}} \to \infty$ in the $\gamma \to \infty$ limit.

1.2.2. Localisation in the Bouchaud trap model. The study of localisation in the Bouchaud trap model has also received considerable attention over the last 10 years. A notable feature of the BTM is that localisation can only occur in dimension one. In higher dimensions, the traps either have negligible effect in the limit (if the tail is integrable, by virtue of the law of large numbers), or are visited in such a way that their overall effect is spatially-homogeneous (see [11] and [4] for a proof of this result in the case of Pareto-like tail decay, although the result is thought to hold more generally for arbitrary non-integrable tail decay).

²This case is extremal in the sense that if $g_{\xi}(x) \sim \gamma \log x$ for $\gamma > d$ or $\gamma = d = 1$ then the solution to equation (1) 'blows-up' in finite time, see [13].

On the other hand, it is known that in dimension one, Pareto-like tail decay $(g_{\sigma}(x) \sim c \log x, c > 0)$ forms the boundary of the localisation universality class. More precisely, if $\log x = O(g_{\sigma}(x))$, it is known that the BTM does not localise in the sense of equation (6) (although it does localise in a certain weaker sense; see, e.g. [11]). On the other hand, it was proven in [22] that for sub-Pareto tail decay $(g_{\sigma}(x) \ll \log x)$, the BTM localises on exactly two-sites in the limit, with a radius of influence (i.e. of the trapping landscape) equal to 0.

1.3. **Overview of our results.** Before detailing our results in full, we first provide a brief overview to highlight salient features; this section is for exposition only, and is not intended to be mathematically rigorous. A complete description of our results follows in Section 1.7 below.

In this initial study of localisation in the BAM, we focus on the case where both potential distribution $\xi(0)$ and trap distribution $\sigma(0)$ have Weibull tail decay

$$\mathbb{P}(\xi(0) > x) = e^{-x^{\gamma}} \quad \text{and} \quad \mathbb{P}(\sigma(0) > x) = e^{-x^{\mu}} \qquad \gamma, \mu > 0 \,.$$

Our results also hold in the $\gamma, \mu \to 0$ limit (with some caveats; see Section 1.6). As we shall see, the BAM with Weibull tail decay turns out to be a natural regime to study, since the interaction between the potential field and trapping landscape exhibits certain phase transitions in (γ, μ) .

1.3.1. *Complete localisation*. Our first main result establishes the complete localisation of the BAM across the entire regime (see Theorem 1.7 below).

Theorem 1.1. There exists a (random) site Z_t such that, as $t \to \infty$,

$$rac{u(t, Z_t)}{U(t)} o 1$$
 in probability.

That the BAM completely localises for some (γ, μ) is expected, since the PAM with Weibull potential also exhibits complete localisation. More surprising, however, is that complete localisation occurs regardless of the presence of very large traps, even in dimension one, since *a priori* it might be thought that large traps would draw probability mass away from the localisation site.

1.3.2. Mutual reinforcement of localisation effects due to the PAM and the BTM. Since complete localisation holds in the entire regime, in order to probe the interaction between the potential field and trapping landscape we need a finer measure of localisation. Such a measure is provided by the radius of influence ρ , which as described above is the smallest integer for which the localisation site Z_t can be determined by maximising a functional on \mathbb{Z}^d that depends on ξ and σ only through their values in balls of radius ρ around each site. Our second main result is to determine the radius of influence ρ , and to prove its optimality (see Theorem 1.7 and part (a) of Theorem 1.10 below).

Theorem 1.2. The radius of influence is

$$\rho := \left[\frac{\gamma - 1}{2}\frac{\mu}{\mu + 1} + \frac{1}{2}\right]^+.$$

Note that ρ is a decreasing function of the strength of both the potential field and trapping landscape (i.e. an increasing function of γ and μ), in other words, the localisation effects due to the PAM and BTM are *mutually reinforcing*.

1.3.3. Reducibility of the BAM to the PAM. We next ask whether the BAM is 'reducible' to the PAM. There are actually two distinct notions of reducibility that are relevant. Strong reducibility describes the situation in which the trapping landscape σ plays no role in determining the localisation site Z_t , and the macroscopic behaviour of the system is adequately approximated by the PAM with potential ξ . Weak reducibility describes the situation in which all necessary information to determine Z_t is contained in the 'net growth rate' $\eta := \xi - \sigma^{-1}$, and moreover, the macroscopic behaviour of the BTM is adequately approximated by the PAM with potential replaced with η . The term 'net growth rate' comes from the interpretation of the BAM as a trapped, branching random walk (see Section 1.5 below). Our third main result is to determine the regimes in which the BAM is strongly and weakly reducibility to the PAM (see parts (c) and (d) of Theorem 1.10 below). These regimes are depicted in Figure 1.



FIGURE 1. Partition of the parameter space of the BAM according to the whether the BAM is 'strongly reducible' to the PAM with the usual potential ξ (left of the dashed line) or 'weakly reducible' to the PAM with the potential replaced with the 'net growth rate' η (left of the bold curve). The boundary curve is $\mu = 1/(\gamma - 2)$.

Theorem 1.3. The BAM is strongly reducible to the PAM if and only if $\gamma < 1$. The BAM is weakly reducible to the PAM if and only if $\rho = 0$ and $\gamma \ge 1$.

1.3.4. Local correlation between the potential field and trapping landscape: The 'fit and stable' hypothesis. Our final result is to establish the local correlation between the potential field and trapping landscape (where 'local' is from the perspective of the localisation site); this is the so-called 'fit and stable' hypothesis that has been predicted numerically in the mathematical biology literature (see, e.g., [6]), but never rigorously confirmed (see Section 1.5 below). Interestingly, the correlation that we observe is *positive* at the localisation site, but *negative* away from the localisation site, providing an unexpected extension to the 'fit and stable' hypothesis.

To describe this correlation, we shall need to define a second, possibly smaller, radius of influence

$$\rho_{\xi} := \left[\frac{\gamma - 1}{2} \frac{\mu}{\mu + 1}\right]^{+} \in \{\rho - 1, \rho\},\,$$

which is the the smallest integer for which the localisation site Z_t can be determined by a maximising a functional on \mathbb{Z}^d that depends on ξ only through its values in balls of radius ρ_{ξ} around each site (note, the functional must still depend on σ through balls of radius at least ρ). For simplicity, we exclude here the 'interface cases', i.e. the points of discontinuity of ρ_{ξ} .

Theorem 1.4. Assume that $\gamma \geq 1$, so that the BAM is not strongly reducible to the PAM. Let Z_t denote the site of complete localisation. Then, as $t \to \infty$ eventually almost surely: (i) the random variables $\xi(Z_t)$ and $\sigma(Z_t)$ are positively correlated; and (ii) for all z such that $0 < |z - Z_t| \leq \rho_{\xi}$, the random variables $\xi(z)$ and $\sigma(z)$ are negatively correlated.

In Theorem 1.9 below we make explicit the nature of this correlation, as well as providing a full description of the localisation site, determining its asymptotic distance from the origin, the local profile of the potential field and trapping landscape, and its ageing behaviour.

1.4. Methods and techniques. Our approach to proving localisation in the BAM is loosely based on existing techniques to prove localisation in the PAM (see, e.g., [1, 12, 26]), although the complex interaction between the potential field and the trapping landscape means that these techniques cannot be trivially adapted. Instead, the presence of the trapping landscape requires the development of existing techniques on two main fronts.

First, proving localisation in the BAM requires the development of the spectral theory of operators of the form $\Delta \sigma^{-1} + \xi$, including path expansions and Feynman-Kac representations for the principal eigenvalue and eigenfunction respectively. To our knowledge this theory has not appeared in the literature before, and may be of independent interest, including in the study of position-dependent mass Schrödinger operators (see Section 1.5 below). In the particular case of the BAM with Weibull tails, we also extend existing techniques to establish the max-class of local eigenvalues; this is necessary in order to extract extra information about the local correlation in the potential field and trapping landscape.

Second, in order to analyse the 'screening effect' of heavy traps, standard percolation estimates are insufficient: in dimension one, because of the geometry; in dimensions higher than one, because of complex dependencies between the potential field, the trapping landscape, and the localisation site Z_t . In dimension one we analyse heavy traps using coarse graining methods; in higher dimensions, we implement new ideas that allow us to apply percolation estimates in the presence of the dependencies.

In addition, our methods provide a new approach to working with 'cluster expansions'. Although these expansions have appeared in the literature before (see, e.g., [1, 10]), the standard approach has been to access them via resolvent formalism. Our techniques provides a purely probabilistic approach to 'cluster expansion', which avoids many of the technicalities of the resolvent formalism. One application would be a simpler, purely probabilistic proof of the localisation results on the PAM found in [10].

1.5. Connections to the literature. Although this is the first work to consider the BAM, there are clear connections between the BAM and other models in the literature. First, the BAM can be interpreted as the thermodynamic limit of a particle system with random branching and trapping mechanisms (given, respectively, by the potential field ξ and the trapping landscape σ). In the probability literature there have been several other analyses of models combining random branching and trapping mechanisms – in particular, trapping mechanisms given by asymmetric transition probabilities [3] and random conductances [28] – although these have not considered the localisation properties of the model, focusing instead on the growth of the total population.

Similar models have also appeared in the mathematical biological literature, where they find an application in the study of population dynamics. Here the branching and trapping rates are recast as the *fitness ('adaptedness')* and *stability ('adaptability')* respectively of individual states (e.g. geographic locations, genetic configurations etc.). While the literature contains several models which allow for randomness in either the fitness [19, 24] or stability [18, 21, 27], most relevant is [6] which considers a model in which *both* these characteristics vary. Indeed, the model considered in [6] is essentially identical to the BAM, except it is defined in a domain without any geometry: when an individual's state changes, the fitness and stability are re-sampled according to their respective distributions.³ The primary observation in [6] (obtained numerically) is the tendency of populations to concentrate on states which are both fit and stable: the 'fit and stable' hypothesis. Our results provide the first rigorous analysis of this phenomenon. Indeed, our results actually suggest a refinement of the hypothesis (for our model with geometry): that populations concentrate on states which are fit and stable, *but also* for which neighbouring states are both fit and *unstable*.

Second, operators of the form $\Delta \sigma^{-1} + \xi$ have important applications in quantum mechanics, since their eigenvalues give the energy levels of a particle whose effective mass is position-dependent (see, e.g., [7, 8, 25]). To make the connection, consider the position-dependent mass Schrödinger equation for a particle with effective mass σ in a potential field ξ . This equation has a Hamiltonian of general form (see [25])

$$\frac{1}{2} \left(\sigma^{-\alpha} \nabla \sigma^{-\beta} \nabla \sigma^{-\gamma} + \sigma^{-\gamma} \nabla \sigma^{-\beta} \nabla \sigma^{-\alpha} \right) + \xi, \qquad \alpha, \beta, \gamma \ge 0, \ \alpha + \beta + \gamma = 1.$$

 $^{^{3}}$ A second minor difference is that the population size is kept constant by the deletion of a uniformly chosen individual at each replication event.

Although there is no canonical choice for α, β, γ , in the discrete setting a natural restriction is $\beta = 0$, which avoids symmetry breaking in the definition of ∇ . Specialising to the case $\alpha = \gamma = 1/2$ gives the Hamiltonian

$$\sigma^{-\frac{1}{2}}\Delta\sigma^{-\frac{1}{2}} + \xi = \sigma^{-\frac{1}{2}} \left(\Delta\sigma^{-1} + \xi\right) \sigma^{\frac{1}{2}}.$$
(7)

We remark that the Hamiltonian in (7) is the 'symmetrised' form of the operator $\Delta \sigma^{-1} + \xi$, and hence has equivalent spectral theory. In Section 3 we develop general theory for operators of the form $\Delta \sigma^{-1} + \xi$, including deriving path expansions and Feynman-Kac representations for the principal eigenvalue and eigenfunction respectively. This section is entirely self-contained, and is completely deterministic, and we expect that it will be of independent interest.

Third, there are connections between the BAM and the PAM in the case where the potential field distribution $\xi(0)$ is allowed to take on highly negative (or even infinitely negative) values, which may be interpreted as 'traps'. Previous work has noted the minimal influence of such 'traps' in $d \geq 2$ (see, e.g. [13, Section 2.4]), essentially due to percolation estimates, an observation that finds echoes in our results and methods. However, there are clear differences between this model and the BAM, primarily due to the fact that the traps in the BAM may coexist with sites of high potential; this coexistence underlies the phenomena of mutual reinforcement and correlation that we observe in the BAM. On the other hand, in dimension one the effect of highly negative potential values in the PAM is significant (see [5]). Indeed, since such sites cannot be avoided, their effect is to 'screen' off the growth that would otherwise occur from sites of high potential, and so the asymptotic growth of the solution depends heavily on the relationship between the upper and lower tails of $\xi(0)$. Again, this is reminiscent of our results in dimension one, which are only valid if the trap distribution decays sufficiently fast to ensure 'screening' effects are negligible.

1.6. The formal set-up for the paper. For the rest of the paper, we make the following assumptions on the potential field ξ and the trapping landscape σ :

Assumption 1.5 (Assumption on the potential field distribution). The random variable $\xi(0)$ is strictly-positive and satisfies

$$\bar{F}_{\xi}(x) = e^{-x^{\gamma}} \,,$$

for some $\gamma > 0$, where $\bar{F}_{\xi}(x) := 1 - F_{\xi}(x) := \mathbb{P}(\xi(0) > x)$.

Assumption 1.6 (Assumptions on the trap distribution). The random variable $\sigma(0)$ satisfies:

(a) No quick sites: The quantity

$$\delta_{\sigma} := \operatorname{essinf} \sigma(0)$$

is strictly positive;

(b) Regularity: The quantity

$$\mu := \lim_{x \to \infty} \frac{\log g_{\sigma}(x)}{\log x}$$

exists and is finite. If $\mu > 0$, then $\sigma(0)$ has a continuous density function $f_{\sigma}(x)$ with a Weibull upper-tail, i.e. for sufficiently large x,

$$\bar{F}_{\sigma}(x) = \exp\{-x^{\mu}\},\,$$

where $\bar{F}_{\sigma}(x) := 1 - F_{\sigma}(x) := \mathbb{P}(\sigma(0) > x)$. If $\mu = 0$, then $\sigma(0)$ has a continuous density function $f_{\sigma}(x)$, with the property that

$$f_{\sigma}(a_x) \sim f_{\sigma}(b_x)$$

for any $a_x, b_x \to \infty$ such that $a_x \sim b_x$ (see Section 1.8 for the asymptotic notation). In both cases, the lower-tail of f_{σ} satisfies, as $x \to 0$,

$$f_{\sigma}(x+\delta_{\sigma}) = o(e^{-1/x})$$

Furthermore, if d = 1, then additionally $\sigma(0)$ satisfies the following two extra conditions:

(c) Sufficiently fast tail decay: As $x \to \infty$ eventually, for some $\varepsilon > 0$,

$$g_{\sigma}(x) > (1 + \varepsilon) \log \log x;$$

(d) **Regularity**: There exists a $c \in (1, \infty]$ such that

$$\lim_{x \to \infty} \frac{g_{\sigma}(x)}{\log \log x} = c$$

with the convergence eventually monotone in the case $c = \infty$.

We wish to briefly comment on the nature of the above assumptions on $\xi(0)$ and $\sigma(0)$. First, we claim that the BAM with Weibull potential field is a natural regime in which to observe the interaction between the localisation effects in the PAM and the BTM. If the potential field is any stronger (indeed if $\gamma < 1$), the BAM is strongly reducible to the PAM⁴. On the other hand, if the potential field is any weaker, the effect of the trapping landscape, while present, is harder to measure. To see why, recall that the PAM with Weibull potential field has been shown to completely localise with a certain finite radius of influence ρ_{PAM} ; it is on the level of this radius that we measure the impact of the trapping landscape σ . Since $\rho_{\text{PAM}} \to \infty$ in the $\gamma \to \infty$ limit, the effect of changes to ρ_{PAM} become harder to quantify, and we leave this study to future work.

Second, the regularity assumption on $\xi(0)$ is imposed mainly for simplicity; weaker regularity assumptions (like those found in [1] and [2] for instance) are possible, although they introduce certain technical difficulties that we wish to avoid. Finally, note that equivalent results for the BAM with Pareto-like potential field can be naturally deduced by considering our results in the $\gamma \rightarrow 0$ limit.

Turning to the assumptions on $\sigma(0)$, first note that the quantity μ measures the 'Weibullness' of the upper-tail of $\sigma(0)$, with the case $\mu = 0$ corresponding to a stronger-than-Weibull trapping landscape. For simplicity, we have chosen not to consider weaker-than-Weibull trapping landscapes in this paper; equivalent results can be naturally deduced by considering our results in the $\mu \to \infty$ limit. As with $\xi(0)$, the regularity assumptions on $\sigma(0)$ are certainly not optimal for our results to hold; they are chosen mainly for simplicity. On the other hand, our assumption that $\sigma(0)$ is bounded away from zero is essential. Indeed we expect that the nature of the localisation behaviour will change if 'quick' sites are present. Finally, the additional tail decay assumption in dimension one is also essential, and our results and methods break down completely without it. Note, however, that this condition is only violated for trap distributions with extremely heavy tails, such as if $\sigma(0)$ is a *log-Pareto* random variable.

1.7. Full description of our results. Here we describe our results in full, expanding on the exposition given in Section 1.3. In order to state our results explicitly, we shall need to introduce some notation. Recall the parameter $\mu \in [0, \infty)$ from Assumption 1.6, which describes the 'Weibull-like' decay parameter of the upper-tail of $\sigma(0)$. Recall also from Section 1.3 the radius of influence

$$\rho := \left[\frac{\gamma - 1}{2}\frac{\mu}{\mu + 1} + \frac{1}{2}\right]^+$$

and the, possibly smaller, radius of influence of the potential field ξ ,

$$\rho_{\xi} := \left[\frac{\gamma - 1}{2}\frac{\mu}{\mu + 1}\right]^+ \in \{\rho - 1, \rho\} \le \rho$$

The relationship between ρ and ρ_{ξ} is depicted in Figure 2; we defer further discussion on ρ and ρ_{ξ} to Remark 1.11.

Next we describe explicitly the localisation site. For each $z \in \mathbb{Z}^d$ and $n \in \mathbb{N}$, define the ball $B(z,n) := \{y \in \mathbb{Z}^d : |y-z| \le n\}$. For each $z \in \mathbb{Z}^d$, define the Hamiltonian

$$\mathcal{H}(z) := \Delta \sigma^{-1} + \xi \mathbb{1}_{B(z,\rho_{\mathcal{E}})}$$

 $^{^{4}}$ Note however that, because of Assumption 1.6, this conclusion does not apply in dimension one if the trapping landscape is sufficiently strong.



FIGURE 2. Partition of the parameter space of the BAM according to the values of ρ (bold lines) and ρ_{ξ} (dashed lines). The boundary curves are of the form $\mu = (2i-1)/(\gamma - 2i)$ and $\mu = (2i)/(\gamma - 2i - 1)$, for $i \in \mathbb{N} \setminus \{0\}$.

restricted to the domain $B(z, \rho)$ with Dirichlet boundary conditions, denoting by $\lambda(z)$ its principal eigenvalue. Note that each $\lambda(z)$ is real since the Hamiltonian $\mathcal{H}(z)$ is similar to the Hermitian operator

$$\sigma^{-\frac{1}{2}} \mathcal{H}(z) \, \sigma^{\frac{1}{2}} = \sigma^{-\frac{1}{2}} \Delta \sigma^{-\frac{1}{2}} + \xi \mathbb{1}_{B(z,\rho_{\varepsilon})} \, .$$

We refer to $\lambda(z)$ as the *local principal eigenvalue at z*, and remark that it is a certain functional of the sets $\xi^{(\rho_{\xi})}(z) := \{\xi(y)\}_{y \in B(z,\rho_{\xi})}$ and $\sigma^{(\rho)}(z) := \{\sigma(y)\}_{y \in B(z,\rho)}$. Note that the random variables $\{\lambda(z)\}_{z \in \mathbb{Z}^d}$ are identically distributed, and have a dependency range bounded by 2ρ , i.e. the random variables $\lambda(y)$ and $\lambda(z)$ are independent if and only if $|y-z| > 2\rho$. Remark also that in the special case $\rho = 0$, $\lambda(z)$ reduces to the 'net growth rate' $\eta(z) = \xi(z) - \sigma^{-1}(z)$.

For any sufficiently large t, define a *penalisation functional* $\Psi_t : \mathbb{Z}^d \to \mathbb{R}$ by

$$\Psi_t(z) := \lambda(z) - \frac{|z|}{\gamma t} \log \log t.$$

Note that Ψ_t has a similar form to the penalisation functional introduced in [10] to prove complete localisation in the PAM with Weibull potential field, representing the trade-off between energetic forces (given by the local principal eigenvalue $\lambda(z)$) and entropic forces (given by a probabilistic penalty which is linear in |z| and decaying in t); see Remark 1.8.

Define a large 'macrobox' $V_t := [-R_t, R_t]^d \cap \mathbb{Z}^d$, with $R_t := t(\log t)^{\frac{1}{\gamma}}$. Fix a constant $0 < \theta < 1/2$ and define the macrobox level $L_t := ((1 - \theta) \log |V_t|)^{\frac{1}{\gamma}}$. Let the subset $\Pi^{(L_t)} := \{z \in \mathbb{Z}^d : \xi(z) > L_t\} \cap V_t$ consist of sites in V_t at which ξ -exceedances of the level L_t occur. Finally, define the random site

$$Z_t := \arg\max_{z \in \Pi^{(L_t)}} \Psi_t(z) \,.$$

The site Z_t is well-defined eventually almost surely since, as we show in Lemma 4.1, the set $\Pi^{(L_t)}$ is non-empty and finite eventually almost surely. Moreover, for t sufficiently large, Z_t almost surely does not depend on the particular choice of θ . We present again (see Theorem 1.1) our main theorem:

Theorem 1.7 (Complete localisation). As $t \to \infty$,

$$\frac{u(t, Z_t)}{U(t)} \to 1$$
 in probability.

Remark 1.8. In order to determine Z_t explicitly, a finite approximation is available for $\lambda(z)$ (see Proposition 5.1 for a precise formulation):

$$\lambda(z) \approx \eta(z) + \sigma^{-1}(z) \sum_{\substack{2 \le k \le 2j \\ p_i \ne z, \ 0 < i < k \\ \text{Set}(p) \subseteq B(z, \rho)}} \prod_{\substack{0 < i < k \\ \text{Set}(p) \subseteq B(z, \rho)}} \prod_{\substack{0 < i < k \\ \text{Set}(p) \subseteq B(z, \rho)}} (2d)^{-1} \frac{\sigma^{-1}(p_i)}{\lambda(z) - \eta_z(p_i)},$$
(8)

where $j := [\gamma - 1]$ and $\eta_z := \xi \mathbb{1}_{B(z,\rho_{\xi})} - \sigma^{-1}$; see Section 1.8 for the definition of the path set $\Gamma_k(z,z)$. This path expansion can be iteratively evaluated to approximate $\Psi_t(z)$ as an explicit function of $\xi^{(\rho_{\xi})}(z)$, $\sigma^{(\rho)}(z)$, |z| and t, which, as we show, is sufficiently precise to determine the localisation site Z_t with overwhelming probability.

Before stating our second and third main results we shall introduce some more notation. First we define exponents that explicitly describe the correlation of the fields ξ and σ around the localisation site Z_t . To this end, define the function $q_{\xi} : \mathbb{N} \to [0, 1]$ and the non-negative constant q_{σ} by

$$q_{\xi}(x) := \begin{cases} \left(1 - \frac{2x}{\gamma - 1} - \frac{1}{\mu + 1}\right)^{+} & \text{if } \gamma > 1, \\ (1 - x)^{+} & \text{else,} \end{cases} \quad \text{and} \quad q_{\sigma} := \left(\frac{\gamma - 1}{\mu + 1}\right)^{+}$$

We shall also need the concept of 'interface cases', which correspond to the values of (γ, μ) where ρ , and respectively ρ_{ξ} , are transitioning from one integer to the next. To this end define the sets

$$\mathcal{B} := \left\{ (\gamma, \mu) : \frac{\gamma - 1}{2} \frac{\mu}{\mu + 1} + \frac{1}{2} = \rho \right\} \text{ and } \mathcal{B}_{\xi} := \left\{ (\gamma, \mu) : \frac{\gamma - 1}{2} \frac{\mu}{\mu + 1} = \rho_{\xi} \right\}.$$

Note that these sets correspond, respectively, to the bold and dashed curves in Figure 2. Finally, define the random time $T_t := \inf\{s > 0 : Z_{t+s} \neq Z_t\}$ and the scales

$$r_t := \frac{t(d\log t)^{\frac{1}{\gamma}-1}}{\log\log t} \quad \text{and} \quad a_t := (d\log t)^{\frac{1}{\gamma}}.$$
(9)

The scales r_t and a_t describe, respectively, the scale of the distance from the origin of the localisation site and the scale of the height of the potential field at the localisation site.

Theorem 1.9 (Description of the localisation site). As $t \to \infty$ the following hold:

(a) (Localisation distance)

$$\frac{Z_t}{r_t} \Rightarrow X \qquad in \ law,$$

where X is a random vector whose coordinates are independent and distributed as Laplace (two-sided exponential) random variables with absolute-moment one.

(b) (Local correlation of the potential field) If $(\gamma, \mu) \notin \mathcal{B}_{\xi}$, then for each $z \in B(0, \rho_{\xi})$ there exists a c > 0 such that

$$\frac{\xi(Z_t+z)}{a_t^{q_\xi(|z|)}} \to c \qquad in \ probability.$$
(10)

If $(\gamma, \mu) \in \mathcal{B}_{\xi}$, then (10) holds for each $z \in B(0, \rho_{\xi} - 1)$, and moreover, for each z such that $|z| = \rho_{\xi}$ there exists a c > 0 such that,

$$f_{\xi(Z_t+z)}(x) \to \frac{e^{cx}f_{\xi}(x)}{\mathbb{E}[e^{c\xi(0)}]},$$

uniformly over $x \in (0, L_t)$, where $f_{\xi(z)}$ is the density of the potential field at site z (see Assumption 1.5).

(c) (Correlation of the trapping landscape at Z_t) If $\mu > 0$ and $\gamma > 1$, then there exists a c > 0 such that

$$\frac{\sigma(Z_t)}{a_t^{q_\sigma}} \to c \qquad in \ probability \,.$$

If $\mu = 0$ and $\gamma > 1$ then, for each $\nu > 0$,

$$\mathbb{P}\left(\frac{\log \sigma(Z_t)}{\log a_t} > q_\sigma - \nu\right) \to 1.$$

If $\gamma = 1$ then,

$$f_{\sigma(Z_t)}(x) \to \frac{e^{-1/x} f_{\sigma}(x)}{\mathbb{E}[e^{-1/\sigma(0)}]},$$

uniformly over x, where $f_{\sigma(Z_t)}$ is the density of the trapping landscape at site Z_t . (d) (Local correlation of the trapping landscape) If $(\gamma, \mu) \notin \mathcal{B}$, then for each $z \in B(0, \rho) \setminus \{0\}$

 $\sigma(Z_t + z) \to \delta_\sigma$ in probability. (11)

If $(\gamma, \mu) \in \mathcal{B}$, then (11) holds for each $z \in B(0, \rho) \setminus \{0\}$ and moreover, for each z such that $|z| = \rho$, there exists a c > 0 such that

$$f_{\sigma(Z_t+z)}(x) \to \frac{e^{c/x} f_{\sigma}(x)}{\mathbb{E}[e^{c/\sigma(0)}]},$$

uniformly over x, where $f_{\sigma(z)}$ is the density of the trapping landscape at site z (see Assumption 1.6).

(e) (Ageing)

$$\frac{T_t}{t} \Rightarrow \Theta \qquad in \ law$$

where Θ is a non-degenerate almost surely positive random variable.

Theorem 1.10 (Optimality results). As $t \to \infty$ the following hold:

(a) (Optimality of the radius of influence) The radius of influence ρ is optimal, in other words, there does not exist a functional ψ_t , depending on ξ and σ only through their values in balls of radius $\rho - 1$ around each site z, such that

$$\mathbb{P}\left(Z_t = \operatorname*{arg\,max}_{z \in \mathbb{Z}^d} \psi_t(z)\right) \to 1.$$
(12)

(b) (Optimality of the radius of influence with respect to the potential field) The radius of influence of the potential field ρ_{ξ} is optimal, in other words, there does not exist a functional ψ_t , depending on ξ only through its values in balls of radius $\rho_{\xi} - 1$ around each site z, such that

$$\mathbb{P}\left(Z_t = \operatorname*{arg\,max}_{z \in \mathbb{Z}^d} \psi_t(z)\right) \to 1.$$
(13)

(c) (Criterion for reduction to the potential ξ) The localisation site is independent of the trapping landscape σ if and only if $\gamma < 1$, in other words, if and only if $\gamma < 1$, there exists a random site $z_t \in \mathbb{Z}^d$, independent of σ , such that,

$$\mathbb{P}\left(Z_t = z_t\right) \to 1. \tag{14}$$

(d) (Criterion for reduction to the 'net growth rate' η) The localisation site Z_t depends on ξ and σ only through the value of η if and only if $\rho = 0$, in other words, if and only if $\rho = 0$, there exists a random site $z_t \in \mathbb{Z}^d$, dependent on ξ and σ only through η , such that,

$$\mathbb{P}\left(Z_t = z_t\right) \to 1. \tag{15}$$

Remark 1.11. We note several interesting features of the radius of influence ρ . As remarked above, ρ is an increasing function of both γ and μ . Moreover, surprisingly it is not necessarily the case that $\rho \to \rho_{\text{PAM}} := [(\gamma - 1)/2]^+$ in the $\mu \to \infty$ limit; indeed, if $\gamma \in (2i, 2i + 1)$ for $i \in \mathbb{N} \setminus \{0\}$, then in fact $\rho \to \rho_{\text{PAM}} + 1$, meaning that influence of the BTM on the BAM is not continuous in the degenerate limit (i.e. as $\sigma(z) \to 1$ simultaneously for each z). On the other hand, $\rho_{\xi} \to \rho_{\text{PAM}}$ in the $\mu \to \infty$ limit, i.e. there is no discontinuity in the effect of the BTM on the BAM on the level of the radius of influence of the potential field ξ . The relationship between ρ , ρ_{ξ} and ρ_{PAM} is depicted in Figure 3.



FIGURE 3. Partition of the parameter space of the BAM according to the relationship between ρ (bold lines) and ρ_{PAM} (dashed lines), and ρ_{ξ} (bold lines) and ρ_{PAM} (dashed lines) respectively, where ρ_{PAM} denotes the radius of influence in the equivalent PAM with identical potential field. The boundary curves are of the form $\mu = (2i-1)/(\gamma - 2i)$ and $\mu = (2i)/(\gamma - 2i - 1)$ respectively, for $i \in \mathbb{N} \setminus \{0\}$.

Remark 1.12. The shape of the local profile of the potential field and trapping landscape in parts (b)–(d) of Theorem 1.9 is derived by considering the path expansion in equation (8) and determining the values of ξ and σ that appropriately balance: (i) the increase in λ gained from favourable realisations of ξ and σ ; and (ii) the probabilistic penalty that results from such favourable realisations of ξ and σ if they are too unlikely. This balance is expressed through a convex function whose integral is asymptotically concentrated in the regions specified in Theorem 1.9. This computation is carried out in the proof of Proposition 5.3, identifying the constants in Theorem 1.9 explicitly.

We also give some heuristic reasons why we must distinguish the cases $(\gamma, \mu) \in \mathcal{B}, \mathcal{B}_{\xi}$ in the correlation results. If $(\gamma, \mu) \notin \mathcal{B}_{\xi}$, then the value of $\xi(Z_t + z)$ is growing (with high probability) as $t \to \infty$ for each $z \in B(0, \rho_{\xi})$. However, if $(\gamma, \mu) \in \mathcal{B}_{\xi}$, this only occurs for $z \in B(0, \rho_{\xi} - 1)$; at the interface of the radius, where $|z| = \rho_{\xi}$, the value of $\xi(Z_t + z)$ instead converges to a certain random variable with law distinct from the law of $\xi(0)$. Similarly, for $(\gamma, \mu) \notin \mathcal{B}, \sigma(Z_t + z)$ converges to δ_{σ} for each $z \in B(0, \rho) \setminus \{0\}$. However, if $(\gamma, \mu) \in \mathcal{B}$, then this is only true for $z \in B(0, \rho - 1) \setminus \{0\}$. If $|z| = \rho$, the value of $\sigma(Z_t + z)$ instead converges to a certain random variable with law distinct from the law of $z \in B(0, \rho) \setminus \{0\}$. These properties are reflective of the fact that the correlation in the fields ξ and σ induced by the localisation site Z_t decays away from the site.

We also explain why the cases $\gamma \leq 1$ and $\mu = 0$ must be further distinguished in our profile for $\sigma(Z_t)$. If $\gamma > 1$ then the value of $\sigma(Z_t)$ is growing, and indeed growing with a deterministic leading order. However, if $\gamma = 1$, this is no longer true and instead $\sigma(Z_t)$ converges to a certain random variable with law distinct from the law of $\sigma(0)$.⁵ The case $\mu = 0$ must be distinguished for a different reason; in this case, the extremes of σ are so large that there are many sites z for which $\sigma^{-1}(z)$ is smaller than the gap in the top statistics of Ψ_t . Past this threshold, differences in the magnitude of σ no longer materially influence the determination of Z_t , and so we lose a degree of certainty about the order of growth of $\sigma(Z_t)$.

Note finally that if (γ, μ) is not in \mathcal{B} and \mathcal{B}_{ξ} respectively, then the probabilities in equations (12) and (13) actually converge to 0 for any such ψ_t ; otherwise, the respective probability will converge to a constant $c \in (0, 1)$. Similarly, if (γ, μ) lies to the right of the dashed or bold line in Figure 1, the probabilities in (14) and (15) respectively converge to 0 for any such z_t ; if (γ, μ)

⁵Of course, in the case $\gamma < 1$, with overwhelming probability σ is independent of the localisation site Z_t (cf. part (c) of Theorem 1.9) and so $\sigma(Z_t)$ has the same law as $\sigma(0)$.

lies on either line, the repective probability instead converges to a constant $c \in (0, 1)$. We do not prove these additional results here.

1.8. Notation. Here we list notation that will be commonly used for the remainder of the paper. Asymptotic notation: For functions f and g we use $f \sim g$ to denote that

$$\lim_{x \to \infty} f(x)/g(x) = 1 \,,$$

and f = o(g) or $f \ll g$ to denote that

$$\lim_{x \to \infty} f(x)/g(x) = 0$$

We use f = O(g) to denote that, as $x \to \infty$, eventually for some constant c > 0,

$$\left|f(x)\right| < c|g(x)|\,.$$

Notation for paths: For an integer k and sites $y, z \in \mathbb{Z}^d$, let $\Gamma_k(y, z)$ be the set of nearest neighbour paths in \mathbb{Z}^d of length k running from y to z, with each $p \in \Gamma_k(y, z)$ indexed as

$$y =: p_0 \rightarrow p_1 \rightarrow p_2 \rightarrow \ldots \rightarrow p_k := z.$$

Similarly, denote

$$\Gamma_{k}(y) := \bigcup_{z \in \mathbb{Z}^{d}} \Gamma_{k}(y, z) , \quad \Gamma(y, z) := \bigcup_{k \in \mathbb{N}} \Gamma_{k}(y, z) ,$$
$$\Gamma(y) := \bigcup_{k \in \mathbb{N}} \Gamma_{k}(y) , \quad \Gamma := \bigcup_{y \in \mathbb{Z}^{d}} \Gamma(y) .$$

For a site $z \in \mathbb{Z}^d$, let n(z) denote the number of shortest paths from the origin to z, i.e.

$$n(z) := |\Gamma_{|z|}(0,z)|.$$

For a path $p \in \Gamma_k(y, z)$ denote $\operatorname{Set}(p) := \{p_0, p_1, \ldots, p_k\}$ and |p| := k. For a nearest neighbour random walk X let $p(X_t) \in \Gamma(X_0)$ denote the geometric path associated with the trajectory of $\{X_s\}_{s \leq t}$ and let $p_k(X) \in \Gamma_k(X_0)$ denote the geometric path associated with the random walk $\{X_s\}_{s \geq 0}$ up to and including its k^{th} jump.

Notation for sets: For a domain $D \in \mathbb{Z}^d$, denote by

$$\partial D = \{y \in D^c : \text{there exists } x \in D \text{ such that } |x - y| = 1\}.$$

For a set $S \in \mathbb{Z}^d$ define $B(S, n) := \bigcup_{z \in S} B(z, n)$ and $\sup(S) := \min_{x \neq y \in S} \{|x - y|\}.$

Notation for solutions of the BAM: For each $y, z \in \mathbb{Z}^d$ define $u_y(t, z)$ to be the solution of the Cauchy problem

$$\frac{\partial u_y(t,z)}{\partial t} = (\Delta \sigma^{-1} + \xi) u_y(t,z), \qquad (t,z) \in [0,\infty) \times \mathbb{Z}^d;$$
$$u_y(0,z) = \mathbb{1}_{\{y\}}(z), \qquad z \in \mathbb{Z}^d;$$

and, for $z \in \mathbb{Z}^d$ and $p \in \Gamma$, define

$$u^{p}(t,z) := \mathbb{E}_{p_{0}}\left[\exp\left\{\int_{0}^{t} \xi(X_{s}) \, ds\right\} \mathbb{1}_{\{X_{t}=z\}} \mathbb{1}_{\{p(X_{t})=p\}}\right], \quad U^{p}(t) := \sum_{z \in \mathbb{Z}^{d}} u^{p}(t,z).$$

Notation for local principal eigenvalues: For each $z \in \mathbb{Z}^d$ and $n \in \mathbb{N}$ define the *n*-local principal eigenvalue $\lambda^{(n)}(z)$ to be the principal eigenvalue of the Hamiltonian

$$\mathcal{H}^{(n)}(z) := \Delta \sigma^{-1} + z$$

restricted to the domain B(z, n) with Dirichlet boundary conditions.

Other notation: For $a, b \in \mathbb{R}$ define $a \wedge b := \min\{a, b\}$.

2. Outline of proof

The main idea of the proof of Theorem 1.7 is that the solution u(t, z) can be decomposed into disjoint components by reference to the trajectories of the underlying BTM in the Feynman-Kac representation in (5). Using such a path decomposition, we prove complete localisation by establishing that: (i) a single component carries the entire non-negligible part of the solution; and (ii) the non-negligible component is localised at Z_t .

To assist in the proof, we introduce the scale

$$d_t := \frac{1}{\gamma} (d\log t)^{\frac{1}{\gamma} - 1} \tag{16}$$

which is the derivative (on the log scale) of the height scale a_t , and naturally examines the gaps in the maximisers of ξ in growing boxes. We also introduce auxiliary scaling functions $f_t, h_t, e_t, b_t \to 0$ and $g_t, s_t \to \infty$ as $t \to \infty$ that are convenient placeholders for negligibly decaying (respectively growing) functions. For technical reasons, we shall need to choose these functions to satisfy certain relationships, as follows. First, let s_t be such that

$$(\log s_t)^2 \ll \log \log t$$
.

Then, choose f_t, h_t, e_t, b_t and g_t satisfying

$$\max\{\bar{F}_{\sigma}(s_t), (\log s_t)^2 / \log \log t, 1 / \log \log s_t\} g_t \ll b_t \ll f_t h_t \ll g_t h_t \ll e_t.$$

$$(17)$$

It is easy to check that such a choice is always possible.

Path decomposition. We explain here how to construct the path decomposition. Recall the definition of V_t from Section 1.7. For a path $p \in \Gamma(0)$ such that $\operatorname{Set}(p) \subseteq V_t$, let

$$z^{(p)} := \underset{z \in \operatorname{Set}(p)}{\operatorname{arg\,max}} \lambda(z)$$

which is well-defined almost surely. Abbreviate

$$B_t := B(0, |Z_t|(1+h_t)) \cap V_t.$$

We partition the path set $\Gamma(0)$ into the following five disjoint components

$$E_t^i := \begin{cases} \left\{ p \in \Gamma(0) : \operatorname{Set}(p) \subseteq B_t, \, z^{(p)} = Z_t \right\}, & i = 1 ; (\operatorname{non-negligible \ component}) \\ \left\{ p \in \Gamma(0) : \operatorname{Set}(p) \subseteq V_t, \, z^{(p)} \in \Pi^{(L_t)} \setminus Z_t \right\}, & i = 2 ; (\operatorname{path \ does \ not \ hit \ best \ site}) \\ \left\{ p \in \Gamma(0) : \operatorname{Set}(p) \subseteq V_t, \, \operatorname{Set}(p) \not\subseteq B_t, \, z^{(p)} = Z_t \right\}, & i = 3 ; (\operatorname{path \ travels \ far}) \\ \left\{ p \in \Gamma(0) : \operatorname{Set}(p) \subseteq V_t, \, z^{(p)} \notin \Pi^{(L_t)} \right\}, & i = 4 ; (\operatorname{path \ avoids \ all \ high \ sites}) \\ \left\{ p \in \Gamma(0) : \operatorname{Set}(p) \not\subseteq V_t \right\}, & i = 5 ; (\operatorname{path \ leaves \ macrobox}) \end{cases}$$

and associate each component E_t^i with a portion of the total mass U(t) of the solution. As such, for each $1 \le i \le 5$, let

$$u^i(t,z) := \sum_{p \in E^i_t} u^p(t,z)$$
 and $U^i(t) = \sum_{z \in \mathbb{Z}^d} u^i(t,z)$.

Our strategy is to establish that each of $U^2(t)$, $U^3(t)$, $U^4(t)$ and $U^5(t)$ are negligible with respect to the total mass U(t) of the solution, in other words that,

$${U^i(t)\over U(t)}
ightarrow 0$$
 in probability, for $i=2,3,4,5$.

To complete the proof of localisation, we also prove that $U^{1}(t)$ is localised at Z_{t} , i.e. that,

$$\frac{u^1(t, Z_t)}{U^1(t)} \to 1 \qquad \text{in probability.}$$

Note that this strategy requires a balance to be struck in how B_t is defined; it must be large enough that $U^3(t)$ is negligible, but small enough to ensure localisation. The scale h_t has been fine-tuned in (17) precisely to ensure this balance is struck correctly. Negligible paths. The negligibility of $U^4(t)$ and $U^5(t)$ are simple to establish; the main difficulty is establishing the negligibility of $U^2(t)$ and $U^3(t)$. Our proof is based on formalising two heuristics. **First heuristic**: Recall the constant $j := [\gamma - 1] \ge \rho$ and the definition of the *j*-local principal eigenvalue $\lambda^{(j)}$ from Section 1.8. We expect that, for a path $p \in \Gamma(0) \setminus E_t^5$,

$$U^{p}(t) \approx \exp\left\{t\lambda^{(j)}(z^{(p)})\right\} a_{t}^{-|p|} , \qquad (18)$$

which represents the balance between (i) the exponential growth of the solution at each site, and (ii) the probabilistic penalty for travelling each step along the path p.

The accuracy of this heuristic relies on some subtle observations about the BAM (and indeed the PAM) which we shall briefly discuss further. First is the claim that the *j*-local principal eigenvalues closely approximate the exponential growth rate of the solution at a site (note that here we could take a slightly smaller constant in place of *j*, but *j* will turn out to be convenient for another reason; see immediately below). This approximation, in turn, is based on the fact that there is a lack of resonance between the top eigenvalues of the operator $\Delta + \xi$ restricted to any finite domain.

Second is the claim that it is never beneficial for a path to visit other sites of high potential, other than $z^{(p)}$. This is proved by way of a 'cluster expansion' (see Lemma 3.13) which bounds the contribution to $U^p(t)$ between the time p visits a site z of high potential until it leaves the ball B(z, j). Crucially, j is chosen precisely to be the smallest integer for which this 'cluster expansion' bound is smaller than the probabilistic penalty associated with the path getting from outside the ball B(z, j) to z (see the proof of Proposition 6.7).

Third is the claim that the probabilistic penalty for travelling along the path p is approximately $1/a_t$ for each step of the path. Implicit in this claim is the highly non-trivial fact that the trapping landscape σ is not an obstacle to the diffusivity of the particle, in other words, that a sufficiently 'quick' path exists from 0 to the site $z^{(p)}$. If $d \ge 2$, this is essentially due to percolation estimates; if d = 1, then this relies crucially on the additional tail decay assumption on the distribution of $\sigma(0)$, and our proofs and methods break down without it.

Second heuristic: We expect that, for i = 1, 2, 3,

$$U^{i}(t) \approx \max_{p \in E^{i}_{t}} U^{p}(t) , \qquad (19)$$

which represents the idea that $U^i(t)$ should be dominated by the contribution from just a single path in the path set E_t^i . This is essentially due to the fact that the number of paths of length kgrows exponentially in k, whereas the probabilistic penalty associated with a path of length kdecays as a_t^{-k} , which dominates since $a_t \to \infty$.

Let us consider what these heuristics imply for U(t). Recall the definition of $\Pi^{(L_t)}$ from Section 1.7. By analogy with Ψ_t and Z_t , define

$$\Psi_t^{(j)}(z) = \lambda^{(j)}(z) - \frac{|z|}{\gamma t} \log \log t$$

and $Z_t^{(j)} := \arg \max_{z \in \Pi^{(L_t)}} \Psi_t^{(j)}$. Note that it will turn out that $Z_t^{(j)} = Z_t$ with overwhelming probability (see Corollary 5.11), so we will interchange between these freely in the discussion that follows. Clearly, by the two heuristics, the dominant contribution to U(t) will come from a path $p \in \Gamma(0)$ that goes directly from the origin to $z^{(p)}$, and so we expect that

$$U(t) \approx \max_{p \in \Gamma(0)} \left\{ \exp\left\{ t\lambda^{(j)}(z^{(p)}) \right\} a_t^{-|z^{(p)}|} \right\} \approx \exp\left\{ t \max_{z \in \mathbb{Z}^d} \Psi_t^{(j)}(z) \right\} = \exp\left\{ t\Psi_t^{(j)}(Z_t^{(j)}) \right\}.$$

Indeed, we formalise this approximation as a lower bound

$$\log U(t) \ge t \Psi_t^{(j)}(Z_t^{(j)}) + O(td_t b_t) \,. \tag{20}$$

Similarly for $U^2(t)$, the heuristics imply that the dominant contribution will come from the path $p \in E_t^2$ that goes directly from the origin to the site

$$Z_t^{(j,2)} = \underset{z \in \Pi^{(L_t)} \setminus \{Z_t^{(j)}\}}{\arg \max} \Psi_t^{(j)}(z) \,,$$

and so

$$U^{2}(t) \approx \exp\left\{t\lambda^{(j)}(Z_{t}^{(j,2)})\right\} a_{t}^{-|Z_{t}^{(j,2)}|} \approx \exp\left\{t\Psi_{t}^{(j)}(Z_{t}^{(j,2)})\right\} \,.$$

We formalise this approximation as an upper bound

$$\log U^{2}(t) \le t \Psi_{t}^{(j)}(Z_{t}^{(j,2)}) + O(td_{t}b_{t})$$

which, together with equation (20), implies that

$$\log U^{2}(t) - \log U(t) \leq -t \left(\Psi_{t}^{(j)}(Z_{t}^{(j)}) - \Psi_{t}^{(j)}(Z_{t}^{(j,2)}) + O(d_{t}b_{t}) \right) \,.$$

Remark that the negligibility of $U^2(t)$ is then a consequence of the gap in the top order statistics of $\Psi_t^{(j)}$ being larger than the error (of order $O(d_t b_t)$) in these bounds.

Finally, the heuristics imply that the dominant contribution to $U^{3}(t)$ will come from a path p that visits Z_{t} but that also ventures outside B_{t} , and so

$$U^{3}(t) \approx \exp\left\{t\lambda^{(j)}(Z_{t})\right\}a_{t}^{-|Z_{t}|(1+h_{t})}$$

We formalise this approximation as an upper bound

$$\log U^{3}(t) \le t\lambda^{(j)}(Z_{t}) - \frac{1}{\gamma}|Z_{t}|(1+h_{t})\log\log t + O(td_{t}b_{t})$$
(21)

which, together with equation (20), implies that

$$\log U^3(t) - \log U(t) \le -\frac{1}{\gamma} |Z_t| h_t \log \log t + O(td_tb_t).$$

Remark that the negligibility of $U^3(t)$ is then a consequence of $|Z_t|h_t \log \log t$ being larger than the error (also of order $O(td_tb_t)$) in these bounds.

In Section 5 we study extremal theory for $\lambda^{(j)}$ and $\Psi_t^{(j)}$, demonstrating, in particular, that

$$\Psi_t^{(j)}(Z_t^{(j)}) - \Psi_t^{(j)}(Z_t^{(j,2)}) > d_t e_t \quad \text{and} \quad |Z_t^{(j)}| h_t \log \log t > t d_t e_t$$

both hold eventually with overwhelming probability. We also show that $Z_t^{(j)} = Z_t$ with overwhelming probability. In the process, we establish the description of the localisation site Z_t that is contained in Theorem 1.9, as well as the optimality results in Theorem 1.10. In Section 6, we show how to formalise the heuristics in equations (18) and (19) into the bounds in equations (20) and (21), and so complete the proof of the negligibility of $U^2(t)$ and $U^3(t)$.

Throughout, we draw on the preliminary results established in Sections 3 and 4. Section 3 contains a compilation of general results on operators of the form $\Delta \sigma^{-1} + \xi$. Section 4 contains general results on the random fields ξ and σ . Of particular concern here is the existence of 'quick' paths through the trapping landscape σ .

Localisation. In Section 7 we complete the proof of Theorem 1.7 by showing that $u^1(t, z)$ is localised at the site Z_t . The main idea is the same as in [12] and [26], namely that: (i) the solution $u^1(t, z)$ is asymptotically approximated by the principal eigenfunction of the operator $\Delta \sigma^{-1} + \xi$ restricted to the domain B_t ; and (ii) the principal eigenfunction decays exponentially away from the site Z_t . Underlying this reasoning is the fact that the domain B_t has been constructed to ensure that $\lambda^{(j)}(Z_t)$ is the largest *j*-local principal eigenvalue in the domain. This in turn allows us to give a Feynman-Kac representation of the principal eigenfunction v_t (see Proposition 7.3), which we analyse probabilistically to establish exponential decay.

3. General theory for the BAM

In this section we develop general theory for operators of the form $\Delta \sigma^{-1} + \xi$ which is valid for arbitrary ξ and positive σ . This section will be entirely self-contained and is completely deterministic, and may be of independent interest. We have chosen to develop the theory in full generality so as to take advantage of the results in future work. Throughout this section, let $D \subset \mathbb{Z}^d$ be a bounded domain and let ξ and σ be arbitrary functions $\xi : \mathbb{Z}^d \to \mathbb{R}$ and $\sigma : \mathbb{Z}^d \to \mathbb{R}^+$, with $\eta := \xi - \sigma^{-1}$. Denote by \mathcal{H} the Hamiltonian

$$\mathcal{H} := \Delta \sigma^{-1} + \xi$$

restricted to the domain D with Dirichlet boundary conditions, and let $\{\lambda_i\}_{i \leq |D|}$ and $\{\varphi_i\}_{i \leq |D|}$ be respectively the (finite) set of eigenvalues and eigenfunctions of \mathcal{H} , with eigenvalues in descending order and eigenfunctions ℓ_2 normalised. Finally, recall that X_s denotes the BTM and define the stopping times

$$\tau_z := \inf\{t \ge 0 : X_t = z\} \text{ and } \tau_{D^c} := \inf\{t \ge 0 : X_t \notin D\}.$$

We start by presenting representations and deriving simple bounds for λ_1 and φ_1 .

Lemma 3.1 (Principal eigenvalue monotonicity). For each $z \in D$ and $\delta > 0$, let $\overline{\lambda}_1$ be the principal eigenvalue of the operator $\mathcal{H} + \delta \mathbb{1}_{\{z\}}$. Then $\overline{\lambda}_1 > \lambda_1$.

Moreover, for each strictly smaller domain $\overline{D} \subset D$, let $\overline{\lambda}_1$ be the principal eigenvalue of \mathcal{H} restricted to the domain \overline{D} with Dirichlet boundary conditions. Then $\overline{\lambda}_1 < \lambda_1$.

Proof. These are general properties of elliptic operators.

Lemma 3.2 (Bounds on the principal eigenvalue).

$$\max_{z \in D} \{\eta(z)\} \le \lambda_1 \le \max_{z \in D} \left\{ \eta(z) + \sum_{|y-z|=1} \frac{1}{2d} \sigma^{-1}(y) \right\} \,.$$

Proof. The lower bound follows from the min-max theorem for the principal eigenvalue; the upper bound follows from the Gershgorin circle theorem. \Box

Proposition 3.3 (Feynman-Kac representation for the principal eigenfunction). For each $y, z \in D$ the principal eigenfunction φ_1 satisfies the Feynman-Kac representation

$$\frac{\varphi_1(y)}{\varphi_1(z)} = \frac{\sigma(y)}{\sigma(z)} \mathbb{E}_y \left[\exp\left\{ \int_0^{\tau_z} \left(\xi(X_s) - \lambda_1 \right) \, ds \right\} \mathbb{1}_{\{\tau_{D^c} > \tau_z\}} \right]. \tag{22}$$

Proof. Consider z fixed and define $v^z(y) := \varphi_1(y)/\varphi_1(z)$. Note that the function v^z satisfies the Dirichlet problem

$$(\Delta \sigma^{-1} + \xi - \lambda_1) v^z(y) = 0, \qquad \qquad y \in D \setminus \{z\};$$
$$v^z(y) = \mathbb{1}_{\{z\}}(y), \qquad \qquad y \notin D \setminus \{z\}.$$

It is easy to check (for instance, by integrating over the first holding time) that the Feynman-Kac representation on the right-hand side of equation (22) also satisfies this Dirichlet problem; hence we are done if there is a unique solution. So assume another non-trivial solution w exists. Then the difference $q := v^z - w$ satisfies the Dirichlet problem

$$(\Delta \sigma^{-1} + \xi - \lambda_1) q(y) = 0, \qquad \qquad y \in D \setminus \{z\};$$

$$q(y) = 0, \qquad \qquad y \notin D \setminus \{z\};$$

which is nonzero if and only if λ_1 is an eigenvalue of the operator $\Delta \sigma^{-1} + \xi$ restricted to the domain $D \setminus \{z\}$ with Dirichlet boundary conditions. By the domain monotonicity of the principal eigenvalue in Lemma 3.1, this is impossible.

Lemma 3.4 (Path-wise evaluation). For each $k \in \mathbb{N}$, $y, z \in D$, $p \in \Gamma_k(z, y)$ such that $p_i \neq y$ for i < k and $\operatorname{Set}(p) \subseteq D$, and $\zeta > \max_{0 \le i \le k} \eta(p_i)$, we have

$$\mathbb{E}_{z}\left[\exp\left\{\int_{0}^{\tau_{y}}(\xi(X_{s})-\zeta)\,ds\right\}\,\mathbb{1}_{\{p_{k}(X)=p\}}\right] = \prod_{i=0}^{k-1}\frac{1}{2d}\frac{\sigma^{-1}(p_{i})}{\zeta-\eta(p_{i})}$$

Proof. This follows by integrating over the holding times at the sites $\{p_i\}_{0 \le i \le k-1}$, which are independent. The restriction on ζ ensures that the resulting integrals are finite.

Proposition 3.5 (Path expansion for the principal eigenvector). For each $y, z \in D$ the principal eigenfunction φ_1 satisfies the path expansion

$$\frac{\varphi_1(y)}{\varphi_1(z)} = \frac{\sigma(y)}{\sigma(z)} \sum_{k \ge 1} \sum_{\substack{p \in \Gamma_k(y,z) \\ p_i \ne z, 0 \le i < k \\ \operatorname{Set}(p) \subseteq D}} \prod_{0 \le i < k} \frac{1}{2d} \frac{\sigma^{-1}(p_i)}{\lambda_1 - \eta(p_i)} \,.$$

Proof. The expectation on the right-hand side of equation (22) can be expanded path-wise using Lemma 3.4, which is valid by the lower bound in Lemma 3.2. \Box

Remark 3.6. Note that the initial factor $\sigma(y)$ in the above path expansion cancels with the term $\sigma^{-1}(p_1)$ appealing in each component of the sum. This turns out to be crucial in establishing the localisation of the eigenfunctions in Sections 6 and 7, since a priori $\sigma(y)$ could be arbitrarily large.

Proposition 3.7 (Path expansion for the principal eigenvalue). For each $z \in D$ the principal eigenvalue has the path expansion

$$\lambda_1 = \eta(z) + \sigma^{-1}(z) \sum_{k \ge 2} \sum_{\substack{p \in \Gamma_k(z,z) \\ p_i \ne z, \ 0 < i < k \\ \operatorname{Set}(p) \subseteq D}} \prod_{0 < i < k} \frac{1}{2d} \frac{\sigma^{-1}(p_i)}{\lambda_1 - \eta(p_i)}.$$

Proof. Recalling that the eigenfunction relation evaluated at a site z gives

$$\lambda_1 = \eta(z) + \sum_{|y-z|=1} \sigma^{-1}(y) \frac{\varphi_1(y)}{\varphi_1(z)},$$

the result follows from Proposition 3.5.

We now study the solution $u_z(t, y)$ to the Cauchy problem

$$\frac{\partial u_z(t,y)}{\partial t} = \mathcal{H} u(t,y) , \qquad (t,y) \in [0,\infty) \times D ; \qquad (23)$$
$$u_z(0,y) = \mathbb{1}_{\{z\}}(y) , \qquad y \in \mathbb{Z}^d .$$

 $\langle \rangle$

In particular, we give the spectral representation of $u_z(t, y)$ and deduce upper and lower bounds.

Proposition 3.8 (Feynman-Kac representation of the solution). For each $y, z \in D$,

$$u_z(t,y) = \mathbb{E}_z \left[\exp\left\{ \int_0^t \xi(X_s) ds \right\} \mathbb{1}_{\{X_t=y\}} \mathbb{1}_{\{\tau_{D^c} > t\}} \right]$$

Proof. It can be directly verified that the Feynman-Kac representation satisfies (23).

Lemma 3.9 (Time-reversal). For each $y, z \in D$,

$$u_z(t, y) \sigma(z) = u_y(t, z) \sigma(y)$$

Proof. Consider the Hermitian operator

$$\tilde{\mathcal{H}} := \sigma^{-\frac{1}{2}} \mathcal{H} \sigma^{\frac{1}{2}} = \sigma^{-\frac{1}{2}} \Delta \sigma^{-\frac{1}{2}} + \xi$$

which can be viewed as the 'symmetrised' form of \mathcal{H} . Since,

$$e^{\tilde{\mathcal{H}}t} = e^{\sigma^{-\frac{1}{2}}\mathcal{H}\sigma^{\frac{1}{2}}t} = \sigma^{-\frac{1}{2}}e^{\mathcal{H}t}\sigma^{\frac{1}{2}}$$

we have, by the fact that $\tilde{\mathcal{H}}$ is Hermitian,

$$u_{z}(t,y) = e^{\mathcal{H}t} \mathbb{1}_{\{z\}}(y) = \left(\frac{\sigma(y)}{\sigma(z)}\right)^{\frac{1}{2}} e^{\tilde{\mathcal{H}}t} \mathbb{1}_{\{z\}}(y) = \left(\frac{\sigma(y)}{\sigma(z)}\right)^{\frac{1}{2}} e^{\tilde{\mathcal{H}}t} \mathbb{1}_{\{y\}}(z)$$
$$= \frac{\sigma(y)}{\sigma(z)} e^{\mathcal{H}t} \mathbb{1}_{\{y\}}(z) = \frac{\sigma(y)}{\sigma(z)} u_{y}(t,z) \,.$$

,

Proposition 3.10 (Spectral representation for the solution). For each $y, z \in D$, the solution $u_z(t, y)$ satisfies the spectral representation

$$u_z(t,y) = \sigma^{-1}(z) \sum_i \frac{e^{\lambda_i t} \varphi_i(z) \varphi_i(y)}{||\sigma^{-\frac{1}{2}} \varphi_i||_{\ell_2}^2}$$

Proof. Recall the Hermitian operator $\tilde{\mathcal{H}}$ from the proof of Lemma 3.9. Note that each (ℓ_2 normalised) eigenfunction $\tilde{\varphi}_i$ of $\tilde{\mathcal{H}}$ satisfies the relation

$$\tilde{\varphi}_i = \frac{\sigma^{-\frac{1}{2}}\varphi_i}{||\sigma^{-\frac{1}{2}}\varphi_i||_{\ell_2}}$$

with λ_i the corresponding eigenvalue for $\tilde{\varphi}_i$. The proof then follows by applying the spectral theorem to $\tilde{\mathcal{H}}$.

Corollary 3.11 (Bounds on the solution). For each $z \in D$ we have the bounds

$$\frac{e^{\lambda_1 t} \sigma^{-1}(z) \varphi_1^2(z)}{||\sigma^{-\frac{1}{2}} \varphi_1||_{\ell_2}^2} \le u_z(t,z) \le e^{\lambda_1 t}$$

Proof. The lower bound follows directly from Proposition 3.10. For the upper bound, first use Proposition 3.10 to write

$$u_{z}(t,z) \leq e^{\lambda_{1}t}\sigma^{-1}(z)\sum_{i}\frac{\varphi_{i}^{2}(z)}{\|\sigma^{-\frac{1}{2}}\varphi_{i}\|_{\ell_{2}}^{2}}.$$

Then, since $u_z(0, z) = 1$, Proposition 3.10 also implies that

$$\sigma^{-1}(z) \sum_{i} \frac{\varphi_i^2(z)}{\|\sigma^{-\frac{1}{2}}\varphi_i\|_{\ell_2}^2} = 1$$

and the result follows.

Proposition 3.12 (Bound on the total mass of the solution). For each $y, z \in D$,

$$\sum_{y \in D} u_z(t, y) \le e^{\lambda_1 t} \sum_{y \in D} \frac{\varphi_1(y)}{\varphi_1(z)}$$

Proof. We write \mathcal{F}_{τ_z} for the σ -algebra generated by the stopping time τ_z . First decompose the Feynman-Kac representation for $u_y(t, z)$ in Proposition 3.8 by conditioning on \mathcal{F}_{τ_z} and using the strong Markov property:

$$\begin{aligned} u_y(t,z) &= \mathbb{E}_{\tau_z} \left[e^{\lambda_1 \tau_z} \mathbb{E}_y \left[\exp\left\{ \int_0^{\tau_z} \left(\xi(X_s) - \lambda_1 \right) \, ds \right\} \, \mathbb{1}_{\{\tau_z < \tau_{D^c}\}} \middle| \mathcal{F}_{\tau_z} \right] \right] \\ &\times \mathbb{E}_z \left[\exp\left\{ \int_0^{t - \tau_z} \xi(X'_s) \, ds \right\} \, \mathbb{1}_{\{X'_{t - \tau_z} = z, \tau'_{D^c} > t - \tau_z\}} \middle| \mathcal{F}_{\tau_z} \right] \, \mathbb{1}_{\{\tau_z \le t\}} \right] \\ &= \mathbb{E}_{\tau_z} \left[e^{\lambda_1 \tau_z} \, \mathbb{E}_y \left[\exp\left\{ \int_0^{\tau_z} \left(\xi(X_s) - \lambda_1 \right) \, ds \right\} \, \mathbb{1}_{\{\tau_z < \tau_{D^c}\}} \middle| \mathcal{F}_{\tau_z} \right] \, u_z(t - \tau_z, z) \, \mathbb{1}_{\{\tau_z \le t\}} \right] \,, \end{aligned}$$

where \mathbb{E}_{τ_z} denotes expectation taken over τ_z , X'_t is an independent copy of X_t , and $\tau'_{D^c} := \inf\{t \ge 0 : X'_t \notin D\}$. Using the upper bound in Corollary 3.11 combined with the Feynman-Kac representation for the principal eigenfunction in Proposition 3.3, we have that

$$u_y(t,z) \le e^{\lambda_1 t} \mathbb{E}_y \left[\exp\left\{ \int_0^{\tau_z} \left(\xi(X_s) - \lambda_1 \right) \, ds \right\} \, \mathbb{1}_{\{\tau_z < \tau_{D^c}\}} \right] = e^{\lambda_1 t} \frac{\varphi_1(y)}{\varphi_1(z)} \frac{\sigma(z)}{\sigma(y)} \, .$$

Finally, applying the time-reversal Lemma 3.9, we have

$$u_z(t,y) = u_y(t,z) \frac{\sigma(y)}{\sigma(z)} \le e^{\lambda_1 t} \frac{\varphi_1(y)}{\varphi_1(z)}$$

which, after summing over $y \in D$, yields the result.

Next we prove a 'cluster expansion' that is useful for bounding expectations of the 'Feynman-Kac type'. It is similar in spirit to [12, Lemma 4.2] and [14, Lemma 2.18], however we will need an additional form of the bound to accommodate the impact of the trapping landscape (see the proof of Lemma 7.4).

Lemma 3.13 (Cluster expansion). For each $z \in D$ and for any $\zeta > \lambda_1$,

$$\mathbb{E}_{z}\left[\exp\left\{\int_{0}^{\tau_{D^{c}}}\left(\xi(X_{s})-\zeta\right)ds\right\}\right] \leq 1+\frac{\max_{w\in D}\left\{\sigma^{-1}(w)\right\}\left|D\right|}{\zeta-\lambda_{1}}$$

and

$$\mathbb{E}_{z}\left[\exp\left\{\int_{0}^{\tau_{D^{c}}}(\xi(X_{s})-\zeta)\,ds\right\}\right] \leq \frac{\sigma^{-1}(z)}{\zeta-\lambda_{1}}\left(1+\frac{\max_{w\in D}\{\sigma^{-1}(w)\}\,|D|}{\zeta-\lambda_{1}}\right)\,.$$

Proof. We proceed by modifying the proofs of [12, Lemma 4.2] and [14, Lemma 2.18]. First abbreviate

$$u(y) := \mathbb{E}_y \left[\exp\left\{ \int_0^{\tau_{D^c}} (\xi(X_s) - \zeta) \, ds \right\} \right]$$

and remark that u solves the boundary value problem

$$(\sigma^{-1}\Delta + \xi - \zeta)u(y) = 0, \qquad \qquad y \in D; \qquad (24)$$
$$u(y) = 1, \qquad \qquad y \notin D.$$

Note that, in contrast to in the proof of Proposition 3.3, in the above boundary value problem the relevant operator is the adjoint of \mathcal{H} , since here we have not weighted the expectation by σ . We make the substitution w := u - 1, where 1 denotes the vector of ones, which turns (24) into

$$\begin{aligned} (\sigma^{-1}\Delta + \xi - \zeta)w(y) &= -(\sigma^{-1}\Delta + \xi - \zeta)\mathbb{1}(y) = \zeta - \xi(y), \qquad & y \in D; \\ w(y) &= 0, \qquad & y \notin D. \end{aligned}$$

Since $\zeta > \lambda_1$, the solution exists and is given by

$$w(y) = \left(\mathcal{R}_{\zeta}(\xi - \zeta)\right)(y)$$

where \mathcal{R}_{ζ} is the resolvent of $\sigma^{-1}\Delta + \xi$ at ζ . By Lemma 3.2 and since $\zeta > \lambda_1$ we have that $\xi(y) - \zeta \leq \sigma^{-1}$ for all $y \in D$, and so by the positivity of the resolvent (guaranteed since \mathcal{H} is elliptic and $\zeta > \lambda_1$) we obtain

$$w(z) \le \left(\mathcal{R}_{\zeta} \sigma^{-1}\right)(y) = \left(\sigma^{-\frac{1}{2}} \tilde{\mathcal{R}}_{\zeta} \sigma^{-\frac{1}{2}}\right)(y) \le \max_{z \in D} \{\sigma^{-1}(z)\} |D| \left\| \tilde{\mathcal{R}}_{\zeta} \right\|,$$

where $\tilde{\mathcal{R}}_{\zeta}$ is the resolvent of the Hermitian operator $\tilde{\mathcal{H}} = \sigma^{-\frac{1}{2}} \Delta \sigma^{-\frac{1}{2}} + \xi$ at ζ and $\|\cdot\|$ denotes the operator norm. By considering the spectral representation of $\tilde{\mathcal{R}}_{\zeta}$ we have $\|\tilde{\mathcal{R}}_{\zeta}\| \leq (\zeta - \lambda_1)^{-1}$ which gives the first bound. For the second bound, consider that (24) implies the identify

$$u(y) = \frac{\sigma^{-1}(y)}{\zeta - \xi(y) + \sigma^{-1}(y)} \sum_{|x-y|=1} \frac{1}{2d} u(x).$$
(25)

Applying the first bound to each u(x) in the sum in (25), the result follows by bounding $\xi(y) - \sigma^{-1}(y)$ in the denominator of (25) from above by λ_1 , valid by the lower bound in Lemma 3.2. \Box

Finally, we give a general way to bound the contribution to the solution $u_z(t, y)$ from paths that hit a certain site $x \in D$ and then stay within a subdomain $E \subseteq D$ that contains x. In particular, we show that this contribution is proportional to the principal eigenfunction of \mathcal{H} restricted to E. This is similar in spirit to [12, Theorem 4.1], and it crucial to establishing complete localisation of the solution.

So fix a domain $E \subseteq D$, a site $x \in E$, and define the operator \mathcal{H}^E to be the restriction of \mathcal{H} to the domain E with Dirichlet boundary conditions, with λ_1^E and φ_1^E respectively its principal eigenvalue and eigenfunction. Define the stopping time

$$\tau_{x,E^c} := \inf\{t \ge \tau_x : X_t \notin E\}.$$

Then the contribution to the solution $u_z(t, y)$ from paths that hit x and then stay within E can be written

$$u_{z}^{x,E}(t,y) := \mathbb{E}_{z} \left[\exp \left\{ \int_{0}^{t} \xi(X_{s}) \, ds \right\} \mathbb{1}_{\{X_{t}=y,\tau_{x} \leq t,\tau_{x,E^{c}} > t,\tau_{D^{c}} > t\}} \right].$$

Proposition 3.14 (Link between solution and principal eigenfunction; see [12, Theorem 4.1]). For each $x \in E$, $y \in E \setminus \{x\}$ and $z \in D$,

$$\frac{u_z^{x,E}(t,y)}{\sum_{y\in D} u_z(t,y)} \le \frac{\sigma(x) \|\sigma^{-\frac{1}{2}}\varphi_1^E\|_{\ell_2}^2}{(\varphi_1^E(x))^3} \varphi_1^E(y) \,.$$

Proof. We proceed by modifying the proof of [12, Theorem 4.1]. The first step is to make use of the time-reversal in Lemma 3.9, suitably adapted to $u_z^{x,E}(t,y)$. In particular, defining

$$u_{y}^{\overleftarrow{x,E}}(t,z) := \mathbb{E}_{y} \left[\exp \left\{ \int_{0}^{t} \xi(X_{s}) \, ds \right\} \mathbb{1}_{\{X_{t}=z,\tau_{x} \leq t, \tau_{x} < \tau_{E^{c}}, \tau_{D^{c}} > t\}} \right]$$

we can write

$$\frac{u_z^{x,E}(t,y)}{\sum_{y\in D} u_z(t,y)} \le \frac{u_z^{x,E}(t,y)}{u_z(t,x)} = \frac{\sigma(y)}{\sigma(x)} \frac{u_y^{\overleftarrow{x,E}}(t,z)}{u_x(t,z)}.$$
(26)

Next we decompose the Feynman-Kac formula for $u_y^{\overleftarrow{x,E}}(t,z)$ as in the proof of Proposition 3.12, by conditioning on the σ -algebra generated by the stopping time τ_x , and using the strong Markov property. More precisely, we write

$$u_{y}^{\overleftarrow{\tau,E}}(t,z) = \mathbb{E}_{\tau_{x}} \left[e^{\tau_{x}\lambda_{1}^{E}} \mathbb{E}_{y} \left[\exp\left\{ \int_{0}^{\tau_{x}} \left(\xi(X_{s}) - \lambda_{1}^{E} \right) \, ds \right\} \mathbb{1}_{\{\tau_{x} < \tau_{E}c\}} \middle| \mathcal{F}_{\tau_{x}} \right] \right]$$

$$\times \mathbb{E}_{x} \left[\exp\left\{ \int_{0}^{t - \tau_{x}} \xi(X_{s}') \, ds \right\} \mathbb{1}_{\{X_{t - \tau_{x}}' = z, \tau_{D}'c > t - \tau_{x}\}} \middle| \mathcal{F}_{\tau_{x}} \right] \mathbb{1}_{\{\tau_{x} \le t\}} \right],$$

$$(27)$$

where \mathbb{E}_{τ_x} , X'_t and τ'_{D^c} are defined as in the proof of Proposition 3.12. Next, note that an application of Corollary 3.11 gives the bound

$$1 \le u_x^{x,E}(w,x) \ \frac{\sigma(x) \|\sigma^{-\frac{1}{2}} \varphi_1^E\|_{\ell_2}^2}{(\varphi_1^E(x))^2} \ e^{-w\lambda_1^E}, \tag{28}$$

and recall the representation

$$u_x^{x,E}(w,x) = \mathbb{E}_x \left[\exp\left\{ \int_0^w \xi(X'_s) \, ds \right\} \mathbb{1}_{\left\{ X'_w = x, \tau'_{E^c} > w \right\}} \right] \,.$$

Combining the bound in (28) with equation (27) (setting $w = \tau_x$), gives

$$\begin{split} u_{y}^{\overleftarrow{\tau,E}}(t,z) &\leq \frac{\sigma(x) \| \sigma^{-\frac{1}{2}} \varphi_{1}^{E} \|_{\ell_{2}}^{2}}{(\varphi_{1}^{E}(x))^{2}} \mathbb{E}_{\tau_{x}} \left[\mathbb{E}_{y} \left[\exp \left\{ \int_{0}^{\tau_{x}} \left(\xi(X_{s}) - \lambda_{1}^{E} \right) \, ds \right\} \mathbb{1}_{\{\tau_{E^{c}} > \tau_{x}\}} \right| \mathcal{F}_{\tau_{x}} \right] \\ &\times \mathbb{E}_{x} \left[\exp \left\{ \int_{0}^{\tau_{x}} \xi(X_{s}') \, ds \right\} \mathbb{1}_{\{X_{\tau_{x}}' = x, \tau_{E^{c}}' > \tau_{x}\}} \right| \mathcal{F}_{\tau_{x}} \right] \\ &\times \mathbb{E}_{x} \left[\exp \left\{ \int_{0}^{t - \tau_{x}} \xi(X_{s}') \, ds \right\} \mathbb{1}_{\{X_{t-\tau_{x}}' = z, \tau_{D^{c}}' > t - \tau_{x}\}} \right| \mathcal{F}_{\tau_{x}} \right] \mathbb{1}_{\{\tau_{x} \leq t\}} \right] \\ &\leq \frac{\sigma(x) \| \sigma^{-\frac{1}{2}} \varphi_{1}^{E} \|_{\ell_{2}}^{2}}{(\varphi_{1}^{E}(x))^{2}} \mathbb{E}_{\tau_{x}} \left[\mathbb{E}_{y} \left[\exp \left\{ \int_{0}^{\tau_{x}} \left(\xi(X_{s}) - \lambda_{1}^{E} \right) \, ds \right\} \mathbb{1}_{\{\tau_{E^{c}} > \tau_{x}\}} \right| \mathcal{F}_{\tau_{x}} \right] \\ &\times \mathbb{E}_{x} \left[\exp \left\{ \int_{0}^{t} \xi(X_{s}') \, ds \right\} \mathbb{1}_{\{X_{t}' = z, \tau_{D^{c}}' > t\}} \right| \mathcal{F}_{\tau_{x}} \right] \mathbb{1}_{\{\tau_{x} \leq t\}} \right] \\ &\leq \frac{\sigma(x)^{2} \| \sigma^{-\frac{1}{2}} \varphi_{1}^{E} \|_{\ell_{2}}^{2}}{\sigma(y) (\varphi_{1}^{E}(x))^{3}} \varphi_{1}^{E}(y) \ u_{x}(t,z) \,, \end{split}$$

where the inequality in the second step results from deleting the condition that $X'_{\tau_x} = x$, and where the last inequality results from deleting the condition that $\tau_x \leq t$, and where we have used the Feynman-Kac representation for φ_1^E given by Proposition 3.3. Combining this with equation (26) gives the result.

4. PROPERTIES OF THE RANDOM ENVIRONMENTS

In this section we establish properties of the i.i.d. fields ξ and σ . In the first part we give asymptotics for the upper order statistics of ξ and σ . The second part is devoted to proving the existence of 'quick paths', which are an essential part of our proof that the trapping landscape does not prevent complete localisation in the BAM.

4.1. Almost sure asymptotics for ξ and σ . For each $a \leq 1$, define the macrobox level $L_{t,a} := ((1-a)\log|V_t|)^{\frac{1}{\gamma}}$ and let the subset $\Pi^{(L_{t,a})} := \{z \in \mathbb{Z}^d : \xi(z) > L_{t,a}\} \cap V_t$ consist of sites in V_t at which ξ -exceedances of the level $L_{t,a}$ occur. Recall that $L_t := L_{t,\theta}$.

Lemma 4.1 (Almost sure asymptotics for ξ). Denote by $\xi_{t,i}$ the *i*th highest value of ξ in V_t . Then for $a \in [0,1)$ and $a' \in (0,1]$, as $t \to \infty$,

$$\xi_{t,[|V_t|^a]} \sim L_{t,a}$$
 and $|\Pi^{(L_{t,a'})}| \sim |V_t|^{a'}$

hold almost surely.

Proof. These follow from well-known results on sequences of i.i.d. random variables; they are proved in a similar way as [17, Lemma 4.7].

Recall that for a set $S \in \mathbb{Z}^d$ we denote by $\operatorname{sep}(S) := \min_{x \neq y \in S} \{|x - y|\}.$

Lemma 4.2 (Almost sure separation of high points of ξ). For any a > 0 and $n \in \mathbb{N}$ let

$$\Pi_n^{(L_t,a)} := \{ z \in B(V_t, n) : \xi(z) > L_{t,a} \}$$

be the set of $L_{t,a}$ exceedences of ξ in the n-extended macrobox $B(V_t, n)$. Then, for any a' < a, as $t \to \infty$

$$sep\left(\Pi_n^{(L_{t,a})} \cup \{0\}\right) > |V_t|^{\frac{1-2a'}{d}}$$

eventually almost surely.

Proof. This result is proved as in [1, Lemma 1].

Remark 4.3. Note that we need the almost sure separation of high points in the *n*-extended macrobox $B(V_t, n)$ rather than just in V_t because each $\lambda^{(n)}(z)$, for $z \in V_t$, depends on the random environments ξ and σ in the ball $B(z, n) \subseteq B(V_t, n)$. This result implies that, eventually almost surely, each $z \in \Pi^{(L_{t,a})}$ has the property that $\xi(y) < L_{t,a}$ for all $y \in B(z, n) \setminus \{z\}$.

Corollary 4.4 (Paths cannot always remain close to high points of ξ). There exists a $c \in (0, 1)$ such that, for each $n \in \mathbb{N}$, all paths $p \in \Gamma(0, z)$ such that $\operatorname{Set}(p) \subseteq V_t$ satisfy, as $t \to \infty$,

$$\left|\left\{i: p_i \notin B(\Pi^{(L_t)}, n)\right\}\right| \ge |z| - \frac{|z|}{t^c},$$

eventually almost surely.

Proof. Abbreviate $N := sep(\Pi^{(L_t)} \cup \{0\})$ and

$$Q := \left| \left\{ i : p_i \notin B(\Pi^{(L_t)}, n) \right\} \right| \,.$$

Suppose a path p passes through m distinct B(x,n) with $x \in \Pi^{(L_t)}$. Then, since there is a minimum distance of (N-2n) between each such ball, the path p satisfies

$$Q \ge m(N-2n) \,.$$

On the other hand, it is clear that that $Q \ge |z| - (2n+1)m$. Therefore

$$Q \ge \min_{m \in \mathbb{N}} \max\left\{m(N-2n-1), |z| - (2n+1)m\right\} \ge \frac{(N-2n-1)|z|}{N} = |z| - \frac{(2n+1)|z|}{N}$$

and the result follows from Lemma 4.2.

Lemma 4.5 (Almost sure asymptotics for σ). Denote by σ_n^1 the largest value among n i.i.d. copies of $\sigma(0)$. Then, under Assumption 1.6, for any c > 1, as $n \to \infty$,

$$g_{\sigma}(\sigma_n^1) \le c \log n$$

eventually almost surely.

Proof. By [9, Theorem 3.5.1] we have equivalence of the statements

$$\left\{ \mathbb{P}\left(g_{\sigma}(\sigma_{n}^{1}) \leq c \log n \text{ ev.}\right) = 1 \right\} \quad \text{and} \quad \left\{ \sum_{n=1}^{\infty} \mathbb{P}\left(g_{\sigma}(\sigma(0)) > c \log n\right) < \infty \right\}.$$

The proof is complete by noticing that, since g_{σ} is continuous by Assumption 1.6, the random variable $\bar{F}_{\sigma}(\sigma(0))$ is uniformly distributed over (0, 1). Hence, for any c > 1,

$$\sum_{n=1}^{\infty} \mathbb{P}(g_{\sigma}(\sigma(0)) > c \log n) = \sum_{n=1}^{\infty} n^{-c} < \infty.$$

4.2. Existence of quick paths. In this section we prove the existence of paths $p \in \Gamma(0, z)$ for certain $z \in V_t$ that have the property that (i) all $\sigma(p_i)$ are relatively small, and (ii) p is not much longer than a direct path to z; what we mean by 'relatively small' and 'not much longer' will depend on the dimension. We shall informally refer to such paths as *quick paths*. The reason we are interested in quick paths is that they are intimately related to the probability that a particle undertaking the BAM reaches a certain site z by time t.

In dimension higher than one, we will additionally require that such paths do not travel too close to a certain well-separated set S_t . The reason for this additional requirement is that we will eventually seek to apply our results to the site Z_t , which depends in a complicated way on $\sigma(z)$ for $z \in B(\Pi^{(L_t)}, \rho)$. We will wish to avoid this dependence, hence our insistence on the fact that the paths do not travel too close to S_t .

4.2.1. Dimension one. In dimension one, there is only one shortest path from 0 to z and this must pass through all intermediate sites. Hence, we seek to show that not too many traps on this path are too large. Clearly, the ability to do this depends on the tail decay of σ , which is the origin of the extra tail decay condition for d = 1 in Assumption 1.6.

To proceed, we must undertake a rather delicate analysis of the trapping landscape σ in the region between 0 and z. We simplify this using coarse graining, essentially placing each site y into a certain 'bin' depending on the value of $\sigma(y)$. We then seek to bound the number of sites in each bin, weighted by the depth of the traps corresponding to each bin. To assist in the coarse graining, we state and prove a technical lemma on the regularity of the upper-tail of $\sigma(0)$.

Lemma 4.6 (Regularity of the upper-tail of $\sigma(0)$). Under Assumption 1.6, let x_t be such that

$$g_{\sigma}(\exp\{\exp\{x_t\}\}) = t\,,$$

which is well-defined by the continuity of g_{σ} . Then, for constants c_1 and c_2 such that $c_2 > c_1 \ge 1$, as $t \to \infty$,

$$g_{\sigma}(\exp\{\exp\{c_2x_t\}\}) > c_1t$$

eventually.

Proof. Let c be the constant in part (d) of Assumption 1.6. In the case where $c < \infty$, for any $\varepsilon > 0$, as $t \to \infty$,

 $t = g_{\sigma}(\exp\{\exp\{x_t\}\}) < x_t (c + \varepsilon)$

eventually. Choosing the $0 < \varepsilon < c(c_2 - c_1)/(c_1 + c_2)$, we have that, as $t \to \infty$,

$$g_{\sigma}(\exp\{\exp\{c_2 x_t\}\}) > c_2 x_t (c-\varepsilon) > t \frac{c_2(c-\varepsilon)}{c+\varepsilon} > c_1 t$$

eventually. On the other hand, in the case where $c = \infty$, then by Assumption 1.6,

$$t = g_{\sigma}(\exp\{\exp\{x_t\}\}) = x_t \kappa_{x_t}$$

for some $\kappa_t \uparrow \infty$. Similarly

$$g_{\sigma}(\exp\{\exp\{c_2 x_t\}\}) = c_2 x_t \kappa_{c_2 x_t} > c_1 x_t \kappa_{x_t} = c_1 t$$

eventually, which completes the proof.

We now define the coarse graining scales that we will use. Let n_t and σ_t be arbitrary functions tending to ∞ as $t \to \infty$.

Lemma 4.7 (Existence of well-spaced coarse graining scales). Let $\varepsilon < 1$ be a constant that satisfies part (c) of Assumption 1.6. Then there exist constants $0 < \delta_1 < \delta_2 < \varepsilon < 1 < c_1$, an integer $I_t = O(\log \log n_t)$ and a set of scaling functions $\{\sigma_t^i\}_{0 \le i \le I_t}$ such that, as $t \to \infty$, the following are all satisfied eventually:

 $\begin{array}{l} (a) \ \ \sigma_t^0 = 0 \,, \, \frac{\log \log \sigma_t^1}{\log \log \sigma_t} \in [1 + \delta_1, 1 + \delta_2] \,, \, \frac{\log \log \sigma_t^i}{\log \log \sigma_t^{i-1}} \in [1 + \delta_1, 1 + \delta_2] \ for \ 2 \le i \le I_t \ ; \\ (b) \ \ g_{\sigma}(\sigma_t^{I_t-1}) \le c_1^{-1} \log n_t \ ; \ and \\ (c) \ \ g_{\sigma}(\sigma_t^{I_t}) \ge c_1 \log n_t \ . \end{array}$

Proof. Choose c_1 , δ_1 and δ_2 such that $1 < c_1^2 < 1 + \delta_2$ and $1 + \delta_1 < (1 + \delta_2)/c_1^2$. Suppose that we define a sequence $\{\bar{\sigma}_t^i\}_{i\geq 0}$ such that

$$\bar{\sigma}_t^0 = 0 \ , \ \frac{\log \log \bar{\sigma}_t^1}{\log \log \sigma_t} = 1 + \delta_1 \quad \text{and} \quad \frac{\log \log \bar{\sigma}_t^i}{\log \log \bar{\sigma}_t^{i-1}} = 1 + \delta_1 \quad \text{for each } i \ge 2 \,,$$

and let I_t be the maximum integer such that

$$g_{\sigma}(\bar{\sigma}_t^{I_t-1}) \le c_1^{-1} \log n_t \; .$$

This satisfies

$$I_t = O(\log \log n_t) \,,$$

since if $I_t > 1$, then eventually

$$(1+\delta_1)^{I_t-2}\log\log\bar{\sigma}_t^1 = \log\log\bar{\sigma}_t^{I_t-1} < g_{\sigma}(\bar{\sigma}_t^{I_t-1}) \le c_1^{-1}\log n_t.$$

Now set $\sigma_t^i = \bar{\sigma}_t^i$ for all $0 \le i \le I_t - 1$, and choose $\sigma_t^{I_t}$ by

$$\begin{cases} \log \log \sigma_t^{I_t} = (1+\delta_2) \log \log \sigma_t^{I_t-1}, & I_t > 1; \\ \log \log \sigma_t^{I_t} = (1+\delta_2) \log \log \sigma_t, & I_t = 1. \end{cases}$$

It remains to check that $g_{\sigma}(\sigma_t^{I_t}) \ge c_1 \log n_t$. By definition,

$$\log \log \sigma_t^{I_t} = \frac{1 + \delta_2}{1 + \delta_1} \log \log \bar{\sigma}_t^{I_t}$$

Then by Lemma 4.6, and the fact that $1 + \delta_1 < (1 + \delta_2)/c_1^2$, as $t \to \infty$,

$$g_{\sigma}(\sigma_t^{I_t}) > c_1^2 g_{\sigma}(\bar{\sigma}_t^{I_t})$$

eventually. Finally, by the definition of I_t ,

$$g_{\sigma}(\bar{\sigma}_t^{I_t}) > c_1^{-1} \log n_t$$

which completes the proof.

Finally, we prove the existence of a quick path. Let c_1 , δ_1 , δ_2 , I_t and $\{\sigma_t^i\}_{0 \le i \le I_t}$ satisfy the conditions in Lemma 4.7. Moreover, for a path $p \in \Gamma_k$ define

$$N_i = \sum_{0 \le j < k} \mathbb{1}_{\{\sigma(p_j) \in (\sigma_t^{i-1}, \sigma_t^i]\}} \quad \text{for each } 1 \le i \le I_t \,.$$

The following proposition essentially bounds the number of sites in each coarse graining scale, weighted by the log of the scale. This will turn out to be the correct definition of a 'quick path' in Section 6.

Proposition 4.8 (Existence of quick paths; d = 1). As $t \to \infty$, each path $p \in \Gamma_{|z|}(0, z)$ with $|z| < n_t$, satisfies

$$\mathbb{P}\left(\sum_{i=1}^{I_t} N_i \log \sigma_t^i < n_t \max\left\{ (\log \sigma_t)^2, \log \log n_t / \log \log \sigma_t \right\} \right) \to 1$$

and

$$\max_{0 \le i < |z|} \sigma(p_i) < \sigma_t^{I_t} \,,$$

eventually almost surely.

Proof. We first prove that the event

$$\mathcal{N}_t := \bigcup_{i=1}^{I_t} \left\{ N_i \le 2n_t \, \bar{F}_\sigma(\sigma_t^{i-1}) \right\}$$

satisfies $\mathbb{P}(\mathcal{N}_t) \to 1$ as $t \to \infty$. Note that each N_i is stochastically dominated by

$$\bar{N}_i \stackrel{d}{=} \operatorname{Binom}(n_t, \bar{F}_{\sigma}(\sigma_t^{i-1}))$$

with $\mathbb{E}\bar{N}_i = n_t \bar{F}_{\sigma}(\sigma_t^{i-1})$ and $\operatorname{Var}\bar{N}_i \leq n_t \bar{F}_{\sigma}(\sigma_t^{i-1})$. By the union bound and Chebyshev's inequality,

$$\mathbb{P}\Big(\bigcup_{i} \{\bar{N}_{i} > 2\mathbb{E}\bar{N}_{i}\}\Big) \leq \sum_{i} \mathbb{P}(\bar{N}_{i} > 2\mathbb{E}\bar{N}_{i}) \leq \sum_{i} \frac{\operatorname{Var}N_{i}}{(\mathbb{E}\bar{N}_{i})^{2}} \leq \sum_{i} \left(n_{t}\bar{F}_{\sigma}(\sigma_{t}^{i-1})\right)^{-1}.$$
 (29)

Since the σ_t^i are increasing in *i*, for any $1 \leq i \leq I_t$,

$$\bar{F}_{\sigma}(\sigma_t^{i-1}) \ge \bar{F}_{\sigma}(\sigma_t^{I_t-1}) \ge n_t^{-c_1^{-1}},$$

by condition (b) of Lemma 4.7. Combining with (29), by the union bound, as $t \to \infty$, eventually

$$\mathbb{P}(\mathcal{N}_t) > 1 - I_t \, n_t^{c_1^{-1} - 1} \to 1$$

since $c_1 > 1$ and $I_t = O(\log \log n_t)$.

So assume the event \mathcal{N}_t holds and split the sum

$$\sum_{i=1}^{I_t} N_i \log \sigma_t^i = N_1 \log \sigma_t^1 + \sum_{i=2}^{I_t} N_i \log \sigma_t^i.$$

For the first term, on the event \mathcal{N}_t and by condition (a) in Lemma 4.7 we have

$$N_1 \log \sigma_t^1 \le 2n_t \bar{F}_{\sigma}(\sigma_t^0) \log \sigma_t^1 = 2n_t \log \sigma_t^1 \le 2n_t (\log \sigma_t)^{1+\delta_2} < n_t (\log \sigma_t)^2/2$$

eventually. Hence it suffices to show that each of the other terms, for $2 \le i \le I_t$, satisfy

$$I_t N_i \log \sigma_t^i < \frac{1}{2} n_t \log \log n_t / \log \log \sigma_t$$

eventually. Recall that by condition (a) in Lemma 4.7, $\log \sigma_t^i \leq (\log \sigma_t^{i-1})^{1+\delta_2}$ for $2 \leq i \leq I_t$. Then, on the event \mathcal{N}_t and by part (c) of Assumption 1.6, eventually,

$$N_i \log \sigma_t^i \le 2n_t \bar{F}_{\sigma}(\sigma_t^{i-1}) \log \sigma_t^i \le 2n_t (\log \sigma_t^{i-1})^{-\varepsilon+\delta_2}$$
$$\le n_t (\log \sigma_t^{i-1})^{-c_2},$$

for some $c_2 > 0$, since $\delta_2 < \varepsilon$. So by monotonicity in *i* and condition (a) in Lemma 4.7,

$$I_t N_i \log \sigma_t^i \le I_t n_t (\log \sigma_t^1)^{-c_2} < n_t \log \log n_t (\log \sigma_t)^{-c_3}$$

eventually, for any $0 < c_3 < c_2(1 + \delta_1)$ which proves the claim.

Finally, the fact that, eventually almost surely,

$$\max_{0 \le i < |z|} \sigma(p_i) < \sigma_t^{I_t}$$

follows from combining condition (c) in Lemma 4.7 with Lemma 4.5.

4.2.2. Dimension higher than one. In dimensions higher than one we use percolation-type estimates to prove the existence of a path $p \in \Gamma(0, z)$ with $z \in S_t$ for some well-separated set S_t that (i) avoids all the deep traps, (ii) has |p| not much more than |z|, and (iii) does not travel too close to sites in S_t . Because we use percolation-type arguments, it will turn out that we need no extra assumption on the tail decay of $\sigma(0)$.

So let us start with the relevant percolation-type estimates; for background on percolation theory see [15]. Consider site percolation on \mathbb{Z}^d with $\mathbb{P}(v \text{ open}) = q$ independently for every $v \in \mathbb{Z}^d$. We say that a subset of \mathbb{Z}^d is *-connected if it is connected with respect to the adjacency relation

$$v \stackrel{*}{\sim} w \Leftrightarrow \max_{1 \le i \le d} |v_i - w_i| = 1,$$

where v_i and w_i denote the coordinate projections of v and w respectively. If $v \stackrel{\sim}{\sim} w$ we say that w is a *-neighbour of v. A *-connected subset of \mathbb{Z}^d is referred to as a *-cluster. The relevance of *-clusters is that they represent the blocking clusters for open paths in \mathbb{Z}^d . For $v \in \mathbb{Z}^d$ a closed site, denote by $\mathcal{C}(v)$ the largest *-cluster of closed sites containing v.

For two sites u, v in \mathbb{Z}^d denote by $d_{\infty}(u, v)$ their chemical distance (also known as the graph distance) with respect to site percolation, defined to be the length of the shortest open path from u to v (and defined to be infinite if no such path exists).

Lemma 4.9 (Expected size and maximum of closed *-clusters). Let $q \in (1 - (3d)^{-1}, 1)$ and suppose u_1, \ldots, u_M are $M \in \mathbb{N}$ distinct closed sites in \mathbb{Z}^d . Then

(i) $\mathbb{E}[|\mathcal{C}(u_1)|] \leq (1-3^d(1-q))^{-1}$, and so in particular $\mathbb{E}[|\mathcal{C}(u_1)|] \to 1$ as $q \to 1$; and (ii) For every $x \in \mathbb{N}$,

$$\mathbb{P}(\max\{|\mathcal{C}(u_1)|, \dots, |\mathcal{C}(u_M)|\} < x) \ge 1 - M(3^d(1-q))^{\lceil \log_{3^d} x \rceil}$$

Proof. Consider performing a breadth-first search on $\mathcal{C}(u_1)$ starting from the site u_1 , by first discovering the closed *-neighbours v_1, \ldots, v_k of u_1 , and then in turn discovering the closed *-neighbours of each of the v_j , $1 \leq j \leq k$, iterating this procedure to explore $\mathcal{C}(u_1)$. Suppose that the site w has just been explored in this procedure. Then the number of closed *-neighbours of w that have not already been discovered is stochastically dominated by a Binom $(3^d - 1, 1 - q)$ random variable. It follows that $|\mathcal{C}(u_1)|$ is stochastically dominated by the total progeny of a branching process with offspring distribution Binom $(3^d, 1 - q)$. Since the expected total progeny of this branching process is $(1 - 3^d(1 - q))^{-1}$, this proves the first statement.

For the second statement, note that by the union bound we have

$$\mathbb{P}(\max\{|\mathcal{C}(u_1)|,\ldots,|\mathcal{C}(u_M)|\} \ge x) \le \sum_{i=1}^M \mathbb{P}(|\mathcal{C}(u_i)| \ge x) = M \mathbb{P}(|\mathcal{C}(u_1)| \ge x).$$

Again by exploring $\mathcal{C}(u_1)$ we have

$$\mathbb{P}(|\mathcal{C}(u_1)| \ge x) \le \mathbb{P}(Z \ge x),$$

where Z is the total progeny of a branching process with offspring distribution $Binom(3^d, 1-q)$. To complete the proof, note that by Markov's inequality we have

$$\mathbb{P}(Z \ge x) \le \mathbb{P}(Z(\lfloor \log_{3^d} x \rfloor) > 0) \le (3^d(1-q))^{\lfloor \log_{3^d} x \rfloor}$$

where Z(n) denotes number of individuals in generation n of the branching process.

Lemma 4.10 (Chemical distance). Fix two sites u, v in \mathbb{Z}^d and a function c := c(q) with $c \to \infty$ as $q \to 1$. Then, as $q \to 1$,

$$\mathbb{P}\left(\frac{d_{\infty}(u,v)}{|u-v|} < 1 + c(1-q)\right) \to 1.$$

Proof. Denote by \mathcal{C}_{∞} the unique infinite open cluster, which exists almost surely for all q sufficiently close to 1 (see [15]). Let $\hat{p} \in \Gamma_{|u-v|}(u, v)$ be any shortest path, denote by K the subset of

 $\operatorname{Set}(\hat{p})$ consisting only of closed sites, and define

$$S := \left| \bigcup_{x \in K} \mathcal{C}(x) \right| \le \sum_{x \in K} |\mathcal{C}(x)|.$$
(30)

By part (i) of Lemma 4.9 and the FKG inequality (see [15], Section 2.2), we have the bound

$$\mathbb{E}[S|\{u, v \in \mathcal{C}_{\infty}\}] \le \frac{\mathbb{E}[|K||\{u, v \in \mathcal{C}_{\infty}\}]}{1 - 3^{d}(1 - q)} \le \frac{|u - v|(1 - q)}{1 - 3^{d}(1 - q)}$$

We now claim that, on the event $\{u, v \in C_{\infty}\}$, it is possible to find a path $p \in \Gamma_k(u, v)$ for some $k \leq |u-v| + (3^d - 1)S$ such that every site in $\operatorname{Set}(p)$ is open. To obtain the required path p take the direct path \hat{p} and divert it around $\mathcal{C}(u)$ for each closed $u \in \operatorname{Set}(\hat{p})$, so that every site in $\operatorname{Set}(p)$ is either in $\operatorname{Set}(\hat{p})$ or in the outer boundary of some $\mathcal{C}(u)$, where by outer boundary we mean the set of sites $\{v \notin \mathcal{C}(u) : \exists u \in \mathcal{C}(u), u \stackrel{*}{\sim} v\}$. This procedure is possible since $u, v \in \mathcal{C}_{\infty}$. Then $\operatorname{Set}(p)$ will consist of just open sites since the outer boundary of each $\mathcal{C}(u)$ is a path of open sites. The bound on |p| follows from the fact that the size of the outer boundary of a *-cluster A is at most $(3^d - 1)|A|$.

We complete the proof of the Lemma with Markov's inequality:

$$\mathbb{P}\left(\frac{d_{\infty}(u,v)}{|u-v|} \ge 1 + c(1-q)\right) \le \mathbb{P}\left(|S| > \frac{c(1-q)|u-v|}{(3^d-1)} \Big| \{u,v\in\mathcal{C}_{\infty}\}\right) + \mathbb{P}\left(\{u,v\in\mathcal{C}_{\infty}\}^c\right)$$
$$\le \frac{3^d}{c(1-3^d(1-q))} + \mathbb{P}\left(\{u,v\in\mathcal{C}_{\infty}\}^c\right).$$

Since $\mathbb{P}(u, v \in \mathcal{C}_{\infty}) \to 1$ as as $q \to 1$, this completes the proof.

We are now ready to show the existence of a quick path in dimensions higher than one. Let $S_t \subseteq \mathbb{Z}^d$ be such that

$$\operatorname{sep}(S_t) > t^{\varepsilon}$$
 and $\min_{u \in S_t} |u| > t^{\varepsilon}$

eventually for some $\varepsilon > 0$. Recall the definition of $j := [\gamma - 1]$. Let σ_t be an arbitrary function tending to infinity as $t \to \infty$. Define the set

$$\mathbb{Z}^d(\sigma_t, S_t) := \{ z \in \mathbb{Z}^d : \sigma(z) \le \sigma_t, z \notin B(S_t, j) \}.$$

For a site $z \in \mathbb{Z}^d$, let $|z|_{\text{chem}}$ be the chemical distance of the ball B(z, j) in this set, that is, the length of the shortest path from the origin to $\partial B(z, j)$ that lies exclusively in this subgraph (setting it as ∞ if such a path does not exist).

Proposition 4.11 (Existence of quick paths; d > 1). Let $z_t \in S_t \cap V_t$ and let c_t be a function such that $c_t \to \infty$ as $t \to \infty$ and $\overline{F}_{\sigma}(\sigma_t)c_t \ll 1$. Then, there exists a constant c > 0 such that, as $t \to \infty$,

$$\mathbb{P}\left(\frac{|z_t|_{\text{chem}}}{|z_t|} \le 1 + \bar{F}_{\sigma}(\sigma_t)c_t + t^{-c}\right) \to 1.$$

Proof. Let $q := 1 - \bar{F}_{\sigma}(\sigma_t)$. By Lemma 4.10, with probability tending to 1 as $t \to \infty$ there exists a path $\hat{p} \in \Gamma_{\ell_t}(0, z_t)$ for some

 $\ell_t \le |z_t| (1 + \bar{F}_{\sigma}(\sigma_t) c_t)$

such that $\sigma(\hat{p}_i) \leq \sigma_t$ for all $0 \leq i < \ell_t$. Let $i = \min\{0 \leq j < \ell_t : \hat{p}_j \in \partial B(z_t, j)\}$ and define $v_t := \hat{p}_i$ to be the first site in $\partial B(z_t, j)$ visited by path \hat{p} . We show how to modify \hat{p} so that we obtain a new path $p \in \Gamma(0, v_t)$ for some $v_t \in \partial B(Z_t, j)$ with $\operatorname{Set}(p) \subseteq \mathbb{Z}^d(\sigma_t, S_t)$.

The modification is done by diverting \hat{p} around the balls of radius j centred on sites in S_t . In doing so, we may encounter new closed sites v, and these too must be avoided if we wish to find a path p with $\operatorname{Set}(p) \subseteq \mathbb{Z}^d(\sigma_t, S_t)$. Formally, the set of these new closed sites is precisely

$$\{x \in \partial B \left(S_t \cap B(\operatorname{Set}(\hat{p}), j), j \right) : \sigma(x) > \sigma_t \}.$$

Denote by M_t the size of this set and its elements as w_1, \ldots, w_{M_t} , and choose $0 < c_1 < \varepsilon$ where ε is the constant appearing in the definition of S_t . Then by the separation of sites in S_t , we have

$$|S_t \cap B(\operatorname{Set}(\hat{p}), j)| \le 2\ell_t t^{-\varepsilon},$$

and so

$$M_t \le 3^d |B(0,j)| \ell_t t^{-\varepsilon} < |z_t| t^{-c_1}$$
(31)

for all t sufficiently large. Choose now $0 < c_2 < c_1$, $\alpha < -1 - (1 - c_1)/c_2$, and t sufficiently large so that

$$\bar{F}_{\sigma}(\sigma_t) < 3^{d\alpha}$$
 .

Applying part (ii) of Lemma 4.9, we deduce that

$$\max\{|\mathcal{C}(w_1)|,\ldots,|\mathcal{C}(w_{M_t})|\} \le t^{c_2}$$

with probability tending to 1 as $t \to \infty$. We claim this implies that, by the separation of sites in S_t and the fact that $c_2 < \varepsilon$, with overwhelming probability there exists a path $p \in \Gamma(0, v_t)$ which avoids all *j*-balls centred on sites in S_t and all closed sites. Indeed to obtain this path we take path \hat{p} and then divert around *j*-balls centred on sites in S_t and then further divert around any new closed *-clusters we encounter. Since we know that no such cluster is too large, they cannot cut the origin off from v_t in $\mathbb{Z}^d(\sigma_t, S_t)$, and furthermore we will not encounter any more sites in S_t on the new path.

We can now bound |p|. The number of additional sites we must visit to obtain p from \hat{p} is at most $3^d M_t(|B(0,j)| + t^{c_2})$ with probability tending to 1 as $t \to \infty$; this comes from counting the diversions around each j-ball and the diversions around each closed cluster we then encounter. Using (31), we can thus choose $0 < c < c_1 - c_2$ to yield the result.

5. Extremal theory for local eigenvalues

In this section, we use point process techniques to study the random variables $Z_t^{(j)}$ and $\Psi_t^{(j)}(Z_t^{(j)})$, and generalisations thereof; the techniques used are similar to those found in [2, 10, 26], although we strengthen the results available in those papers. In the process, we complete the proof of Theorems 1.9 and 1.10. Throughout this section, let ε be such that $0 < \varepsilon < \theta$.

5.1. Upper-tail properties of the local principal eigenvalues. The first step is to give uppertail asymptotics for the distribution of the local principal eigenvalues $\lambda^{(n)}(z)$ for $z \in \Pi^{(L_t)}$ and $n \in \mathbb{N}$. These will allow us to study the random variables $Z_t^{(j)}$ and $\Psi_t^{(j)}(Z_t^{(j)})$ via point process techniques. For technical reasons, we shall actually consider an i.i.d. set of *punctured* versions of $\lambda^{(n)}(0)$ which can be coupled to coincide with $\lambda^{(n)}(z)$ eventually almost surely for each $z \in \Pi^{(L_t)}$.

To this end, let $\{\tilde{\xi}_z\}_{z \in V_t}$ be an independent collection of i.i.d. potential fields $\tilde{\xi}_z : \mathbb{Z}^d \to \mathbb{R}$ distributed as

$$\tilde{\xi}_{z}(y) \stackrel{d}{=} \begin{cases} \xi(0) , & \text{if } y = z ;\\ \xi(0) \mathbb{1}_{\{\xi(0) < L_{t}\}} , & \text{otherwise} , \end{cases}$$
(32)

and define

$$\tilde{\Pi}^{(L_t)} := \{ z \in V_t : \tilde{\xi}_z(z) > L_t \} \quad \text{and} \quad \tilde{\Pi}^{(L_{t,\varepsilon})} := \{ z \in V_t : \tilde{\xi}_z(z) > L_{t,\varepsilon} \}$$

by analogy with $\Pi^{(L_t)}$ and $\Pi^{(L_{t,\varepsilon})}$. Similarly, let $\{\tilde{\sigma}_z\}_{z \in V_t}$ be an independent collection of i.i.d. trapping landscapes $\tilde{\sigma}_z : \mathbb{Z}^d \to \mathbb{R}$ distributed as $\tilde{\sigma}_z(y) \stackrel{d}{=} \sigma(0)$, and abbreviate $\tilde{\eta}_z(z) = \tilde{\xi}_z(z) - \tilde{\sigma}_z^{-1}(z)$. Then, for each $z \in V_t$ and $n \in \mathbb{N}$, let $\tilde{\lambda}^{(n)}(z)$ be the principal eigenvalue of the *punctured* Hamiltonian

$$\tilde{\mathcal{H}}^{(n)}(z) := \Delta \tilde{\sigma}_z^{-1} + \tilde{\xi}_z$$

restricted to the domain B(z, n) with Dirichlet boundary conditions, and observe that $\{\tilde{\lambda}^{(n)}(z)\}_{z \in V_t}$ are i.i.d. by construction. Note that we have suppressed the explicit *t*-dependence in $\tilde{\xi}_z$ (and hence $\tilde{\lambda}^{(n)}(z)$ etc.).

We now proceed to study the upper-tail asymptotics for the distribution of each $\tilde{\lambda}^{(n)}(z)$.

Proposition 5.1 (Path expansion for $\tilde{\lambda}^{(n)}$). For each $n \in \mathbb{N}$ and $z \in \tilde{\Pi}^{(L_{t,\varepsilon})}$ uniformly, as $t \to \infty$,

$$\tilde{\lambda}^{(n)}(z) = \tilde{\eta}_{z}(z) + \tilde{\sigma}_{z}^{-1}(z) \sum_{\substack{2 \le k \le 2j \\ p_{i} \ne z, \ 0 < i < k \\ \operatorname{Set}(p) \subseteq B(z,n)}} \sum_{\substack{0 < i < k \\ \operatorname{Set}(p) \subseteq B(z,n)}} \prod_{\substack{0 < i < k \\ 0 < i < k}} (2d)^{-1} \frac{\tilde{\sigma}_{z}^{-1}(p_{i})}{\tilde{\lambda}^{(n)}(z) - \tilde{\eta}_{z}(p_{i})} + o(d_{t}e_{t}),$$

 $= \tilde{\eta}_z(z) + O(a_t^{-1}) \,.$

Proof. Applying Proposition 3.7 we have that

$$\tilde{\lambda}^{(n)}(z) = \tilde{\eta}_{z}(z) + \tilde{\sigma}_{z}^{-1}(z) \sum_{k \ge 2} \sum_{\substack{p \in \Gamma_{k}(z,z) \\ p_{i} \ne z, \ 0 < i < k \\ \operatorname{Set}(p) \subseteq B(z,n)}} \prod_{0 < i < k} (2d)^{-1} \frac{\tilde{\sigma}_{z}^{-1}(p_{i})}{\tilde{\lambda}^{(n)}(z) - \tilde{\eta}_{z}(p_{i})} + \frac{\tilde{\sigma}_{z}^{-1}(p_{i})}{\tilde{\lambda}^{(n)}(z) - \tilde{\tau}_{z}(p_{i})} + \frac{\tilde{\sigma}_{z}^{-1}(p_{i})}{\tilde{\lambda}^{(n)}(z) - \tilde{\tau}_{z}(p_{i})} + \frac{\tilde{\sigma}_{z}^{-1}(p_{i})}{\tilde{\lambda}^{(n)}(z) - \tilde{\tau}_{z}(p_{i})} + \frac{\tilde{\sigma}_{z}^{-1}(p_{i})}{\tilde{\lambda}^{(n)}(z) - \tilde{\tau}_{z}(p_{i})} + \frac{\tilde{\sigma}_$$

Now observe that, by Lemma 3.2 and the truncation in $\tilde{\xi}_z$, for each $p_i \in B(z,n) \setminus \{z\}$,

$$\tilde{\lambda}^{(n)}(z) - \tilde{\eta}_z(p_i) > L_{t,\varepsilon} - L_t - \delta_{\sigma}^{-1} \sim (\theta - \varepsilon)a_t.$$

Moreover, each $\tilde{\sigma}_z^{-1}(p_i)$ is bounded above by δ_{σ}^{-1} . Finally, as $t \to \infty$,

$$a_t^{-(2j+2)} = o(d_t e_t),$$

by the definition of j. This means that, up to the error $o(d_t e_t)$, we can truncate the sum at paths with 2j steps. It also means that the total contribution from the sum over paths $p \in \Gamma_k(z, z)$ is $O(a_t^{-1})$.

Proposition 5.2 (Extremal theory for $\tilde{\lambda}^{(n)}$; see [2, Section 6], [10, Proposition 4.2]). For each $n \in \mathbb{N}$, there exists a scaling function $A_t = a_t + O(1)$ such that, as $t \to \infty$ and for each fixed $x \in \mathbb{R}$,

$$t^d \mathbb{P}\left(\tilde{\lambda}^{(n)}(0) > A_t + xd_t\right) \to e^{-x}$$

Moreover, there exists a c > 0 such that, as $t \to \infty$ and uniformly for x > 0,

$$t^d \mathbb{P}\left(\tilde{\lambda}^{(n)}(0) > A_t + xd_t\right) < e^{-cx^{\min\{1,\gamma\}}}.$$

Proof. To ease notation, we abbreviate the fields $\tilde{\xi}_0$ and $\tilde{\sigma}_0$ by ξ and σ respectively (although these should not be confused with the original fields ξ and σ , since the law of $\tilde{\xi}_0(z)$ has been truncated at $z \neq 0$). First remark that, by Lemma 3.2, as $t \to \infty$,

$$\tilde{\lambda}^{(n)}(0) > A_t + x d_t$$
 implies that $\xi(0) > L_{t,\varepsilon}$,

eventually, which means that we can apply the path expansion in Proposition 5.1 to $\tilde{\lambda}^{(n)}(0)$. Let A_t be an arbitrary scale such that $A_t = a_t + O(1)$, and define the function

$$Q(A_t;\xi,\sigma) := \sigma^{-1}(0) + \sigma^{-1}(0) \sum_{\substack{2 \le k \le 2j \\ p_i \ne 0, \ 0 \le i < k \\ \operatorname{Set}(p) \subseteq B(z,j)}} \prod_{\substack{0 < i < k \\ 0 \le i < k}} (2d)^{-1} \frac{\sigma^{-1}(p_i)}{A_t - \eta(p_i)},$$

if $\xi(y) < L_t$ for each $y \in B(0, j) \setminus \{0\}$ and $Q(A_t; \xi, \sigma) := 0$ otherwise. Then, since $\tilde{\lambda}^{(n)}(0)$ is strictly increasing in $\xi(0)$ we have that, as $t \to \infty$,

$$\mathbb{P}\left(\tilde{\lambda}^{(n)}(0) > A_t + xd_t\right) \sim \mathbb{P}\left(\xi(0) > A_t + xd_t + Q(A_t + xd_t;\xi,\sigma)\right) \\
\sim \mathbb{P}\left(\xi(0) > A_t + xd_t + Q(A_t;\xi,\sigma)\right)$$
(33)

$$\sim t^{-d} e^{-x} \int_{\xi,\sigma} \exp\left\{a_t^{\gamma} - \left(A_t + Q(A_t;\xi,\sigma)\right)^{\gamma}\right\} d\mu_{\xi} d\mu_{\sigma}$$
(34)

where the first asymptotic accounts for the error in the path expansion Proposition 5.1, the second and third asymptotics result from Taylor expansions, and are uniform in ξ and σ , and

where μ_{ξ} and μ_{σ} stand for the joint probability densities of $\{\xi(y)\}_{y \in B(0,n) \setminus \{0\}}$ and $\{\sigma(y)\}_{y \in B(0,n)}$ respectively. Consider then the integral in (34), which we abbreviate as

$$I(A_t) := \int_{\xi,\sigma} f(A_t;\xi,\sigma) \, d\mu_\xi \, d\mu_\sigma \,, \quad f(A_t;\xi,\sigma) := \exp\left\{a_t^\gamma - \left(A_t + Q(A_t;\xi,\sigma)\right)^\gamma\right\}.$$

Note that $Q(A_t; \xi, \sigma)$ is uniformly bounded as $t \to \infty$ (by $2\delta_{\sigma}^{-1}$ for instance). Hence, for C sufficiently large, as $t \to \infty$ eventually

$$f(a_t + C; \xi, \sigma) < 1 < f(a_t - C; \xi, \sigma)$$

uniformly in $\{\xi(y)\}_{y \in B(0,n) \setminus \{0\}}$ and $\{\sigma(y)\}_{y \in B(0,n)}$. This implies that $I(A_t - C) < 1 < I(a_t + C)$. Moreover, since $f(A_t; \xi, \sigma)$ is continuous in A_t uniformly for each $\{\xi(y)\}_{y \in B(0,n) \setminus \{0\}}$ and $\{\sigma(y)\}_{y \in B(0,n)}$, the function $I(A_t)$ is continuous in A_t . Hence, by the intermediate value theorem function, there exists an $A_t = a_t + O(1)$ such that, as $t \to \infty$ eventually $I(A_t) = 1$, which gives the first result. For the second, instead of (33) we bound $Q(A_t + xd_t; \xi, \sigma)$ above, uniformly in x > 0, by $Q(A_t; \xi, \sigma)$, which produces the bound

$$t^{-d} \int_{\xi,\sigma} \exp\left\{a_t^{\gamma} - \left(A_t + Q(A_t;\xi,\sigma)\right)^{\gamma} \left(1 + \frac{x}{\gamma}(\log t)^{-1}\right)^{\gamma}\right\} d\mu_{\xi} d\mu_{\sigma}.$$

In the case $\gamma \geq 1$, we bound this expression above uniformly in x > 0 by

$$t^{-d} \int_{\xi,\sigma} \exp\left\{a_t^{\gamma} - \left(A_t + Q(A_t;\xi,\sigma)\right)^{\gamma} \left(1 + \frac{x}{\gamma}(\log t)^{-1}\right)\right\} d\mu_{\xi} d\mu_{\sigma} \sim e^{-\frac{x}{\gamma}(1+o(1))},$$

using the definition of A_t and the fact that $A_t + Q(A_t; \xi, \sigma) \sim a_t$ in the last step. The case $\gamma < 1$ is simpler, since then we have simply

$$\mathbb{P}\bigg(\xi(0) > A_t + xd_t + \sigma^{-1}(0)\bigg) = \mathbb{P}\bigg(\xi(0) > a_t + xd_t + O(1)\bigg)$$

and the bound follows from the regularity of Weibull tail of $\xi(0)$ in Assumption 1.5.

We now define the set-up we shall need to examine the correlation of the potential field and trapping landscape near sites of high $\tilde{\lambda}^{(n)}$; since the nature of this correlation differs depending on (γ, μ) , so does our set-up. Fix a constant $\nu \in (0, 1)$. Recalling the definition of the 'interface cases' \mathcal{B} and \mathcal{B}_{ξ} , define the 'interface sites'

$$\mathcal{F} := \begin{cases} y \in \mathbb{Z}^d : |y| = \rho, & \text{if } (\gamma, \mu) \in \mathcal{B}, \\ \varnothing, & \text{else}, \end{cases} \quad \text{and} \quad \mathcal{F}_{\xi} := \begin{cases} y \in \mathbb{Z}^d : |y| = \rho_{\xi}, & \text{if } (\gamma, \mu) \in \mathcal{B}_{\xi}, \\ \varnothing, & \text{else}. \end{cases}$$

Recalling the definition of n(y), for each $y \in \mathbb{Z}^d$ define the positive constants

$$c_{\sigma} := \begin{cases} \left(\frac{\gamma}{\mu}\right)^{\frac{1}{\mu+1}}, & \text{if } q_{\sigma} > 0, \\ 0, & \text{else}, \end{cases}, \quad c_{\xi}(y) := \begin{cases} \left(n(y)^{2}(2d)^{-1}\delta_{\sigma}^{-1}c_{\sigma}^{-1}\right)^{\frac{1}{\gamma-1}}, & \text{if } q_{\xi}(|y|) > 0, \\ 0, & \text{else}, \end{cases}$$

 $\bar{c}_{\sigma}(y) := n(y)^2 (2d)^{-1} \gamma c_{\sigma}^{-1}$ and $\bar{c}_{\xi}(y) := \bar{c}_{\sigma}(y) \delta_{\sigma}^{-1}$. For each $n \in \mathbb{N}$, if $\mu > 0$ and $\gamma > 1$, define the rectangles

$$\begin{split} E_{\xi} &:= \prod_{y \in (B(0, n \wedge \rho_{\xi}) \setminus \{0\}) \setminus \mathcal{F}_{\xi}} (-f_t, f_t) \ \times \prod_{y \in (B(0, n) \setminus B(0, n \wedge \rho_{\xi})) \cup \mathcal{F}_{\xi}} (f_t, g_t) \,, \\ E_{\sigma} &:= (-f_t, f_t) \ \times \prod_{y \in (B(0, n) \setminus \{0\}) \setminus \mathcal{F}} (0, f_t) \ \times \prod_{y \in (B(0, n) \setminus B(0, n \wedge \rho)) \cup \mathcal{F}} (0, g_t) \,, \\ S_{\xi} &:= \prod_{y \in (B(0, n \wedge \rho_{\xi}) \setminus \{0\}) \setminus \mathcal{F}_{\xi}} a_t^{q_{\xi}(|y|)} (c_{\xi}(y) - f_t, c_{\xi}(y) + f_t) \ \times \prod_{y \in (B(0, n) \setminus B(0, n \wedge \rho_{\xi})) \cup \mathcal{F}_{\xi}} (f_t, g_t) \,, \end{split}$$

and

$$S_{\sigma} := a_t^{q_{\sigma}}(c_{\sigma} - f_t, c_{\sigma} + f_t) \times \prod_{y \in (B(0,n) \setminus \{0\}) \setminus \mathcal{F}} (\delta_{\sigma}, \delta_{\sigma} + f_t) \times \prod_{y \in (B(0,n) \setminus B(0,n \wedge \rho)) \cup \mathcal{F}} (0, g_t) \,.$$

If $\mu = 0$ and $\gamma > 1$, define instead

$$E_{\sigma} := (a_t^{-\nu}, \infty) \times \prod_{y \in B(0,n) \setminus \{0\}} (0, g_t) \text{ and } S_{\sigma} := a_t^{\gamma - 1}(a_t^{-\nu}, \infty) \times \prod_{y \in B(0,n) \setminus \{0\}} (0, g_t) + \sum_{y \in B(0,n) \setminus \{0\}} (0,$$

whereas if $\gamma \leq 1$, maintain the definition of E_{σ} but define instead

$$S_{\sigma} := \prod_{y \in B(0,n)} (0,g_t)$$

For each $n \in \mathbb{N}$ and $z \in \mathbb{Z}^d$, define the event

$$\mathcal{S}_t^{(n)}(z) := \left\{ \{ \tilde{\xi}_z(y) \}_{y \in B(z,n) \setminus \{z\}} \in S_{\xi} , \ \{ \tilde{\sigma}_z(y) \}_{y \in B(z,n)} \in S_{\sigma} \right\} ,$$

and, for each $x \in \mathbb{R}$ and the scaling function A_t from Proposition 5.2, further define the event

$$\mathcal{A}_t := \left\{ \tilde{\lambda}^{(n)}(0) > A_t + x d_t \right\} \,.$$

Proposition 5.3 (Correlation of potential field and trapping landscape). For each $n \in \mathbb{N}$, as $t \to \infty$,

$$\mathbb{P}\left(\mathcal{S}_t^{(n)}(0)\big|\mathcal{A}_t\right) \to 1\,.$$

Moreover, as $t \to \infty$,

$$f_{\tilde{\xi}_0(y)|\mathcal{A}_t}(x) \to \frac{e^{\bar{c}_{\xi}(y)x}f_{\xi}(x)}{\mathbb{E}[e^{\bar{c}_{\xi}(y)\xi(0)}]} , \quad \text{for each } y \in \mathcal{F}_{\xi},$$
(35)

uniformly over $x \in (0, L_t)$, and

$$f_{\tilde{\sigma}_0(y)|\mathcal{A}_t}(x) \to \frac{e^{\bar{c}_\sigma(y)/x} f_\sigma(x)}{\mathbb{E}[e^{\bar{c}_\sigma(y)/\sigma(0)}]}, \quad \text{for each } y \in \mathcal{F},$$
(36)

uniformly over x. Finally, if $\gamma = 1$, then for each $x \in \mathbb{R}^+$, as $t \to \infty$,

$$f_{\tilde{\sigma}_0(0)|\mathcal{A}_t}(x) \to \frac{e^{-1/x} f_\sigma(x)}{\mathbb{E}[e^{-1/\sigma(0)}]},$$
(37)

 $uniformly \ over \ x.$

Proof. Once again, we abbreviate $\tilde{\xi}_0$ and $\tilde{\sigma}_0$ as ξ and σ respectively. Define a field $s : B(0,n) \setminus \{0\} \cup B(0,n) \to \mathbb{R}$ with projections s_{ξ} and s_{σ} onto $B(0,n) \setminus \{0\}$ and B(0,n) respectively. For a scale $C_t \sim a_t$ define the function

$$Q_t(C_t;s) := a_t^{-q_{\sigma}} (c_{\sigma} + s_{\sigma}(0))^{-1} - a_t^{-q_{\sigma}} (c_{\sigma} + s_{\sigma}(0))^{-1} \\ \times \sum_{\substack{2 \le k \le 2j \\ p_i \ne 0, 0 < i < k \\ \text{Set}(p) \subseteq B(0,n)}} \sum_{\substack{0 < i < k \\ O(n)}} \prod_{\substack{0 < i < k \\ O(n)}} \frac{(2d)^{-1} (\delta_{\sigma} + s_{\sigma}(p_i))^{-1}}{(C_t - a_t^{q_{\xi}(|p_i|)} (c_{\xi}(p_i) + s_{\xi}(p_i)) + (\delta_{\sigma} + s_{\sigma}(p_i))^{-1}},$$

if, for each $y \in B(0, n) \setminus \{0\}$,

$$a_t^{q_{\xi}(|y|)}(c_{\xi}(y) + s_{\xi}(y)) \in (0, L_t)$$
, $s_{\sigma}(y) > 0$ and $a_t^{q_{\sigma}}(c_{\sigma} + s_{\sigma}(0)) > 0$
are satisfied and $Q_t(C_t; s) := 0$ otherwise. Define further the function

are satisfied, and
$$Q_t(C_t, s) := 0$$
 otherwise. Define further the function

$$\begin{split} R_t(C_t;s) &:= a_t^{\gamma} - (C_t + Q_t(C_t;s))^{\gamma} + \sum_{y \in B(0,n)} \left(\log f_{\xi} \left(a_t^{q_{\xi}(|y|)}(c_{\xi}(y) + s_{\xi}(y)) \right) + \log a_t^{q_{\xi}(|y|)} \right) \\ &+ \log f_{\sigma} \left(a_t^{q_{\sigma}}(c_{\sigma} + s_{\sigma}(0)) + \log a_t^{q_{\sigma}} + \sum_{y \in B(0,n) \setminus \{0\}} \log f_{\sigma} \left(\delta_{\sigma} + s_{\sigma}(y) \right) \right). \end{split}$$

To motivate these definitions, consider that, similarly to the above, we can write

$$\mathbb{P}\left(\tilde{\lambda}^{(n)}(0) > A_t + xd_t\right) \sim t^{-d} e^{-x} \int_{\mathbb{R}^{2|B(0,n)|-1}} \exp\left\{R_t(A_t;s)\right\} \, ds \,. \tag{38}$$

It remains to show that the integral in (38) is asymptotically concentrated on the set $E_{\xi} \times E_{\sigma}$ and that equations (35)–(37) are satisfied. This fact can be checked by a somewhat lengthy computation which we only sketch here. We shall treat separately three cases: (i) $\gamma > 1$ and $\mu > 0$; (ii) $\gamma > 1$ and $\mu = 0$; and (iii) $\gamma \le 1$. We begin with case (i), which is the most delicate.

We must analyse the variables $s_{\sigma}(0)$, $\{s_{\sigma}(y)\}_{y \in B(0,n) \setminus \{0\}}$, and $\{s_{\xi}(y)\}_{y \in B(0,n) \setminus \{0\}}$ separately; we start with $s_{\sigma}(0)$. In what follows abbreviate $R_t(A_t; s)$ by $R_t(s)$. Fix an arbitrary choice of the components of s and consider how $R_t(s)$ varies with $s_{\sigma}(0)$. Notice that the function $R_t(s)$ can be decomposed into two parts, one of which decreases as $s_{\sigma}(0)$ increases (through Q_t) and another which increases as $s_{\sigma}(0)$ increases (through f_{σ}). The first part is analysed by Taylor expanding $(A_t + Q_t(A_t; s))^{\gamma}$, from which it can be seen that the dependence on $s_{\sigma}(0)$ is, as $t \to \infty$,

$$\gamma a_t^{-q_\sigma} a_t^{\gamma-1} (c_\sigma + s_\sigma(0))^{-1} (1 + o(1))$$

where the error term o(1) is uniform in s. The second part is given by $-\log f_{\sigma}(a_t^{q_{\sigma}}(c_{\sigma} + s_{\sigma}(0)))$ which is, if $\mu > 0$, eventually

$$a_t^{q_\sigma\mu}(c_\sigma+s_\sigma(0))^\mu.$$

Hence, since we defined q_{σ} precisely so that

$$-q_{\sigma} + \gamma - 1 = q_{\sigma}\mu \,,$$

the function R_t has the asymptotic form, as $t \to \infty$,

$$R_t(s) = f_1(t;s) + a_t^{\kappa_1} \left(g_1(s_\sigma(0) + o(1)) \right)$$

where $f_1(t; s)$ is some function not depending on $s_{\sigma}(0)$, κ_1 is some positive constant, the function $g_1(x)$ satisfies

$$g_1(x) := -\gamma (c_{\sigma} + x)^{-1} - (c_{\sigma} + x)^{\mu}$$

and the error term o(1) is uniform in s. Then we have, uniformly in s, as $t \to \infty$,

$$\int_{\mathbb{R}} e^{R_t(s)} \, ds_{\sigma}(0) \sim e^{f_1(t;s)} \int_{\mathbb{R}} \exp\left\{a_t^{\kappa_1} g_1(s_{\sigma}(0))\right\} \, ds_{\sigma}(0) \,. \tag{39}$$

Remark that $g_1(x)$ achieves a unique maximum at 0 (by the construction of c_{σ}). Therefore, by the Laplace method, the above integral is eventually asymptotically concentrated around 0 on the order $a_t^{\kappa_1}$, and hence the integral is concentrated on the domain $s_{\sigma}(0) \in (-f_t, f_t)$.

Consider now the variables $\{s_{\sigma}(y)\}_{y \in B(0,n) \setminus \{0\}}$. Fix an $s_{\sigma}(0) \in (-f_t, f_t)$ and an arbitrary choice of the remaining components of s. Again, similarly to the above, the function $R_t(s)$ can be decomposed into two parts, one whose dependence on $s_{\sigma}(y)$ is, as $t \to \infty$,

$$n(y)^{2} (2d)^{-1} \gamma c_{\sigma}^{-1} a_{t}^{\gamma - 2|y|} a_{t}^{-q_{\sigma}} (\delta_{\sigma} + s_{\sigma}(y))^{-1} (1 + o(1))$$

uniformly in s, and another whose dependence is

$$-\log f_{\sigma}(\delta_{\sigma}+s_{\sigma}(y)).$$

Then we have, uniformly in s, as $t \to \infty$,

$$\int_{\mathbb{R}} e^{R_t(s)} ds_{\sigma}(y) \sim e^{f_2(t;s)} \int_{\mathbb{R}} \exp\left\{\gamma c_{\sigma}^{-1} a_t^{\kappa_2} (\delta_{\sigma} + s_{\sigma}(y))^{-1}\right\} f_{\sigma}(\delta_{\sigma} + s_{\sigma}(y)) ds_{\sigma}(y) ,$$

where $f_2(t; s)$ is some function not depending on $s_{\xi}(y)$, κ_2 is some non-negative constant with $\kappa_2 > 0$ if and only if $y \in B(0, \rho) \setminus \mathcal{F}$, and where the error term o(1) is uniform in s. Hence, if $y \in B(0, \rho) \setminus \mathcal{F}$, then along with the lower-tail assumption in 1.6, it is clear that the above integral is asymptotically concentrated on $s_{\sigma}(y) \in (0, f_t)$. On the other hand, if $y \in \mathcal{F}$, then the integrand is asymptotically

$$e^{\bar{c}_{\sigma}(y)/(s_{\sigma}(y)+\delta_{\sigma})}f_{\sigma}(s_{\sigma}(y)+\delta_{\sigma}),$$

uniformly over $s_{\sigma}(y)$, which establishes (36). Trivially, if $y \notin B(0, \rho)$, then the integral is concentrated on $s_{\sigma}(y) \in (f_t, g_t)$.

Finally, consider the variables $\{s_{\xi}(y)\}_{y \in B(0,n) \setminus \{0\}}$ and fix $s_{\sigma}(0) \in (-f_t, f_t)$, $s_{\sigma}(y) \in (0, f_t)$ for each $y \in B(0, \rho) \setminus \mathcal{F}$, and an arbitrary choice of the remaining components of s. The function $R_t(s)$ can be decomposed into two parts, one whose dependence on $s_{\xi}(y)$ is of order, as $t \to \infty$,

$$n(y)^{2} (2d)^{-1} (\delta_{\sigma} + s_{\sigma}(y))^{-1} \gamma c_{\sigma}^{-1} a_{t}^{q_{\xi}(|y|)} a_{t}^{\gamma - 1 - 2|y|} a_{t}^{-q_{\sigma}} \left(c_{\xi}(y) + s_{\xi}(y) \right) \left(1 + o(1) \right) \,,$$

uniformly in s, another whose dependence is

$$a_t^{q_{\xi}(|y|)\gamma}(c_{\xi}(y)+s_{\xi}(y))^{\gamma}$$

Hence, since we defined $q_{\xi}(|y|)$ precisely so that

$$q_{\xi}(|y|) + \gamma - 1 - 2|y| - q_{\sigma} = q_{\xi}(|y|)\gamma$$

if $y \in B(0, \rho_{\xi})$, the function R_t has the asymptotic form, as $t \to \infty$,

$$R_t(s) = f_3(t;s) + a_t^{\kappa_3} \left(g_3(s_{\xi}(y)) + o(1) \right)$$

where $f_3(t;s)$ is some function not depending on $s_{\xi}(y)$, κ_3 is some non-negative constant with $\kappa_3 > 0$ if any only if $y \in B(0, \rho_{\xi}) \setminus \mathcal{F}_{\xi}$, the function $g_3(x)$ satisfies

$$g_3(x) := \gamma n(y)^2 (2d)^{-1} \delta_{\sigma}^{-1} c_{\sigma}^{-1} (c_{\xi}(y) + x) - (c_{\xi}(y) + x)^{\gamma},$$

and where the error term o(1) is uniform in s. Then we have, uniformly in s, as $t \to \infty$,

$$\int_{\mathbb{R}} e^{R_t(s)} ds_{\xi}(y) \sim e^{f_3(t;s)} \int_{\mathbb{R}} \exp\left\{a_t^{\kappa_3} g_3(s_{\xi}(y))\right\} ds_{\xi}(y) \,.$$

If $y \in B(0, \rho_{\xi}) \setminus \mathcal{F}_{\xi}$, and since $g_3(x)$ achieves a unique maximum at 0 (by the construction of $c_{\xi}(y)$), again by the Laplace method this integral is also asymptotically concentrated on $s_{\xi}(y) \in (-f_t, f_t)$. On the other hand, if $y \in \mathcal{F}_{\xi}$, then the integrand is asymptotically

$$e^{\bar{c}_{\xi}(y)s_{\xi}(y)}f_{\xi}(s_{\xi}(y)),$$

uniformly over $s_{\xi}(y)$, which establishes (35). Trivially, if $y \notin B(0, \rho_{\xi})$, then the integral is concentrated on $s_{\xi}(y) \in (f_t, g_t)$. Since we have now shown that each component of (38) is asymptotically concentrated on the respective component of the set $E_{\xi} \times E_{\sigma}$, integrating first over $s_{\xi}(y)$ and $s_{\sigma}(y)$ for $y \in B(0, n) \setminus \{0\}$, and then over $s_{\sigma}(0)$, we have the result.

We now turn to case (ii). In this case the integral over $s_{\sigma}(0)$ in (39) becomes

$$e^{f_1(t;s)} \int_{\mathbb{R}} e^{-\gamma s_{\sigma}^{-1}(0)} f_{\sigma} \left(A_t^{\gamma-1} s_{\sigma}(0) \right) \, ds_{\sigma}(0) \sim e^{f_1(t;s)} \int_{\mathbb{R}} e^{-\gamma s_{\sigma}^{-1}(0)} f_{\sigma} \left(a_t^{\gamma-1} s_{\sigma}(0) \right) \, ds_{\sigma}(0) \,,$$

where we used the regularity in 1.6 in the last step. On the region $(0, a_t^{-\nu})$, this integral can be bounded above as

$$\int_{0}^{a_{t}^{-\nu}} e^{-\gamma s_{\sigma}^{-1}(0)} f_{\sigma} \left(a_{t}^{\gamma-1} s_{\sigma}(0) \right) \, ds_{\sigma}(0) \le \int_{0}^{a_{t}^{-\nu}} e^{-\gamma s_{\sigma}^{-1}(0)} \, ds_{\sigma}(0) \le e^{-\gamma a_{t}^{\nu}}$$

On the other hand, for any $0 < c < \nu$, the integral is bounded below by

$$\int_{a_t^{-c}}^{\infty} e^{-\gamma s_{\sigma}^{-1}(0)} f_{\sigma}\left(a_t^{\gamma-1} s_{\sigma}(0)\right) \, ds_{\sigma}(0) \ge e^{-\gamma a_t^{c}} \bar{F}_{\sigma}\left(a_t^{\gamma-1-c}\right) \gg e^{-\gamma a_t^{\nu}}$$

with the final asymptotic following since $\mu = 0$. Hence the integral in (39) is asymptotically concentrated on $s_{\sigma}(0) \in (a_t^{-\nu}, \infty)$. Finally, notice that for fixed $s_{\sigma}(0) \in (a_t^{-\nu}, \infty)$ we have that, as $t \to \infty$,

$$Q_t(A_t; s) = a_t^{1-\gamma} s_{\sigma}^{-1}(0) + o(d_t)$$

since $\nu < 1$, with the error uniform in $s_{\sigma}(0)$. Hence, for $s_{\sigma}(0) \in (a_t^{-\nu}, \infty)$, as $t \to \infty$,

$$\exp\{R_t(A_t;s)\} \sim t^{-d} a_t^{\gamma-1} e^{-\gamma s_{\sigma}^{-1}(0)} \prod_{s_{\xi}} f_{\xi}(s_{\xi}) \prod_{s_{\sigma}} f_{\sigma}(s_{\sigma})$$

and so the integral in (38) is asymptotically concentrated on $E_{\xi} \times E_{\sigma}$.

Case (iii) is easier to handle. Now the integral in (39) becomes

$$e^{f_1(t;s)} \int_{\mathbb{R}} \exp\{-\gamma a_t^{\gamma-1} s_{\sigma}^{-1}(0) + o(1)\} f_{\sigma}(s_{\sigma}(0)) ds_{\sigma}(0),$$

with the error uniform in s. If $\gamma < 1$, then this integral is clearly concentrated on $s_{\sigma}(0) \in (0, g_t)$. If $\gamma = 1$, then the integrand of this integral is asymptotically

$$e^{s_{\sigma}^{-1}(0)}f_{\sigma}(s_{\sigma}(0)),$$

uniformly over $s_{\sigma}(0)$, which establishes (37). The remainder of the proof is identical.

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5.2. Constructing the point process. The existence of asymptotics for the (punctured) local principal eigenvalues allows us to establish scaling limits for the penalisation functional $\Psi_t^{(j)}$. We start by constructing a point set from the pair $(z, \Psi_t^{(j)}(z))$ which will converge to a point process in the limit. Here we make use of the length scale r_t defined in equation (9), and the first and second order scales A_t and d_t for the extremes of the local principal eigenvalues defined in Proposition 5.2 and equation (16) respectively. Note that the appropriate rescaling functions for the point set are actually A_{r_t} and d_{r_t} , although since $d_{r_t} \sim d_t$ we shall eventually end up substituting these.

For technical reasons, we shall actually need to consider a certain generalisation of the functional $\Psi_t^{(j)}$. More precisely, for each $c \in \mathbb{R}$, define the functional $\Psi_{t,c}^{(j)}: V_t \to \mathbb{R}$ by

$$\Psi_{t,c}^{(j)}(z) := \lambda^{(j)}(z) - \frac{|z|}{\gamma t} \log \log t + c \frac{|z|}{t}.$$

For each $z \in \Pi^{(L_t)}$ define

$$Y_{t,c,z}^{(j)} := \frac{\Psi_{t,c}^{(j)}(z) - A_{r_t}}{d_{r_t}} \quad \text{and} \quad \mathcal{M}_{t,c}^{(j)} := \sum_{z \in \Pi^{(L_t)}} \mathbb{1}_{(zr_t^{-1}, Y_{t,c,z}^{(j)})}.$$

Finally, for each $\tau \in \mathbb{R}$ and $\alpha > -1$ let

$$\hat{H}^{\alpha}_{\tau} := \{ (x, y) \in \dot{\mathbb{R}}^{d+1} : y \ge \alpha |x| + \tau \}$$

where $\dot{\mathbb{R}}^{d+1}$ is the one-point compactification of \mathbb{R}^{d+1} .

Proposition 5.4 (Point process convergence). For each $\tau, c \in \mathbb{R}$ and $\alpha > -1$, as $t \to \infty$,

$$\mathcal{M}_{t,c}^{(j)}|_{\hat{H}_{\tau}^{\alpha}} \Rightarrow \mathcal{M} \quad in \ law,$$

where \mathcal{M} is a Poisson point process on \hat{H}^{α}_{τ} with intensity measure $\nu(dx, dy) = dx \otimes e^{-y - |x|} dy$.

Proof. The idea of the proof is to replace the set $\{\lambda^{(j)}(z)\}_{z\in\Pi^{(L_t)}}$ with the set of i.i.d. punctured principal eigenvalues $\{\tilde{\lambda}^{(j)}\}_{z\in V_t}$ and then apply standard results in i.i.d. extreme value theory to show convergence to \mathcal{M} in \hat{H}^{α}_{τ} .

To this end, define $\tilde{\Psi}_{t,c}^{(j)}(z)$ and $\tilde{Y}_{t,c,z}^{(j)}$ equivalently to $\Psi_{t,c}^{(j)}(z)$ and $Y_{t,c,z}^{(j)}$ after replacing $\lambda^{(j)}(z)$ everywhere with $\tilde{\lambda}^{(j)}(z)$ and further define

$$\tilde{\mathcal{M}}_{t,c}^{(j)} = \sum_{v \in V_t} \mathbb{1}_{(zr_t^{-1}, \tilde{Y}_{t,c,z}^{(j)})}.$$

Recall that $\{\tilde{\lambda}_t^{(j)}\}_{z \in V_t}$ are i.i.d. with tail asymptotics and uniform tail decay governed by Proposition 5.2. By applying an identical argument as in [26, Lemma 3.1] and [16, Lemma 4.3], we have that, as $t \to \infty$,

$$\tilde{\mathcal{M}}_{t,c}^{(j)}\big|_{\hat{H}_{\tau}^{\alpha}} \Rightarrow \mathcal{M} \quad \text{in law}\,.$$

Note that the uniform tail decay is necessary for the point process convergence to hold since \hat{H}^{α}_{τ} is a non-compact set (see [16, Lemma 4.3]). We claim that if $z \in V_t$ is such that

$$(zr_t^{-1}, \tilde{Y}_{t,c,z}^{(j)}) \in \hat{H}_{\tau}^{\alpha}$$

then $z \in \tilde{\Pi}^{(L_t)}$ eventually almost surely. This is since $(zr_t^{-1}, \tilde{Y}_{t,c,z}^{(j)}) \in \hat{H}_{\tau}^{\alpha}$ is equivalent to

$$\tilde{\lambda}^{(j)}(z) \ge A_{r_t} + \frac{\alpha |z| d_{r_t}}{r_t} + \frac{|z|}{\gamma t} \log \log t - \frac{c|z|}{t} + \tau d_{r_t}$$

which implies that, as $t \to \infty$,

$$\tilde{\lambda}^{(j)}(z) \ge a_t(1+o(1)) + (\alpha+1+o(1))\frac{|z|}{\gamma t}\log\log t + O(d_t) \\\ge a_t(1+o(1)) + O(d_t)$$

since $A_{r_t} \sim a_{r_t} \sim a_t$, $d_{r_t} \sim d_t$ and $\alpha > -1$. The claim then follows by the upper bound in Lemma 3.2. As a consequence, we have that, as $t \to \infty$,

$$\sum_{c \in \tilde{\Pi}^{(L_t)}} \mathbb{1}_{(zr_t^{-1}, \tilde{Y}_{t,c,z}^{(j)})} \big|_{\hat{H}_{\tau}^{\alpha}} \Rightarrow \mathcal{M} \quad \text{in law} \,.$$

$$\tag{40}$$

To complete the proof we show how to construct a coupling, valid eventually almost surely, with the property that

$$\left\{\lambda^{(j)}(z)\right\}_{z\in\Pi^{(L_t)}} = \left\{\tilde{\lambda}^{(j)}(z)\right\}_{z\in\tilde{\Pi}^{(L_t)}}.$$
(41)

To do this, note that by Lemma 4.2 there almost surely exists a t_0 such that, for all $t > t_0$, we have $r(\Pi^{(L_t)}) > 2j$ and, for each $z \in \Pi^{(L_t)}$ and $y \in B(z, j) \setminus \{z\}$, it holds that $\xi(y) < L_t$. For such t we define the coupling as follows: for $z \in \Pi^{(L_t)}$ and $y \in B(z, j)$ set $\tilde{\xi}_z(y) = \xi(y)$ and $\tilde{\sigma}_z(y) = \sigma(y)$; for $z \notin \Pi^{(L_t)}$ independently sample $\tilde{\xi}_z(z)$ as $\xi(0)$ conditioned on $\xi(0) < L_t$; and otherwise independently sample $\tilde{\xi}_z(y)$ and $\tilde{\sigma}_z(y)$ according to their respective law. By the separation guaranteed for $t > t_0$, $\{\tilde{\xi}_z\}_{z \in V_t}$ and $\{\tilde{\sigma}_z\}_{z \in V_t}$ are indeed i.i.d. fields with law as in (32), and moreover (41) holds by construction. Combining with (40) completes the proof.

Remark 5.5. Although we state Proposition 5.4 for arbitrary $c \in \mathbb{R}$, we shall only apply it to c = 0 and one other value of c that will be determined in Section 6.

We now use the point process \mathcal{M} to analyse the joint distribution of top two statistics of the functional $\Psi_{t,c}^{(j)}$. So let

$$Z_{t,c}^{(j)} := \underset{z \in \Pi^{(L_t)}}{\arg \max} \Psi_{t,c}^{(j)}(z) \quad \text{and} \quad Z_{t,c}^{(j,2)} := \underset{\substack{z \in \Pi^{(L_t)}\\z \neq Z_{t,c}^{(j)}}}{\arg \max} \Psi_{t,c}^{(j)} .$$

Note that eventually these are well-defined almost surely, since $\Pi^{(L_t)}$ is finite and non-zero by Lemma 4.1.

Corollary 5.6. For each $c \in \mathbb{R}$, as $t \to \infty$,

$$\left(\frac{Z_{t,c}^{(j)}}{r_t}, \frac{Z_{t,c}^{(j,2)}}{r_t}, \frac{\Psi_{t,c}^{(j)}(Z_{t,c}^{(j)}) - A_{r_t}}{d_{r_t}}, \frac{\Psi_{t,c}^{(j)}(Z_{t,c}^{(j,2)}) - A_{r_t}}{d_{r_t}}\right)$$

converges in law to a random vector with density

$$p(x_1, x_2, y_1, y_2) = \exp\{-(y_1 + y_2) - |x_1| - |x_2|) - 2^d e^{-y_2}\} \mathbb{1}_{\{y_1 > y_2\}}$$

Proof. This follows from the point process density in Proposition 5.4 using the same computation as in [26, Proposition 3.2]. \Box

5.3. Properties of the localisation site. In this subsection we use the results from the previous subsection to analyse the localisation sites $Z_{t,c}^{(j)}$ and Z_t , and in the process complete the proof of Theorems 1.9 and 1.10. For each $c \in \mathbb{R}$, introduce the events

$$\mathcal{G}_{t,c} := \{ \Psi_{t,c}^{(j)}(Z_{t,c}^{(j)}) - \Psi_{t,c}^{(j)}(Z_{t,c}^{(j,2)}) > d_t e_t \},\$$
$$\mathcal{H}_t := \{ r_t f_t < |Z_t^{(j)}| < r_t g_t \} \quad \text{and} \quad \mathcal{I}_t := \{ a_t (1 - f_t) < \Psi_t^{(j)}(Z_t^{(j)}) < a_t (1 + f_t) \},\$$
event

and the event

$$\mathcal{E}_{t,c} := \mathcal{S}_t^{(j)}(Z_t^{(j)}) \cap \mathcal{G}_{t,0} \cap \mathcal{G}_{t,c} \cap \mathcal{H}_t \cap \mathcal{I}_t$$
(42)

which acts to collect the relevant information that we shall later need.

Proposition 5.7. For each $c \in \mathbb{R}$, as $t \to \infty$,

$$\mathbb{P}(\mathcal{E}_{t,c}) \to 1$$
.

Proof. This follows from Proposition 5.2, Corollary 5.6, the coupling in the proof of Proposition 5.4, and since $A_{r_t} \sim a_t$ and $d_{r_t} \sim d_t$.

In the next few propositions, we prove that the sites $Z_{t,c}^{(j)}$ and $Z_t^{(j)}$ are both equal to the localisation site Z_t with overwhelming probability.

Proposition 5.8. For each $c \in \mathbb{R}$, on the event $\mathcal{E}_{t,c}$, as $t \to \infty$,

$$Z_{t,c}^{(j)} = Z_t^{(j)}$$

holds eventually.

Proof. Assume that $Z_{t,c}^{(j)} \neq Z_t^{(j)}$ and recall that $1/\log\log t < e_t/g_t$ eventually by (17). On the event $\mathcal{E}_{t,c}$, the statements

$$\Psi_t^{(j)}(Z_t^{(j)}) - \Psi_t^{(j)}(Z_{t,c}^{(j)}) > d_t e_t \quad \text{and} \quad \Psi_{t,c}^{(j)}(Z_{t,c}^{(j)}) - \Psi_{t,c}^{(j)}(Z_t^{(j)}) > d_t e_t$$

and, eventually,

$$|\Psi_t^{(j)}(Z_t^{(j)}) - \Psi_{t,c}^{(j)}(Z_t^{(j)})| = |c| \frac{|Z_t^{(j)}|}{t} < \gamma \frac{d_t g_t}{\log \log t} < d_t e_t$$

all hold, giving a contradiction.

Lemma 5.9. For each $c \in \mathbb{R}$, on the event $\mathcal{E}_{t,c}$, as $t \to \infty$,

$$\lambda^{(j)}(Z_t^{(j)}) \ge \lambda(Z_t^{(j)}) \quad and \quad \lambda^{(j)}(Z_t) \ge \lambda(Z_t)$$

and

$$\lambda^{(j)}(Z_t^{(j)}) - \lambda(Z_t^{(j)}) < d_t e_t$$

all hold eventually.

Proof. The first two statements follow from the domain monotonicity of the principal eigenvalue in Lemma 3.1. For the third statement, remark that the event $\mathcal{E}_{t,c}$ implies that $Z_t^{(j)} \in \Pi^{(L_{t,\varepsilon})}$, that $\xi(y) < L_t$ for all $y \in B(Z_t^{(j)}, \rho)$, that $\xi(y) < g_t$ for all y such that $j \ge |y - Z_t^{(j)}| > \rho_{\xi}$, and that $\sigma(Z_t^{(j)}) > a_t^{q_\sigma} f_t$. Hence, by considering the path expansion in Proposition 5.1 (valid by the coupling in the proof of Proposition 5.4), we have that for some C > 0,

$$\lambda^{(j)}(Z_t^{(j)}) - \lambda(Z_t^{(j)}) < \frac{Ca_t^{-(\frac{\gamma-1}{\mu+1})^+}g_t}{f_t(L_{t,\epsilon} - L_t)^{2\rho+1}} < d_t e_t$$
(43)

eventually, with the last equality holding since

$$-2\rho - 1 - \left(\frac{\gamma - 1}{\mu + 1}\right)^+ < 1 - \gamma.$$
⁽⁴⁴⁾

Remark 5.10. Note that ρ is precisely the smallest integer such that (44) holds.

Corollary 5.11 (Equivalence of $Z_t^{(j)}$ and Z_t). For each $c \in \mathbb{R}$, on the event $\mathcal{E}_{t,c}$, as $t \to \infty$,

$$Z_t^{(j)} = Z_t$$

eventually.

Proof. Assume that $Z_t^{(j)} \neq Z_t$. On the event $\mathcal{E}_{t,c}$, Lemma 5.9 implies that

$$\begin{pmatrix} \Psi_t^{(j)}(Z_t^{(j)}) - \Psi_t(Z_t^{(j)}) \end{pmatrix} - \left(\Psi_t^{(j)}(Z_t) - \Psi_t(Z_t) \right) \\ = \left(\lambda^{(j)}(Z_t^{(j)}) - \lambda(Z_t^{(j)}) \right) - \left(\lambda^{(j)}(Z_t) - \lambda(Z_t) \right) < d_t e$$

holds eventually. On the other hand, on the event $\mathcal{E}_{t,c}$, and by the definition of Z_t and $Z_t^{(j)}$ as the argmax of Ψ_t and $\Psi_t^{(j)}$ respectively,

$$\Psi_t^{(j)}(Z_t^{(j)}) - \Psi_t^{(j)}(Z_t) > d_t e_t \quad \text{and} \quad \Psi_t(Z_t) - \Psi_t(Z_t^{(j)}) > 0$$

also hold, giving a contradiction.

Finally, we prove a criterion for the independence of Z_t from the trapping landscape σ . Define $\psi_t(z) := \xi(z) - \frac{|z|}{\gamma t} \log \log t$, and let $z_t := \arg \max_{z \in \Pi^{(L_t)}} \psi_t(z)$. Note that z_t is independent of σ .

Proposition 5.12 (Criterion for the independence of Z_t from the trapping landscape σ). If $\gamma < 1$, then as $t \to \infty$,

$$\mathbb{P}\left(Z_t = z_t\right) \to 1$$

Proof. By Proposition 5.7 we may assume $\mathcal{E}_{t,c}$ holds. Observe that, on $\mathcal{E}_{t,c}$ and by Proposition 5.1, any $z \in \Pi^{(L_{t,\varepsilon})} \setminus \{Z_t^{(j)}\}$ satisfies

$$\psi_t(Z_t^{(j)}) > \Psi_t^{(j)}(Z_t^{(j)}) > \Psi_t^{(j)}(z) + d_t e_t > \psi_t(z) + O(1) + d_t e_t$$
.

Moreover, by Lemma 3.2 and on $\mathcal{E}_{t,c}$, any $z \in \Pi^{(L_t)} \setminus \Pi^{(L_{t,\varepsilon})}$ also satisfies

$$\psi_t(Z_t^{(j)}) > \Psi_t^{(j)}(Z_t^{(j)}) > \psi_t(z) + O(1) + d_t e_t$$
.

Since $d_t e_t \to \infty$ if $\gamma < 1$, this implies that $Z_t^{(j)} = \arg \max_{z \in \Pi^{(L_t)}} \psi_t(z) =: z_t$. Corollary 5.11 completes the proof.

5.4. **Proof of Theorem 1.9.** We prove Theorem 1.9 on the event $\mathcal{E}_{t,c}$, since by Proposition 5.7 this event holds with overwhelming probability eventually. Part (a) is implied directly by the definition of the event $\mathcal{E}_{t,c}$. Parts (b)–(d) follow by combining the definition the event $\mathcal{E}_{t,c}$ with Proposition 5.3 and applying the coupling in the proof of Proposition 5.4. Finally, part (e) is a consequence of the point process convergence, and is proved in an identical manner to the corresponding results in [10, 26].

5.5. **Proof of Theorem 1.10.** Consider parts (a) and (b). By definition, Z_t depends only on the values of ξ and σ in balls of radius ρ_{ξ} and ρ respectively around each site, and so the radii ρ_{ξ} and ρ are certainly sufficient. To show necessity, consider that the results in parts (b)–(d) of Theorem 1.9 establish the correlation of the fields ξ and σ at a distance ρ_{ξ} and ρ respectively around Z_t . Hence these radii are necessary as well.

Consider then part (c). The sufficient condition for the reduction to ξ follows directly from Proposition 5.12. To show necessity, consider that the results in part (c) of Theorem 1.9 establish that, if $\gamma \geq 1$, the value of $\sigma(Z_t)$ is not an independent copy of $\sigma(0)$, and hence Z_t must depend on σ .

It remains to prove part (d). If $\rho = 0$ then Z_t depends only on η by definition. On the other hand, suppose $\rho \geq 1$ and, for the purposes of contradiction, that there exists a random site z_t , depending only on ξ and σ through η , such that, as $t \to \infty$,

$$\mathbb{P}(Z_t = z_t) \to 1.$$

Fix a site y and a constant $c > \delta_{\sigma}$. We establish a contradiction by considering two bounds on the probability of the event

$$\{\sigma(y) < c, |Z_t - y| = 1\}.$$

We first consider the case $(\gamma, \mu) \notin \mathcal{B}_{\sigma}$. Then by part (d) of Theorem 1.9, conditionally on event $\{|Z_t - y| = 1\}$, we have that $\sigma(y) \to \delta_{\sigma}$ in probability as $t \to \infty$. This implies that there exists some $c_1 > 0$ such that

$$\mathbb{P}(\sigma(y) < c, |Z_t - y| = 1) > (\mathbb{P}(\sigma(y) < c) + c_1) \mathbb{P}(|Z_t - y| = 1)$$
(45)

eventually. In the case $(\gamma, \mu) \in \mathcal{B}_{\sigma}$, conditionally on event $\{|Z_t - y| = 1\}$ and again by part (d) of Theorem 1.9,

$$f_{\sigma(y)}(x) \to c_2 e^{\bar{c}_{\sigma}/x} f_{\sigma}(x)$$

for some $c_2 > 0$, and so (45) holds in this case as well.

We now work on the event $\{Z_t = z_t\}$ and show how to obtain a lower bound on the probability of the event $\{\sigma(y) < c, |z_t - y| = 1\}$. Let $\bar{\eta} = \{\eta(v) : v \neq y\}$. Remark first that, since $z_t \in \Pi^{(L_t)}$, by Proposition 5.1 we have that $\lambda_t(z_t)$ is increasing in $\eta(y)$ for $|y - z_t| = 1$. Hence there exists a function $\beta_t : \bar{\eta} \to \mathbb{R} \cup \{\infty\}$ such that, conditionally on $\bar{\eta}$,

$$\{|z_t - y| = 1\}$$
 and $\{\eta(y) \ge \beta_t(\bar{\eta})\}$

agree almost surely. To see this, set $\beta_t(\bar{\eta})$ to be the minimum $\eta(y)$ such that with such a value of $\eta(y)$, we have $|z_t - y| = 1$ (and setting it to be infinity if no such value exists). Then clearly, if

 $\eta(y) < \beta_t(\bar{\eta})$ we cannot have $|z_t - y| = 1$, and on the other hand we claim that if $\eta(y) \ge \beta_t(y)$ we have $|z_t - y| = 1$. This follows by the almost-sure separation of Lemma 4.2, which ensures that $\{y = z_t\}$ has probability 0. Denote by $\mathcal{F}_{\bar{\eta}}$ the σ -algebra generated by $\bar{\eta}$. Then, eventually almost surely,

$$\begin{split} \mathbb{P}(\sigma(y) < c, \, |z_t - y| = 1) &= \mathbb{E}_{\bar{\eta}} \left[\mathbb{E}[\mathbb{1}_{\{|z_t - y| = 1\}} \mathbb{1}_{\{\sigma(y) < c\}} | \mathcal{F}_{\bar{\eta}}] \right] \\ &= \mathbb{E}_{\bar{\eta}} \left[\mathbb{E}[\mathbb{1}_{\{\eta(y) > \beta_t(\bar{\eta})\}} \mathbb{1}_{\{\sigma(y) < c\}} | \mathcal{F}_{\bar{\eta}}] \right] \\ &\leq \mathbb{E}_{\bar{\eta}} \left[\mathbb{E}[\mathbb{1}_{\{\eta(y) > \beta_t(\bar{\eta})\}} | \mathcal{F}_{\bar{\eta}}] \mathbb{E}[\mathbb{1}_{\{\sigma(y) < c\}} | \mathcal{F}_{\bar{\eta}}] \right] \\ &= \mathbb{P}(\sigma(y) < c) \mathbb{P}(|z_t - y| = 1) \,, \end{split}$$

where the second equality uses the fact that z_t depends on σ only through η , and the inequality holds since, conditionally on $\mathcal{F}_{\bar{\eta}}$, the events $\{\eta(y) > \beta_t(\bar{\eta})\}$ and $\{\sigma(y) < c\}$ are negatively correlated. Since $z_t = Z_t$ with probability going to one, combining with (45) gives the required contradiction.

6. Negligible paths

In this section we show that the contribution to the total mass U(t) from the components $U^2(t)$, $U^3(t)$, $U^4(t)$ and $U^5(t)$ are all negligible. We proceed in two parts: first we prove a lower bound on the total mass U(t), and then we bound from above the contribution to the total mass from each $U^i(t)$. Throughout this section, let ε be such that $0 < \varepsilon < \theta$.

6.1. **Preliminaries.** We begin by proving a general result on eigenfunction decay around sites of high potential, which will be used in both the lower and upper bound. For each $z \in \Pi^{(L_{t,\varepsilon})}$, let φ_1 denote the principal eigenfunction of the Hamiltonian $\mathcal{H}^{(j)}(z)$.

Proposition 6.1. For each $z \in \Pi^{(L_{t,\varepsilon})}$ uniformly, as $t \to \infty$, almost surely

$$\sum_{y \in B(z,j) \setminus \{z\}} \varphi_1(y) \to 0 \qquad and \qquad \sum_{y \in B(z,j) \setminus \{z\}} \frac{\sigma(y)^{-\frac{1}{2}} \varphi_1(y)}{||\sigma^{-\frac{1}{2}} \varphi_1||_{\ell_2}} \to 0 \,.$$

Proof. By Proposition 3.5, we have the path expansion

$$\frac{\varphi_1(y)}{\varphi_1(z)} = \frac{\sigma(y)}{\sigma(z)} \sum_{k \ge 1} \sum_{\substack{p \in \Gamma_k(y,z)\\ p_i \ne z, \ 0 \le i < k\\ \operatorname{Set}(p) \subseteq \overline{B}(z,j)}} \prod_{\substack{0 \le i < k\\ \operatorname{Set}(p) \subseteq \overline{B}(z,j)}} (2d)^{-1} \frac{\sigma^{-1}(p_i)}{\lambda^{(j)}(z) - \eta(p_i)} , \quad y \in B(z,j) \setminus \{z\}.$$
(46)

Since, by Lemmas 4.2 and 3.2, for each $y \in B(z, j) \setminus \{z\}$, almost surely

$$\lambda^{(j)}(z) - \eta(y_i) > L_{t,\varepsilon} - L_t - \delta_{\sigma}^{-1},$$

and moreover since $\sigma^{-1}(y) < \delta_{\sigma}^{-1}$ for all $y \in B(z, j)$, the result follows.

Corollary 6.2 (Bound on total mass of the solution). For each $z \in \Pi^{(L_{t,\varepsilon})}$ uniformly and any c > 1, as $t \to \infty$, almost surely

$$\mathbb{E}_{z}\left[e^{\int_{0}^{t}\xi(X_{s})\,ds}\mathbb{1}_{\{\tau_{B(z,j)^{c}}>t\}}\right] < c\,e^{t\lambda^{(j)}(z)}$$

eventually.

Proof. This follows by combining Propositions 6.1 and 3.12.

6.2. Lower bound on the total mass U(t). Recall that by the discussion in Section 2, the total mass U(t) can be approximated by considering both the benefit of being near a site of high potential and the probabilistic penalty from diffusing to that site. To formalise a lower bound for U(t) we need a bound on both of these terms.

We begin by bounding from below the benefit to the solution from paths that start and end at a site of high potential.

Lemma 6.3. For each $z \in \Pi^{(L_{t,\varepsilon})}$ uniformly,

$$\log u_z(t,z) \ge t\lambda^{(j)}(z) + o(1)$$

eventually almost surely.

Proof. Recall the Feynman-Kac formula for the solution $u_z(t, z)$ (see, e.g., Proposition 3.8), and note that the expectation is larger than the corresponding expectation taken only over paths that do not leave B(z, j). Using Corollary 3.11, we then have that

$$u_{z}(t,z) \geq \frac{e^{\lambda^{(j)}(z)t}\sigma^{-1}(z)\varphi_{1}^{2}(z)}{||\sigma^{-\frac{1}{2}}\varphi_{1}||_{\ell_{2}}^{2}}$$

where φ_1 denotes the principal eigenfunction of the Hamiltonian $\mathcal{H}^{(j)}(z)$. Since the domain B(z, j) is finite, the fact that the eigenfunction $\sigma^{-\frac{1}{2}}\varphi_1$ is localised at z (by Proposition 6.1) ensures that the square eigenfunction $\sigma^{-1}\varphi_1^2$ is also localised at z, and the result follows.

The next step is to bound from above the probabilistic penalty incurred by diffusing to a certain site. This will be a function both of the distance of the site from the origin, as well as the size of the traps on paths from the origin to the site. Here we use the existence of quick paths that we established in a general setting in Section 4.

Recall the scaling function s_t , which satisfies the properties in (17). If d = 1, for $\sigma_t := s_t$ and $n_t := r_t g_t$, recall the definitions of I_t and $\{\sigma_t^i\}$ from Proposition 4.8. Let $p \in \Gamma_{|Z_t|}(0, Z_t)$ be the (unique) shortest path from 0 to Z_t and define

$$N_i^p := \sum_{0 \le l < |Z_t|} \mathbb{1}_{\{\sigma(p_l) \in (\sigma_t^{i-1}, \sigma_t^i]\}}, \quad i = 1, \dots, I_t.$$

If $d \ge 2$, for $z_t := Z_t$, $\sigma_t := s_t$ and $S_t := \Pi^{(L_t)}$, recall the definition of $|Z_t|_{\text{chem}}$ from Proposition 4.11. Denote by Θ_t^d the event

$$\Theta_t^d := \begin{cases} \left\{ \sum_{i=1}^{I_t} N_i^p \log \sigma_t^i < td_t b_t \,, \, \max_{0 \le l < |Z_t|} \sigma(p_l) < \sigma_t^{I_t} \right\} \,, & d = 1 \,; \\ \left\{ |Z_t|_{\text{chem}} < |Z_t| + r_t b_t \right\} \,, & d \ge 2 \,. \end{cases}$$

Proposition 6.4 (Existence of quick paths). For each $c \in \mathbb{R}$, as $t \to \infty$,

$$\mathbb{P}(\Theta_t^d, \mathcal{E}_{t,c}) \to 1.$$

Proof. Recall that on event $\mathcal{E}_{t,c}$ we have that $|Z_t| < r_t g_t$. Suppose d = 1. Then the result follows immediately from Proposition 4.8 and the properties of the scaling function s_t in (17), since

$$\log \log r_t g_t \sim \log \log t$$
.

Suppose then $d \geq 2$. Note that conditioning on ξ determines $\Pi^{(L_t)}$ and also that, by Lemma 4.2, eventually almost surely $\Pi^{(L_t)}$ satisfies the properties required by the set S_t . Since $Z_t \in \Pi^{(L_t)}$, conditioning on the values of σ in $B(\Pi^{(L_t)}, j)$ therefore determines Z_t . Given Z_t and $\Pi^{(L_t)}$, the event Θ_t^d is fully determined by the values of σ in $\mathbb{Z}^d \setminus B(\Pi^{(L_t)}, j)$. Hence we can apply Proposition 4.11 with $z_t = Z_t$, $\sigma_t = s_t$ and $S_t = \Pi^{(L_t)}$, to deduce that there exists a $c_1 < 1$ such that, for all functions $c_t \to \infty$ such that $\overline{F}_{\sigma}(s_t)c_t \ll 1$,

$$|Z_t|_{\text{chem}} < |Z_t|(1 + \bar{F}_{\sigma}(s_t)c_t + t^{-c_1})|$$

with probability tending to 1. By (17), we can pick a c_t such that

$$r_t g_t F_\sigma(s_t) c_t \ll r_t b_t \,,$$

and so we have the result.

We are now ready to prove the lower bound.

Proposition 6.5. For each $c \in \mathbb{R}$, on the events $\mathcal{E}_{t,c}$ and Θ_t^d , as $t \to \infty$,

$$\log U(t) \ge t\lambda^{(j)}(Z_t) - \frac{|Z_t|}{\gamma} \log \log t + O(td_tb_t)$$

almost surely.

Proof. In the following proof set $z = Z_t$ and abbreviate $\tau = \tau_z$. We first consider the case of $d \ge 2$. By the Feynman-Kac formula (5), the total mass U(t) can be written as

$$U(t) = \mathbb{E}_0 \left[\exp \left\{ \int_0^t \xi(X_s) ds \right\} \right]$$

Using the non-negativity of ξ and by the strong Markov property, we have, for each $r \in (0, 1)$,

$$U(t) \geq \mathbb{E}_{0} \Big[\exp \Big\{ \int_{0}^{t} \xi(X_{s}) ds \Big\} \mathbb{1}_{\{\tau < rt\}} \Big]$$

$$\geq \mathbb{E}_{0} \Big[\exp \Big\{ \int_{\tau}^{t-(rt-\tau)} \xi(X_{s}) ds \Big\} \mathbb{1}_{\{\tau < rt\}} \Big] = \mathbb{E}_{0} \Big[\exp \Big\{ \int_{\tau}^{t-(rt-\tau)} \xi(X_{s}) ds \Big\} \Big] \mathbb{P}_{0}(\tau < rt)$$

$$\geq u_{z}((1-r)t, z) \mathbb{P}_{0}(\tau < rt).$$
(47)

We now seek to bound $\mathbb{P}_0(\tau < rt)$. Since we are on event Θ_t^d , there exists a path

$$p \in \bigcup_{y \in \partial B(z,j)} \Gamma_{\ell_t}(0,y)$$

for some $\ell_t < |z| + r_t b_t$ such that $\sigma(x) < s_t$ for all $x \in \operatorname{Set}(p)$. Moreover, since we are on event $\mathcal{E}_{t,c}$, each $\sigma(x) \in B(z,j) \setminus \{z\}$ is such that $\sigma(x) < a_t^{\nu}$ for some $\nu \in (0,1)$. We shall denote by $\{\tilde{X}_t\}_{t \in \mathbb{R}^+}$ a random walk with generator $\Delta \tilde{\sigma}^{-1}$, where $\tilde{\sigma}(x) = s_t$ for all $x \in \operatorname{Set}(p)$, $\tilde{\sigma}(x) = a_t^{\nu}$ for all $x \in B(z,j) \setminus \{z\}$, and $\tilde{\sigma}(x) = \sigma(x)$ otherwise. By a simple coupling argument we have that

$$\mathbb{P}_0(\tau < rt) \ge \mathbb{P}_0(\tilde{\tau} < rt), \tag{48}$$

where $\tilde{\tau}$ is the first hitting time of z by \tilde{X} . Using a similar calculation as in [20][Proposition 4.2], for any $r_1 + r_2 \leq r$,

$$\begin{aligned} \mathbb{P}_{0}(\tilde{\tau} < rt) > (2d)^{-\ell_{t}-j} \mathbb{P}\left(\operatorname{Poi}(r_{1}ts_{t}^{-1}) = \ell_{t}\right) \mathbb{P}\left(\operatorname{Poi}(r_{2}ta_{t}^{-\nu}) = j\right) \\ = (2d)^{-\ell_{t}-j} e^{-r_{1}ts_{t}^{-1}} \frac{(r_{1}ts_{t}^{-1})^{\ell_{t}}}{(\ell_{t})!} e^{-r_{2}ta_{t}^{-\nu}} \frac{(r_{2}ta_{t}^{-\nu})^{j}}{j!}. \end{aligned}$$

Applying Stirling's formula, we obtain

$$\log \mathbb{P}_0(\tilde{\tau} < rt) \ge -r_1 t s_t^{-1} - r_2 t a_t^{-\nu} - \ell_t \log \left(\frac{2d\,\ell_t}{er_1 t s_t^{-1}}\right) + j \log r_2 + O(\log t)\,. \tag{49}$$

Now note that on the event $\mathcal{E}_{t,c}$ we have that $Z_t \in \Pi^{(L_{t,\varepsilon})}$. Hence we can combine equations (47)–(49) and Lemma 6.3 to get that

$$\log U(t) \ge (1 - r_1 - r_2)t\lambda^{(j)}(z) - r_1ts_t^{-1} - r_2ta_t^{-\nu} - \ell_t \log\left(\frac{2d\,\ell_t}{er_1ts_t^{-1}}\right) + j\log r_2 + O(\log t)\,.$$

Use the bound $\ell_t < |z| + r_t b_t$ and choose $r = r_1 + r_2$ to maximise this equation, that is, set

$$r_1 := \frac{|z| + r_t b_t}{t(\lambda^{(j)}(z) + s_t^{-1})}$$
 and $r_2 := \frac{j}{t(\lambda^{(j)}(z) + a_t^{-\nu})}$

It is clear that on event $\mathcal{E}_{t,c}$ we have $r \in (0,1)$. With these values of r_1 and r_2 we obtain

$$\log U(t) \ge t\lambda^{(j)}(z) - (|z| + r_t b_t) \Big\{ \log \Big(\frac{\lambda^{(j)}(z) + s_t^{-1}}{s_t^{-1}} \Big) + O(1) \Big\} + O(\log t) \,.$$

On event $\mathcal{E}_{t,c}$ we have that $\lambda^{(j)}(z) < a_t(1+f_t)$. Since also $|z| < r_t g_t$ on event $\mathcal{E}_{t,c}$ we find that

$$\log U(t) \ge t\lambda^{(j)}(z) - |z|\log(\lambda^{(j)}(z)) - r_t b_t \log(\lambda^{(j)}(z)) + O\left(r_t g_t \log(s_t)\right)$$
$$\ge t\lambda^{(j)}(z) - \frac{|z|}{\gamma}\log\log t + O\left(td_t b_t\right)$$

by the choice of the scaling functions s_t in equation (17).

Next, we turn to the case d = 1. Denote by $\{\bar{X}_t\}_{t \in \mathbb{R}^+}$ a random walk with generator $\Delta \bar{\sigma}^{-1}$ where $\bar{\sigma}(x) = \sigma_t^i$ if $\sigma(x) \in (\sigma_t^{i-1}, \sigma_t^i]$. Again, by a simple coupling argument

$$\mathbb{P}_0(\tau < rt) \ge \mathbb{P}_0(\bar{\tau} < rt)$$

where $\bar{\tau}$ is the first hitting time of z by \bar{X} and $r \in (0, 1)$. Furthermore, we have

$$\mathbb{P}_0(\bar{\tau} < rt) > 2^{-|Z_t|} \prod_{i=1}^{n_t} \mathbb{P}(\operatorname{Poi}(r_i t(\sigma_t^i)^{-1}) = N_i^p),$$

for any $\{r_i\}_{1 \le i \le I_t}$ satisfying $\sum_i r_i \le r$. By a similar calculation to the $d \ge 2$ case, we have

$$\log U(t) \ge t(1-r)\lambda^{(j)}(z) + \sum_{i=1}^{I_t} \left(-r_i t(\sigma_t^i)^{-1} - N_i^p \log(2N_i^p / (er_i t(\sigma_t^i)^{-1})) \right) + O(\log t) \,.$$

Choose r and $\{r_i\}$ to maximise this equation, that is, set

$$r_i = \frac{N_i^p}{t(\lambda^{(j)}(z) + (\sigma_t^i)^{-1})} \quad \text{and} \quad r = \sum_i r_t$$

noting that $r \in (0, 1)$ for the same reason as in the $d \ge 2$ case. Then,

$$\log U(t) \ge t\lambda^{(j)}(z) + \sum_{i=1}^{I_t} \left(-N_i^p \left(\log \left(\lambda^{(j)}(z)\sigma_t^i \right) \right) + O(1) \right) + O(\log t)$$

= $t\lambda^{(j)}(z) - |z| \log \left(\lambda^{(j)}(z) \right) - \sum_{i=1}^{I_t} \left(N_i^p \log \sigma_t^i + O(|z|) \right) + O(\log t).$

The result follows since we are on event Θ_t^d .

6.3. Contribution from each $U^{i}(t)$ is negligible. In this section we prove that the contribution to U(t) from the each of the components $U^{i}(t)$, for i = 2, 3, 4, 5, is negligible. The most difficult step is bounding the contribution from the components $U^{2}(t)$ and $U^{3}(t)$.

The difficulty with these components is that paths are permitted to visit sites of high potential that are not Z_t . Away from these sites, there is a probabilistic penalty associated with each step of the path; this is easy to bound. However, close to these sites, the maximum contribution from the path may come from a complicated sequence of return cycles to the site. This motivates our set-up, which groups paths into equivalence classes depending only on their trajectory away from sites of high potential.

For each t, we define a partition of paths into equivalence classes as follows. Suppose $p, \bar{p} \in \Gamma$ are two finite paths in \mathbb{Z}^d . Define inductively, $r^0 = 0$, and

$$s^{\ell} := \min\{i \ge r^{\ell-1}: \, p_i \in \Pi^{(L_t)}\} \quad \text{and} \quad r^{\ell} := \min\{i > s^{\ell}: \, p_i \in \partial B(p_{s^{\ell}}, j)\}$$

for each $\ell \in \mathbb{N}$, setting each to be ∞ if no such minimum *i* exists, and define similarly $(\bar{s}^{\ell}, \bar{r}^{\ell})_{\ell \geq 1}$ for path \bar{p} . Then we say that p and \bar{p} are in the same equivalence class if and only if, for all $\ell \geq 0$,

$$s^{\ell+1} - r^{\ell} = \bar{s}^{\ell+1} - \bar{r}^{\ell} \quad \text{and} \quad p_{r^{\ell}+i} = \bar{p}_{\bar{r}^{\ell}+i} , \quad \text{for each } i \in \{0, 1, \dots, s^{\ell+1} - r^{\ell}\}.$$

Note that although s^{ℓ} and r^{ℓ} depend on t (through the set $\Pi^{(L_t)}$), we suppress this dependence for clarity. If p and \bar{p} are in the same equivalence class at time t we write $p \sim \bar{p}$. Denote by $P(p) := \{\bar{p} \in \Gamma : p \sim \bar{p}\}$. Informally, the equivalence class P(p) consists of paths that have identical trajectory except for when they are in balls of radius j around sites $z \in \Pi^{(L_t)}$ (or, more accurately, when they first hit a site $z \in \Pi^{(L_t)}$ until when they leave the ball B(z, j)).

It is natural to group these equivalence classes P(p) according to (i) how many balls of radius j around sites $z \in \Pi^{(L_t)}$ the path visits; and (ii) the total length of the path outside such balls. So for $m, n \in \mathbb{N}$, let $\mathcal{P}_{n,m}$ be the set of equivalence classes P(p) of paths p that satisfy

$$\max\{\ell : r^{\ell} < \infty\} = m \quad \text{and} \quad \sum_{\ell=0}^{m-1} (s^{\ell+1} - r^{\ell}) + s^{m+1} \mathbb{1}_{\{s^{m+1} < \infty\}} + |p| \mathbb{1}_{\{s^{m+1} = \infty\}} - r^m = n.$$

Note that if a path p satisfies these two properties for some m and n then any other path $\bar{p} \in P_p$ will also satisfy these properties for the same m and n and hence $\mathcal{P}_{n,m}$ is well-defined. The quantity m counts the number of balls of radius j around $z \in \Pi^{(L_t)}$ that the path *exits* (which is easier to work with than the number of balls the path enters); the quantity n counts the total length of the path between leaving each of these balls and hitting the next site $z \in \Pi^{(L_t)}$.

Recalling the definitions of $p(X_t)$, define the event

$$\{p(X) \in P(p)\} := \bigcup_{s \ge 0} \{p(X_s) \in P(p)\},\$$

and remark that we have the relationship

$$\{p(X_t) \in P(p)\} \subseteq \{p(X) \in P(p)\}.$$
(50)

Denote by

$$U^{P(p)}(t) = \mathbb{E}_0 \left[\exp\left\{ \int_0^t \xi(X_s) \, ds \right\} \, \mathbb{1}_{\{p(X_t) \in P(p)\}} \right].$$

the contribution to the total solution U(t) from the path equivalence class P(p).

The following lemma bounds the contribution of each $P(p) \in \mathcal{P}_{n,m}$ in terms of m and n. The key fact motivating our set-up is that the contribution is decreasing in n.

Lemma 6.6 (Bound on the contribution from each equivalence class). Let $m, n \in \mathbb{N}$ and $p \in \Gamma(0)$ such that $\operatorname{Set}(p) \subseteq V_t$ and $P(p) \in \mathcal{P}_{n,m}$. Define $z^{(p)} := \operatorname{arg} \max_{z \in \operatorname{Set}(p)} \lambda^{(j)}(z)$ and let $\zeta > \max\{\lambda^{(j)}(z^{(p)}), L_{t,\varepsilon}\}$. Then there exist constants $c_1, c_2 > 0$ such that, uniformly in m, n, p and ζ , as $t \to \infty$,

$$U^{P(p)}(t) \le e^{\zeta t} \left(c_1(\zeta - L_t) \right)^{-n} \left(1 + c_2 \left(\zeta - \lambda^{(j)}(z^{(p)}) \right)^{-1} \right)^n$$

eventually almost surely.

Proof. The strategy of the proof is to split $U^{P(p)}(t)$ into three components, corresponding to the contribution: (i) from when X_s is outside $B(\Pi^{(L_t)}, j)$ until X_s hits a site $z \in \Pi^{(L_t)}$; (ii) from when X_s hits $z \in B(\Pi^{(L_t)}, j)$ until when X_s leaves the ball B(z, j); and (iii) if X_s hits $z \in \Pi^{(L_t)}$ and does not subsequently leave B(z, j), from this component separately. To bound the contribution from these components, we make use of Corollary 6.2, Lemma 3.13 and Lemma 3.4 respectively.

There are two cases to consider, depending on whether the event described in (iii) occurs, that is, if $s^{m+1} < \infty$. We begin with this case. To simplify notation in the following we abbreviate

$$I_a^b := \exp\left\{\int_a^b (\xi(X_s) - \zeta) \, ds\right\}$$

Recall the definition of $(s^{\ell}, r^{\ell})_{\ell \in \mathbb{N}}$ and define the stopping times

$$R^0 := 0 \ , \quad S^\ell := \inf\{s \ge R^{\ell-1}: \, X_s = p_{s^\ell}\} \quad \text{and} \quad R^\ell := \inf\{s \ge S^\ell: \, X_s = p_{r^\ell_t}\}$$

for each $\ell \in \{1, \ldots, m\}$, and similarly define S^{m+1} since $s^{m+1} < \infty$. We can then write

$$U^{P(p)}(t) = \mathbb{E}_{0} \left[e^{\int_{0}^{t} \xi(X_{s}) \, ds} \mathbb{1}_{\{p(X_{t}) \in P(p)\}} \right] = e^{\zeta t} \mathbb{E}_{0} \left[I_{0}^{t} \mathbb{1}_{\{p(X_{t}) \in P(p)\}} \right]$$
$$= e^{\zeta t} \mathbb{E}_{0} \left[\left(\prod_{\ell=0}^{m} I_{R^{\ell}}^{S^{\ell+1}} \right) \left(\prod_{\ell=1}^{m} I_{S^{\ell}}^{R^{\ell}} \right) I_{S^{m+1}}^{t} \mathbb{1}_{\{p(X_{t}) \in P(p)\}} \right].$$

Note that, conditionally on $\mathcal{F}_{S^{m+1}}$ (the σ -algebra generated by S^{m+1}), the quantity $I_{S^{m+1}}^t$ is independent of all other I_a^b in this expectation. Thus we have

$$U^{P(p)}(t) = e^{\zeta t} \mathbb{E} \left\{ \mathbb{E}_{0} \left[\left(\prod_{\ell=0}^{m} I_{R^{\ell}}^{S^{\ell+1}} \right) \left(\prod_{\ell=1}^{m} I_{S^{\ell}}^{R^{\ell}} \right) \mathbb{1}_{\{p(X_{t}) \in P(p)\}} \middle| \mathcal{F}_{S^{m+1}} \right] \times \mathbb{E}_{0} \left[I_{S^{m+1}}^{t} \mathbb{1}_{\{p(X_{t}) \in P(p)\}} \middle| \mathcal{F}_{S^{m+1}} \right] \right\}.$$
(51)

We use Corollary 6.2 to bound the expectation on the second line of (51); in the calculation that follows, abbreviate $s := s^{m+1}$ and $S := S^{m+1}$. We obtain, for some C > 1,

$$\mathbb{E}_{0}\left[I_{S}^{t}\mathbb{1}_{\{p(X_{t})\in P_{t}(p)\}} \middle| S\right] \leq \mathbb{1}_{\{S\leq t\}} \mathbb{E}_{p_{s}}\left[I_{0}^{t-S}\mathbb{1}_{\{\tau_{B(p_{s},j)}>t-S\}} \middle| S\right] \leq Ce^{(t-S)(\lambda^{(j)}(p_{s})-\zeta)} \leq Ce^{(t-S)(\lambda^{(j)}(p_{$$

almost surely, since $\zeta > \lambda^{(j)}(p_s)$. Combining with (51) and using equation (50) we obtain

$$U^{P(p)}(t) \leq C e^{\zeta t} \mathbb{E}_{0} \left[\left(\prod_{\ell=0}^{m} I_{R^{\ell}}^{S^{\ell+1}} \right) \left(\prod_{\ell=1}^{m} I_{S^{\ell}}^{R^{\ell}} \right) \mathbb{1}_{\{p(X)\in P(p)\}} \right]$$
$$= C e^{\zeta t} \mathbb{E}_{0} \left[\left(\prod_{\ell=0}^{m} I_{R^{\ell}}^{S^{\ell+1}} \right) \mathbb{1}_{\{p(X)\in P(p)\}} \right] \mathbb{E}_{0} \left[\left(\prod_{\ell=1}^{m} I_{S^{\ell}}^{R^{\ell}} \right) \mathbb{1}_{\{p(X)\in P(p)\}} \right].$$
(52)

Let $\xi_{\max}^{(\ell)} = \max_{r^{\ell} \leq k < s^{\ell+1}} \xi(p_k)$, for $\ell = \{0, 1, \dots, m\}$. By Lemma 3.4, which we can apply here since $\zeta > L_{t,\varepsilon} > L_t \geq \max_{0 \leq l \leq m} \xi_{\max}^{(\ell)}$,

$$\mathbb{E}_{0}\left[\left(\prod_{\ell=0}^{m} I_{R^{\ell}}^{S^{\ell+1}}\right) \mathbb{1}_{\{p(X)\in P(p)\}}\right] \leq (2d)^{-n} \prod_{\ell=0}^{m} \prod_{k=r^{\ell}}^{s^{\ell+1}-1} \left(1 + \sigma(p_{k})(\zeta - \xi_{\max}^{(\ell)})\right)^{-1} \qquad (53)$$

$$\leq (2d)^{-n} \left(1 + \delta_{\sigma}(\zeta - L_{t})\right)^{-n},$$

almost surely, using the definition of n and the lower bound on σ . Making the new abbreviation $s := s^{\ell}$, we have

$$\mathbb{E}_{0}\left[\left(\prod_{\ell=1}^{m} I_{S^{\ell}}^{R^{\ell}}\right) \mathbb{1}_{\{p(X)\in P(p)\}}\right] = \prod_{\ell=1}^{m} \mathbb{E}_{p_{s}}\left[I_{0}^{\tau_{B(p_{s},j)}} \mathbb{1}_{\{p(X)\in P(p)\}}\right] \leq \prod_{\ell=1}^{m} \mathbb{E}_{p_{s}}\left[I_{0}^{\tau_{B(p_{s},j)}}\right].$$

Since $\zeta > \lambda^{(j)}(z^{(p)})$, we can apply the first bound in the cluster expansion in Lemma 3.13 to deduce that

$$\prod_{\ell=1}^{m} \mathbb{E}_{p_s} \left[I_0^{\tau_{B(p_s,j)}} \right] \le \left(1 + \frac{\delta_{\sigma}^{-1} |B(0,j)|}{\zeta - \lambda^{(j)}(z^{(p)})} \right)^m.$$
(54)

Using these two estimates, we obtain from equation (52) the desired bound.

We now deal with the case that $s^{m+1} = \infty$. Similarly to the above, we condition on \mathcal{F}_{R^m} (the σ -algebra generated by R^m) to write $U^{P(p)}(t)$ as

$$e^{\zeta t} \mathbb{E}\bigg\{\mathbb{E}_0\left[\left(\prod_{\ell=0}^m I_{R^\ell}^{S^{\ell+1}}\right)\left(\prod_{\ell=1}^m I_{S^\ell}^{R^\ell}\right)\mathbb{1}_{\{p(X)\in P(p)\}}\Big|\mathcal{F}_{R^m}\right]\mathbb{E}_0\left[I_{R^m}^t\mathbb{1}_{\{R^m\leq t\}}\Big|\mathcal{F}_{R^m}\right]\bigg\}.$$

Set $l := |p| - r^m > 0$ and $\tau_{end} := \inf\{s > 0 : X_s = X_t\}$. Observe that, since $\zeta > L_{t,\varepsilon} > L_t \ge \xi(X_t)$, almost surely

$$\mathbb{E}_0\left[I_{R^m}^t \mathbb{1}_{\{p(X_t)\in P(p)\}} \middle| \mathcal{F}_{R^m}\right] \le \mathbb{E}_0\left[I_{R^m}^{\tau_{\text{end}}} \mathbb{1}_{\{p(X_t)\in P(p)\}} \middle| \mathcal{F}_{R^m}\right]$$

and applying Lemma 3.4 (valid by Lemma 3.2) we get that

$$\mathbb{E}_{0}\left[I_{R^{m}}^{t}\mathbb{1}_{\{p(X_{t})\in P(p)\}} \middle| \mathcal{F}_{R^{m}}\right] \leq (2d)^{-l}(1+\delta_{\sigma}(\zeta-L_{t}))^{-l}$$

almost surely. The rest of the proof proceeds similarly to the previous case.

We can use Lemma 6.6 to bound the contribution to the total mass U(t) from $U^2(t)$ and $U^3(t)$.

Proposition 6.7 (Upper bound on $U^2(t)$). There exists a constant c such that, as $t \to \infty$,

$$\log U^{2}(t) \le t \max_{z \in \Pi^{(L_{t})} \setminus \{Z_{t}\}} \Psi_{t,c}^{(j)}(z) + O(td_{t}b_{t})$$

almost surely.

Proof. Recall the path set $E_t^2,$ and for each $m,n\in\mathbb{N}$ define

$$\mathcal{P}_{n,m}^2 := \bigcup_{p \in E_t^2} P_t(p) \cap \mathcal{P}_{n,m}$$

Note that $|\mathcal{P}_{n,m}^2| \leq \kappa^{n+m}$, with $\kappa = \max\{2d, |\partial B(0,j)|\}$. We observe that

$$U^{2}(t) = \sum_{n,m} U^{\mathcal{P}^{2}_{n,m}}(t) \leq \sum_{n,m} \kappa^{n+m} \max_{P \in \mathcal{P}^{2}_{n,m}} \left\{ U^{P}(t) \right\} = \sum_{n,m} \kappa^{-n-m} \max_{P \in \mathcal{P}^{2}_{n,m}} \left\{ \kappa^{2(n+m)} U^{P}(t) \right\}$$
$$\leq \max_{n,m} \max_{P \in \mathcal{P}^{2}_{n,m}} \left\{ \kappa^{2(n+m)} U^{P}(t) \right\} \sum_{n,m} \kappa^{-n-m} .$$

For each $P \in \mathcal{P}_{n,m}^2$, denote by $z^{(P)}$ the site $y \in \Pi^{(L_t)}$ on a given path $p \in P$ which maximises $\lambda^{(j)}(y)$, remarking that this a class property of P eventually almost surely by Lemma 4.2. Using Lemma 6.6, for each $P \in \mathcal{P}_{n,m}^2$ and for any $\zeta > \max\{\lambda^{(j)}(z^{(P)}), L_{t,\varepsilon}\}$, we have that there exist constants $c_1, c_2, c_3 > 0$ such that, eventually almost surely,

$$\kappa^{2(n+m)} U^P(t) \le e^{\zeta t} \left(c_1(\zeta - L_t) \right)^{-n} \left(c_2 + c_3(\zeta - \lambda^{(j)}(z^{(P)}))^{-1} \right)^m.$$

Set $\zeta = \max\{\lambda^{(j)}(z^{(P)}), L_{t,\varepsilon}\} + d_t b_t$. To lower bound n, observe that the number of steps between exiting a *j*-ball and hitting another site in $\Pi^{(L_t)}$ is at least j + 1. We apply Corollary 4.4 to the balls $B(\Pi^{(L_t)}, j+1)$ to deduce that, eventually almost surely

$$n > m(j+1) + |z^{(P)}| - |z^{(P)}|^{c_4}, (55)$$

for some $c_4 < 1$. Then, by monotonicity in n,

$$\kappa^{2(n+m)} U^{P}(t) \leq e^{t(\lambda^{(j)}(z^{(P)}) + d_{t}b_{t})} (c_{1}(L_{t,\varepsilon} - L_{t}))^{-|z^{(P)}| + |z^{(P)}|^{c_{4}}} \times (c_{1}(L_{t,\varepsilon} - L_{t}))^{-j-1} (c_{2} + c_{3}d_{t}b_{t})^{-1}))^{m}$$

eventually almost surely. Note that j was chosen precisely to be the smallest integer such that

$$(j+1)\log a_t + \log(d_t) \to \infty \tag{56}$$

which implies, since $b_t \gg 1/\log\log t$ by (17), that

$$(j+1)\log a_t + \log(c_2 + c_3d_tb_t) \to \infty$$

By Lemma 4.2, for $z \in \Pi^{(L_t)}$, as $t \to \infty$,

$$|z|^{c_4} < td_t b_t$$

eventually almost surely. Moreover,

$$\log\left(L_{t,\varepsilon} - L_t\right) > \log a_t + c_5$$

eventually for some positive c_5 . So there exists a constant c such that

$$2(n+m)\log\kappa + \log U^{P}(t) \le c|z^{(P)}| + \lambda^{(j)}(z^{(P)})t - \frac{1}{\gamma}|z^{(P)}|\log\log t + td_{t}b_{t}$$

eventually almost surely, which yields the result.

Proposition 6.8 (Upper bound on $U^{3}(t)$). There exists a constant c such that, as $t \to \infty$,

$$\log U^{3}(t) \le t \Psi_{t,c}^{(j)}(Z_{t}) - h_{t} \frac{1}{\gamma} |Z_{t}| \log \log t + O(td_{t}b_{t})$$

almost surely.

Proof. Recall the set of paths E_t^3 and define $\mathcal{P}_{n,m}^3$ by analogy with $\mathcal{P}_{n,m}^2$. The proof then follows as for Proposition 6.7 after strengthening the bound in (55) to give that for each $p \in E_t^3$ and for some $c_1 < 1$, eventually almost surely

$$n > m(j+1) + (1+h_t)\frac{1}{\gamma}|Z_t|\log\log t - |Z_t|^{c_1}.$$

Proposition 6.9 (Upper bound on $U^4(t)$). For all $t \ge 0$,

$$U^4(t) < e^{tL_t}$$

Proof. This follows trivially from the definition of $U^4(t)$.

Proposition 6.10 (Negligibility of $U^{5}(t)$). As $t \to \infty$, almost surely,

$$\frac{U^5(t)}{U(t)} \to 0 \,.$$

Proof. The equivalent statement for the PAM with Weibull potential is proved in [14, Section 2.5], and is a consequence of a large probabilistic penalty for diffusing outside the macrobox V_t . The assumption that $\sigma(0) > \delta_{\sigma}$ ensures that the proof applies equally well in our case.

Corollary 6.11. There exists a constant c such that, as $t \to \infty$,

$$\frac{U^2(t) + U^3(t) + U^4(t) + U^5(t)}{U(t)} \mathbb{1}_{\mathcal{E}_{t,c}} \mathbb{1}_{\Theta_t^d} \to 0$$

almost surely.

Proof. Let c be the maximum of the constants appearing in Propositions 6.7 and 6.8. Combining Propositions 6.5 and 6.7, and recalling that $Z_{t,c}^{(j)} = Z_t$ eventually by Proposition 5.8 and Corollary 5.11, we have that, on the events $\mathcal{E}_{t,c}$ and Θ_t^d , eventually almost surely

$$\log U^{2}(t) - \log U(t) \le t \left(\Psi_{t,c}^{(j)}(Z_{t,c}^{(j,2)}) - \Psi_{t,c}^{(j)}(Z_{t,c}^{(j)}) \right) + c|Z_{t}| + O(td_{t}b_{t}).$$

Using the gap in the maximisers of $\Psi_{t,c}^{(j)}$ and since $|Z_t| < r_t g_t$, we have that, as $t \to \infty$,

$$\log U^2(t) - \log U(t) \le -td_t e_t + O(r_t g_t) + O(td_t b_t) \to -\infty$$

by the properties of the scaling functions in (17). Similarly, combining Propositions 6.5 and 6.8, we have that, on the events $\mathcal{E}_{t,c}$ and Θ_t^d , eventually almost surely

$$\log U^{3}(t) - \log U(t) \leq -h_{t} \frac{1}{\gamma} |Z_{t}| \log \log t + c|Z_{t}| + O(td_{t}b_{t})$$

and so, using that $|Z_t| > r_t f_t$ on the event $\mathcal{E}_{t,c}$, as $t \to \infty$,

$$\log U^{3}(t) - \log U(t) \leq -r_{t}f_{t}h_{t}\frac{1}{\gamma}\log\log t + O(td_{t}b_{t}) \to -\infty$$

by the properties in (17). Finally, combining Propositions 6.5, 6.9 and 6.10, we get the result. \Box

7. LOCALISATION

In this section we complete the proof of Theorem 1.7; that is, we show that the non-negligible component of the total solution, $u^1(t,z)$, is eventually localised at Z_t . Recall the idea of the proof that was outlined in Section 2, that: (i) the solution $u^1(t,z)$ is closely approximated by the principal eigenfunction of $\Delta \sigma^{-1} + \xi$ restricted to the domain

$$B_t := B(0, |Z_t|(1+h_t)) \cap V_t$$

and; (ii) the principal eigenfunction decays exponentially away from Z_t . Throughout this section, fix the constant c > 0 from Corollary 6.11.

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7.1. Approximating the solution with the principal eigenfunction. Let λ_t and v_t denote, respectively, the principal eigenvalue and eigenfunction of the Hamiltonian $\Delta \sigma^{-1} + \xi$ restricted to the domain B_t with Dirichlet boundary conditions, renormalising v_t so that $v_t(Z_t) = 1$.

Lemma 7.1 (Gap in *j*-local principal eigenvalues in B_t). On the event $\mathcal{E}_{t,c}$, each $z \in B_t \setminus \{Z_t\}$ satisfies

$$\lambda^{(j)}(Z_t) - \lambda^{(j)}(z) > d_t e_t + o(d_t e_t).$$

Proof. On the event $\mathcal{E}_{t,c}$, we have that $\lambda^{(j)}(Z_t) > a_t(1-f_t)$ and so the claim is true for $z \notin \Pi^{(L_t)}$ by Lemma 3.2. On the other hand, if $z \in \Pi^{(L_t)}$ then

$$d_t e_t < \Psi_t^{(j)}(Z_t) - \Psi_t^{(j)}(z) = \lambda^{(j)}(Z_t) - \lambda^{(j)}(z) + \frac{|z| - |Z_t|}{\gamma t} \log \log t \,.$$

To complete the proof, notice that, for each $z \in B_t$,

$$\frac{z|-|Z_t|}{\gamma t}\log\log t < \frac{r_t g_t h_t}{\gamma t}\log\log t = d_t g_t h_t \ll d_t e_t$$

since $g_t h_t \ll e_t$ by (17).

Corollary 7.2. Eventually on the event $\mathcal{E}_{t,c}$, each $z \in B_t \setminus \{Z_t\}$ satisfies

$$\lambda_t > \lambda^{(j)}(z) + d_t e_t + o(d_t e_t)$$

Proof. First note that, on the event $\mathcal{E}_{t,c}$, the ball $B(Z_t, j) \subseteq B_t$. Hence, by the domain monotonicity in Lemma 3.1, we have $\lambda_t \geq \lambda^{(j)}(Z_t)$, and so the result follows from Lemma 7.1.

Proposition 7.3 (Feynman-Kac representation for the principal eigenfunction). Eventually on the event $\mathcal{E}_{t,c}$,

$$v_t(z) = \frac{\sigma(z)}{\sigma(Z_t)} \mathbb{E}_z \left[\exp\left\{ \int_0^{\tau_{Z_t}} \left(\xi(X_s) - \lambda_t \right) \, ds \right\} \, \mathbb{1}_{\{\tau_{B_t^c} > \tau_{Z_t}\}} \right] \,,$$

where

$$\tau_{Z_t} := \inf\{t \ge 0 : X_t = Z_t\} \quad and \quad \tau_{B_t^c} := \inf\{t \ge 0 : X_t \notin B_t\}.$$

Proof. This is an application of Proposition 3.3, valid precisely because of Corollary 7.2. \Box

7.2. Exponential decay of the principal eigenfunction. Recall the partition of paths into equivalence classes in Section 6, the quantities r^{ℓ} and s^{ℓ} associated to each equivalence class, and, for $m, n \in \mathbb{N}$, the set of equivalence classes $\mathcal{P}_{n,m}$. Recall also the event $\{p(X) \in P(p)\}$.

Define the path set

$$\bar{E}_t^1 := \left\{ p \in E_t^1 : |p| = \min \left\{ i : p_i = Z_t \right\} \right\}$$

and for each $m, n \in \mathcal{N}$ define

$$\bar{\mathcal{P}}^1_{n,m} := \bigcup_{p \in \bar{E}^1_t} P_t(p) \cap \mathcal{P}_{n,m} \,.$$

Further, for each $P \in \overline{\mathcal{P}}_{n,m}^1$ and $y \in B_t$ define

$$v_t^P(y) := \frac{\sigma(y)}{\sigma(Z_t)} \mathbb{E}_y \left[\exp\left\{ \int_0^{\tau_{Z_t}} \left(\xi(X_s) - \lambda_t \right) \, ds \right\} \mathbb{1}_{\{p(X) \in P\}} \right].$$
(57)

For each $P \in \overline{\mathcal{P}}_{n,m}^1$ denote by $z^{(P)}$ the site $y \in \Pi^{(L_t)}$ on a given path $p \in P$, excluding the site Z_t , which maximises $\lambda^{(j)}(y)$, setting $z^{(P)} = \emptyset$ (and $\lambda^{(j)}(\emptyset) = 0$) if no such y exists. Remark that, whenever $z^{(P)}$ is defined, it is a class property of P eventually almost surely, by Lemma 4.2.

Lemma 7.4 (Bound on the contribution from each equivalence class). Let $m, n \in \mathbb{N}$ and $P \in \overline{\mathcal{P}}_{n,m}^1$. Then there exist constants $c_1, c_2 > 0$ such that, for each m, n, P and $y \in B_t \setminus \Pi^{(L_t)}$ uniformly, as $t \to \infty$,

$$v_t^P(y) \sigma(Z_t) \le (c_1(\lambda_t - L_t))^{-n} \left(1 + c_2(\lambda_t - \lambda^{(j)}(z^{(P)}))^{-1}\right)^{m-1}$$

and, for every $y \in \Pi^{(L_t)}$ uniformly,

$$v_t^P(y)\,\sigma(Z_t) \le \left(\lambda_t - \lambda^{(j)}(z^{(P)})\right)^{-1} \left(c_1(\lambda_t - L_t)\right)^{-n} \left(1 + c_2(\lambda_t - \lambda^{(j)}(z^{(P)}))^{-1}\right)^{m-1}$$

both hold eventually almost surely.

Proof. Starting with the Feynman-Kac representation for $v_t^P(y)$ in equation (57), the proof follows similarly as in Lemma 6.6 for $\zeta = \lambda_t$, which is a valid setting for ζ because of Corollary 7.2. Two modifications are necessary to adapt the proof.

The first modification comes from the observation that, for any $p \in P$, the final site Z_t gives no contribution to the expectation, and hence we have m-1 instead of the m in Lemma 6.6.

The second modification is necessary to take into account the additional $\sigma(y)$ factor present in the Feynman-Kac representation in equation (57), which a priori could be arbitrarily large. How we take this into account depends on whether p starts at a site of high potential. If $y \notin \Pi^{(L_t)}$, we simply modify equation (53) by pulling out the factor $\sigma(y)$ and bounding the right-hand side by

$$(2d)^{-n}\sigma^{-1}(y)(\lambda_t - L_t)^{-1} (1 + \delta_\sigma(\lambda_t - L_t))^{-n+1}$$

and the claimed result follows. If $y \in \Pi^{(L_t)}$, we instead modify equation (54) by using the second bound in Lemma 3.13 on the product factor for $\ell = 1$, which yields (abbreviating $s := s^{\ell}$)

$$\mathbb{E}_{y}[I_{0}^{\tau_{B}(y,j)}]\prod_{\ell=2}^{m-1}\mathbb{E}_{p_{s}}\left[I_{0}^{\tau_{B}(p_{s},j)}\right] \leq \sigma^{-1}(y)(\lambda_{t}-\lambda^{(j)}(z))^{-1}\left(1+\frac{\delta_{\sigma}^{-1}|B(0,j)|}{\lambda_{t}-\lambda^{(j)}(z^{(P)})}\right)^{m-1},$$
rain the claimed result follows.

and again the claimed result follows.

Proposition 7.5 (Exponential decay of principal eigenfunction). On the event $\mathcal{E}_{t,c}$ there exists a constant C > 0 such that, for each $y \in B_t$ uniformly, as $t \to \infty$,

$$\log v_t(y) + \log \sigma(Z_t) \le -C|y - Z_t| \log \log t$$

eventually almost surely.

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Proof. As in Proposition 6.7, we observe that there exists $\kappa > 1$ such that

$$v_t(y) = \sum_{n,m} \sum_{P \in \bar{\mathcal{P}}^1_{n,m}} v_t^P(y) \le \max_{n,m} \max_{P \in \mathcal{P}^1_{n,m}} \left\{ \kappa^{2(n+m)} v_t^P(y) \right\} \sum_{n,m} \kappa^{-n-m} \,.$$

Suppose $y \in B_t \setminus \Pi^{(L_t)}$. Then for each $P \in \overline{\mathcal{P}}_{n,m}^1$, by Lemma 7.4 there exist $c_1, c_2, c_3 > 0$ such that

$$\kappa^{2(n+m)}\sigma(Z_t)v_t^P(y) \le (c_1(\lambda_t - L_t))^{-n}(c_2 + c_3(\lambda_t - \lambda^{(j)}(z^{(P)}))^{-1})^{m-1}$$

eventually almost surely. Note also that by Corollary 4.4 (similarly to (55)), eventually almost surely

$$n > (m-1)(j+1) + c_4|y - Z_t|$$

for any $c_4 < 1$. Then, for any $0 < \varepsilon < \theta$,

$$\kappa^{2(n+m)}\sigma(Z_t)v_t^P(y) \le (c_1(L_{t,\varepsilon} - L_t))^{-c_4|y-Z_t|} \left((c_1(L_{t,\varepsilon} - L_t))^{-j-1}(c_2 + c_3(d_te_t)^{-1}) \right)^{m-1}$$

eventually almost surely by monotonicity in n and Corollary 7.2, and so, applying equation (56), there exists a C > 0 such that

$$2(n+m)\log\kappa + \log v_t^P(y) + \log\sigma(Z_t) \le -C|y-Z_t|\log\log t$$

eventually almost surely. Suppose then that $y \in \Pi^{(L_t)}$. Here we proceed similarly, but we now need the stronger bound $n > m(j+1) + c_4 |y - Z_t|$ for any $c_4 < 1$, valid eventually almost surely for $y \in \Pi^{(L_t)}$ by Lemma 4.2. Then,

$$\kappa^{2(n+m)}\sigma(Z_t)v_t^P(y) \le \left((c_1(L_{t,\varepsilon} - L_t))^{-j-1}(d_t e_t)^{-1} \right) (c_1(L_{t,\varepsilon} - L_t))^{-c_4|y-Z_t|} \\ \times \left((c_1(L_{t,\varepsilon} - L_t))^{-j-1}(c_2 + c_3(d_t e_t)^{-1}) \right)^{m-1},$$

and the rest of the proof follows as before.

7.3. Completion of the proof of Theorem 1.7. We are now in a position to establish Theorem 1.7. First, remark that Proposition 7.5 implies that, as $t \to \infty$,

$$\mathbb{1}_{\mathcal{E}_{t,c}} \, \sigma(Z_t) \sum_{z \in B_t \setminus \{Z_t\}} v_t(z)^2 \to 0$$

almost surely, and so in particular $\mathbb{1}_{\mathcal{E}_{t,c}} \|v_t\|_{\ell_2}^2 \to 1$, since we know $\sigma(Z_t) > \delta_{\sigma}$. Hence since

$$\mathbb{1}_{\mathcal{E}_{t,c}} \,\sigma(Z_t) \,\|\sigma^{-\frac{1}{2}} v_t\|_{\ell_2}^2 \,\sum_{z \in B_t \setminus \{Z_t\}} v_t(z) \,\leq\, \mathbb{1}_{\mathcal{E}_{t,c}} \,\delta_{\sigma}^{-1} \,\|v_t\|_{\ell_2}^2 \,\sigma(Z_t) \,\sum_{z \in B_t \setminus \{Z_t\}} v_t(z) \,, \tag{58}$$

the left-hand side of equation (58) also converges to zero almost surely. To finish the proof, we apply Proposition 3.14, which gives that

$$\mathbb{1}_{\mathcal{E}_{t,c}} \frac{1}{U(t)} \sum_{z \in B_t \setminus \{Z_t\}} u_1(t,z) \to 0$$

almost surely. Combining the above with the negligibility results already established in Corollary 6.11 on events $\mathcal{E}_{t,c}$ and Θ_t^d , and the fact that the events $\mathcal{E}_{t,c}$ and Θ_t^d hold eventually with overwhelming probability by Proposition 6.4, we have established Theorem 1.7.

References

- A. Astrauskas. Poisson-type limit theorems for eigenvalues of finite-volume Anderson Hamiltonian. Acta Appl. Math., 96:3–15, 2007.
- [2] A. Astrauskas. Extremal theory for spectrum of random discrete Schrödinger operator. I. Asymptotic expansion formulas. J. Stat. Phys., 131:867–916, 2008.
- [3] C. Bartsch, M. Kochler, and N. Gantert. Survival and growth of a branching random walk in random environment. Markov Process. Related Fields, 15(4):525–548, 2009.
- [4] G. Ben Arous and J. Černý. Dynamics of trap models. Math. Stat. Physics Lecture Notes Les Houches Summer School, 83, 2006.
- [5] M. Biskup and W. König. Screening effect due to heavy lower tails in one-dimensional parabolic Anderson model. J. Stat. Phys., 102(5-6):1253-1270, 2001.
- [6] T. Brotto, G. Bunin, and J. Kurchan. Population aging through survival of the fit and stable. J. Stat. Mech., 2016.
- [7] G. Chen and Z. Chen. Exact solutions of the position-dependent mass Schrödinger equation in D dimensions. *Phys. Lett. A.*, 331(5):312–315, 2004.
- [8] H. Eleuch, P. Jha, and Y. Rostovtsev. Analytical solution to position dependent mass for 3d-Schrödinger equation. Math. Sci. Lett., 1:1–6, 2012.
- [9] P. Embrechts, C. Klüppelberg, and T. Mikosch. Modelling extremal events, volume 33 of Applications of Mathematics. Springer-Verlag, 1997.
- [10] A. Fiodorov and S. Muirhead. Complete localisation and exponential shape of the parabolic Anderson model with Weibull potential. *Electron. J. Probab.*, 19(58):1–27, 2014.
- [11] L.R.G. Fontes, M. Isopi, and C.M. Newman. Random walks with strongly inhomogeneous rates and singular diffusions: Convergence, localization and aging in one dimension. Ann. Probab., 30(2):579–604, 2002.
- [12] J. Gärtner, W. König, and S.A. Molchanov. Geometric characterization of intermittency in the parabolic Anderson model. Ann. Probab., 35(2):439–499, 2007.
- [13] J. Gärtner and S.A. Molchanov. Parabolic problems for the Anderson model. I. Intermittency and related topics. Comm. Math. Phys., 132:613–655, 1990.
- [14] J. Gärtner and S.A. Molchanov. Parabolic problems for the Anderson model. II. Second-order asymptotics and structure of high peaks. *Probab. Theory Related Fields*, 111:17–55, 1998.
- [15] G. Grimmett. Percolation, volume 321 of Grundlehren der Mathematischen Wissenschaften. Springer-Verlag, 1999.
- [16] R. Van der Hofstad, W. König, and P. Mörters. The universality classes in the parabolic Anderson model. Commun. Math. Phys., 267:307–353, 2006.
- [17] R. Van der Hofstad, P. Mörters, and N. Sidorova. Weak and almost sure limits for the parabolic Anderson model with heavy tailed potentials. Ann. Appl. Probab., 18(6):2450–2494, 2008.
- [18] K. Ishii, H. Matsuda, Y. Iwasa, and A. Sasaki. Evolutionarily stable mutation rate in a periodically changing environment. *Genetics*, 121(1):163–74, 1989.
- [19] J. F. C. Kingman. A simple model for the balance between selection and mutation. J. Appl. Prob., 15:1–12, 1978.
- [20] W. König, H. Lacoin, P. Mörters, and N. Sidorova. A two cities theorem for the parabolic Anderson model. Ann. Probab., 37(1):347–392, 2009.
- [21] E. G. Leigh. Natural selection and mutability. The American Naturalist, 104(937):301–305, 1970.

- [22] S. Muirhead. Two-site localisation in the Bouchaud trap model with slowly varying traps. Electron. Commun. Probab., 20(25):115, 2014.
- [23] S. Muirhead and R. Pymar. Localisation in the Bouchaud–Anderson model. Stoch. Proc. Appl., 126:3402–3462, 2016.
- [24] S. C. Park, D. Simon, and J. Krug. The speed of evolution in large asexual populations. J. Stat. Phys., 138:381–410, 2010.
- [25] O. von Roos. Position-dependent effective masses in semiconductor theory. Phys. Rev. B, 27(7545), 1983.
- [26] N. Sidorova and A. Twarowski. Localisation and ageing in the parabolic Anderson model with Weibull potential. Ann. Probab., 42(4):1666–1698, 2014.
- [27] F. Taddei, M. Radman, J. Maynard-Smith, B. Toupance, P. H. Gouyon, and B. Godelle. Role of mutator alleles in adaptive evolution. *Nature*, 387(6634):700-702, 1997.
- [28] T. Wolff. Random walk local times, Dirichlet energy and effective conductivity in the random conductance model. PhD Thesis, TU Berlin, 2013.